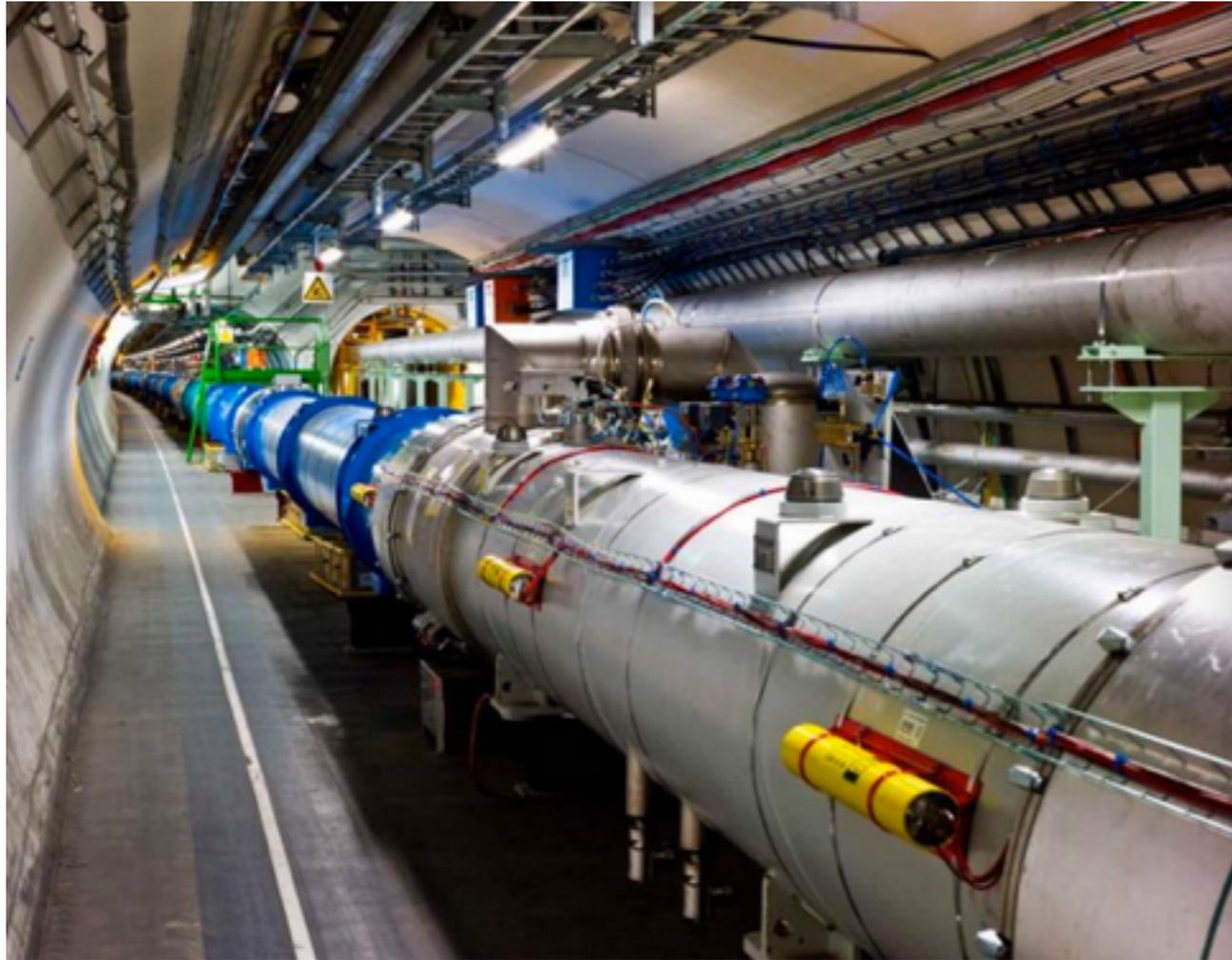


# **Strongly coupled gauge theories: What can lattice calculations teach us?**

**Anna Hasenfratz  
University of Colorado Boulder  
SCGT-2015, Nagoya, Mar 3, 2015**

# What is Beyond the Standard Model ?



The LHC will restart this month - will it reveal the nature of the Higgs ?

“With this new energy level, the LHC will open new horizons for physics and for future discoveries,” says CERN Director-General Rolf Heuer. “I’m looking forward to seeing what nature has in store for us”.  
(Feb 2015)



# Composite Higgs

is viable possibility:

Higgs is a  $\bar{q}q$  bound state (possibly  $qq$  )

- What models are compatible with EW data?
  - Most likely strongly coupled
- What are the generic properties of strongly coupled models?
  - is walking necessary ?
  - spectrum : where is  $M_{0^{++}}$  compared to  $M_\rho$  ?



# Strongly coupled models

We need only three Goldstones — 2 massless fermions will do

- $N_f = 2$  SU(3) : QCD
- $N_f = 2$  SU(2) adjoint : conformal
- $N_f = 2$  SU(3) sextet : popular but is it indeed chirally broken?
  - **Poster:** RG  $\beta$  function with Wilson fermions disagree with staggered
  - discrepancy could be due to rooting (?) or strong coupling effects
    - **needs better understanding**



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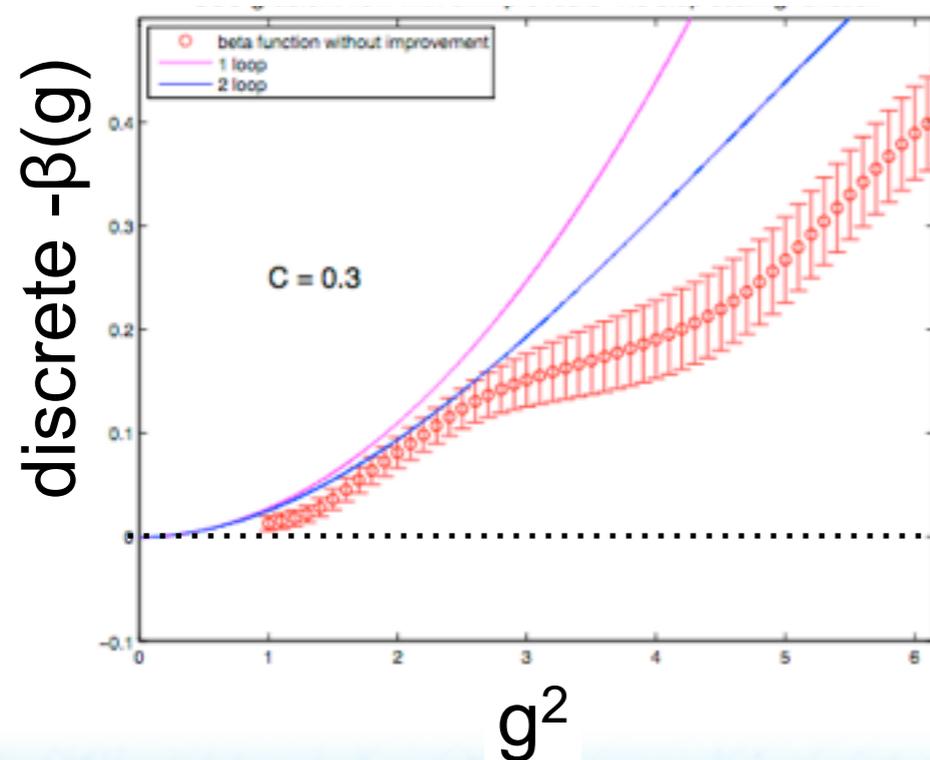
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# Strongly coupled models

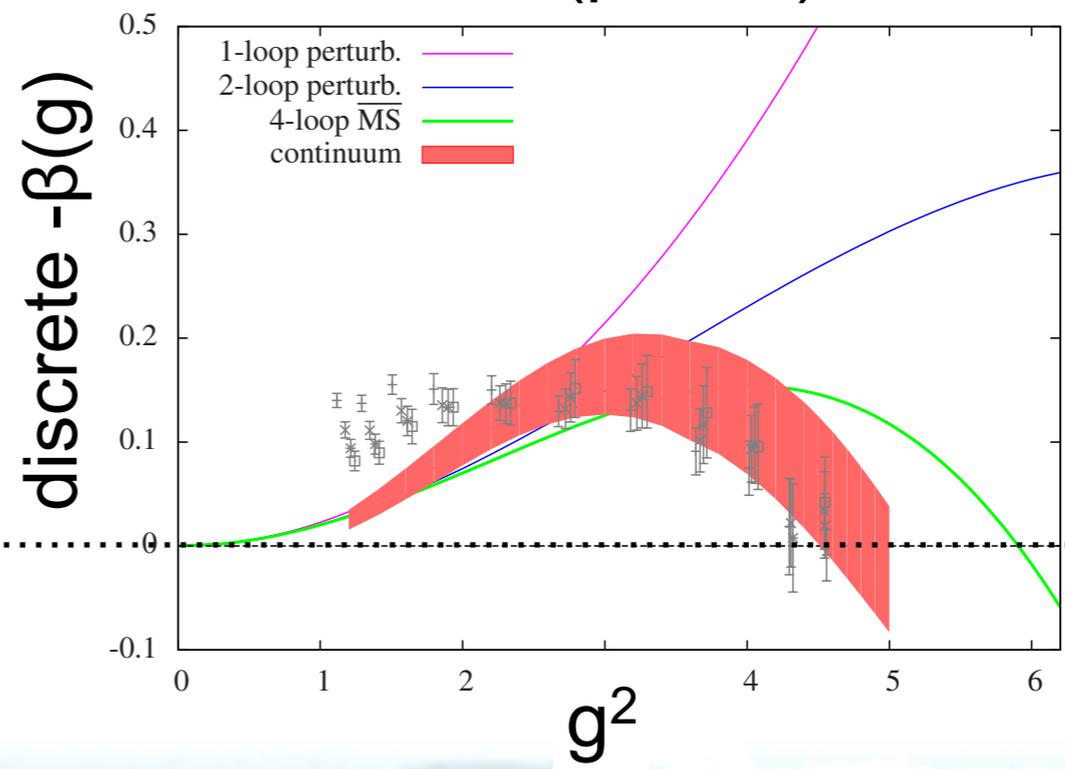
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Staggered (1502.00028)



Wilson (poster)



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??

If not  $N_f = 2$  :

- $N_f = 6$  SU(2) fundamental (1313.4889 - LSD)
- $N_f = 8$  SU(3) fundamental : seems to be close to the conformal window : **E. Rinaldi talk; D. Schaich finite T poster**

We need some mechanism to break flavor

$$SU(8) \times SU(8) \rightarrow SU(2) \times SU(2)$$

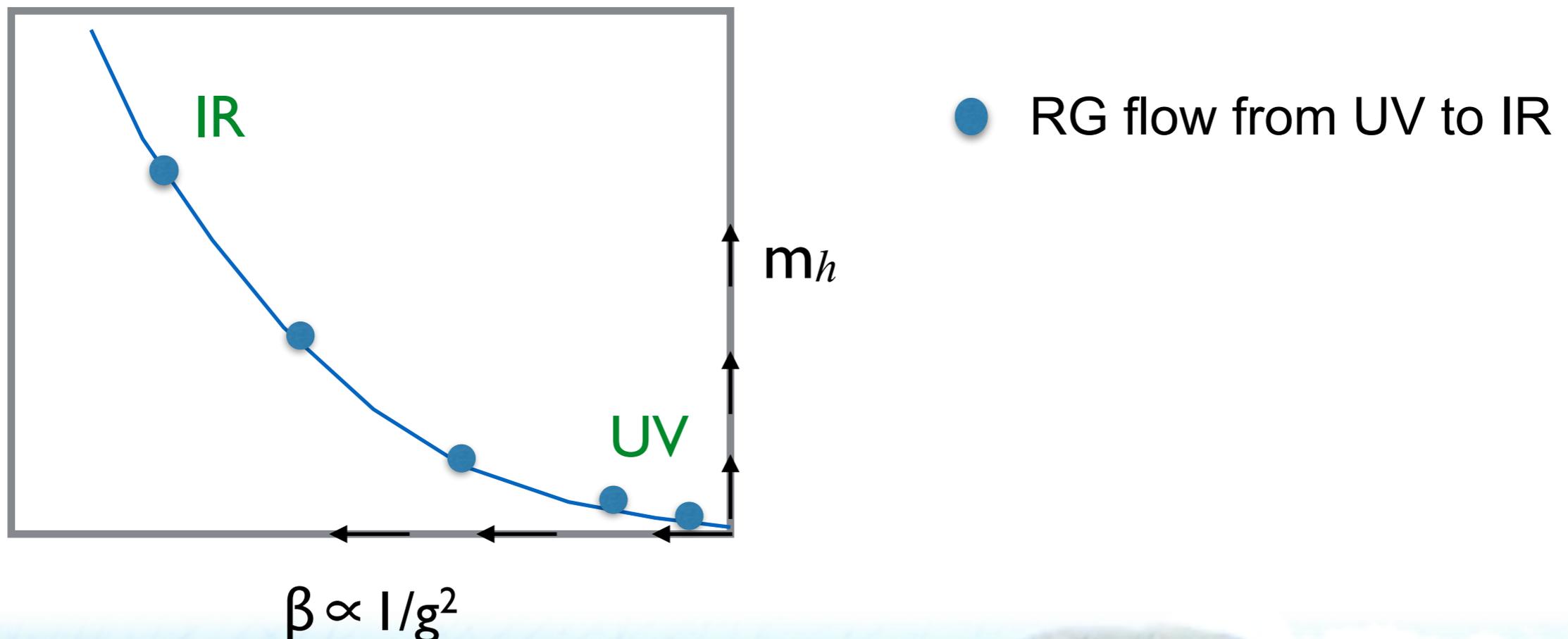
What is the remnant of the many flavors in the IR?



# Simple model - I

SU( $N_c$ ) gauge with  $N_\ell$  light ( $m_\ell \approx 0$ ) and  $N_h$  heavy ( $m_h$ ) fermions  
In the IR the heavy flavors decouple,  $N_\ell$  light remain

$N_\ell + N_h = \text{small}$ : gauge coupling runs fast, heavy flavors have limited effect on the IR (QCD)

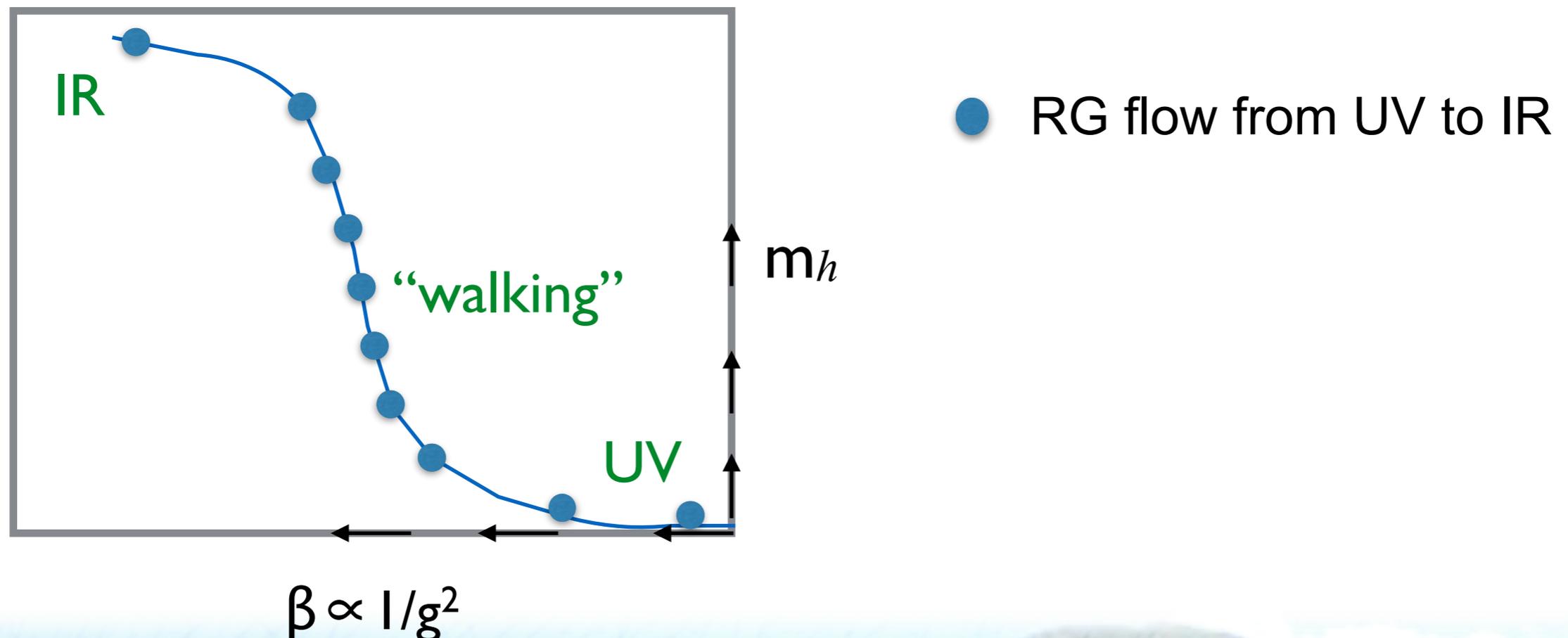


## Simple model - II

SU( $N_c$ ) gauge with  $N_\ell$  light ( $m_\ell \approx 0$ ) and  $N_h$  heavy ( $m_h$ ) fermions

$N_\ell + N_h = \text{near}$  but below the conformal window

**IF** the gauge coupling is “walking” the IR can be very different

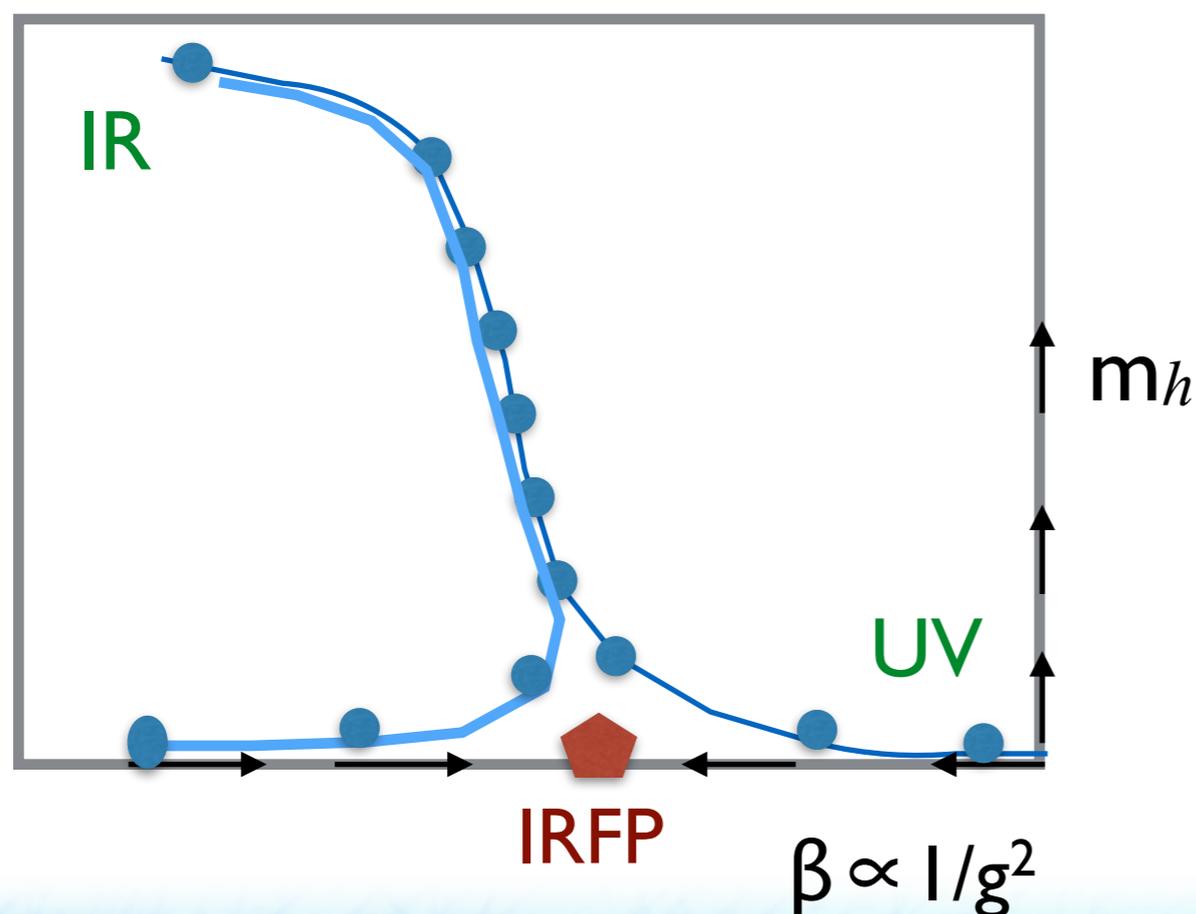


# Simple model - III

SU( $N_c$ ) gauge with  $N_\ell$  light ( $m_\ell \approx 0$ ) and  $N_h$  heavy ( $m_h$ ) fermions

$N_\ell + N_h =$  above the conformal window

the gauge coupling will be “walking”;  
the IR will be very different

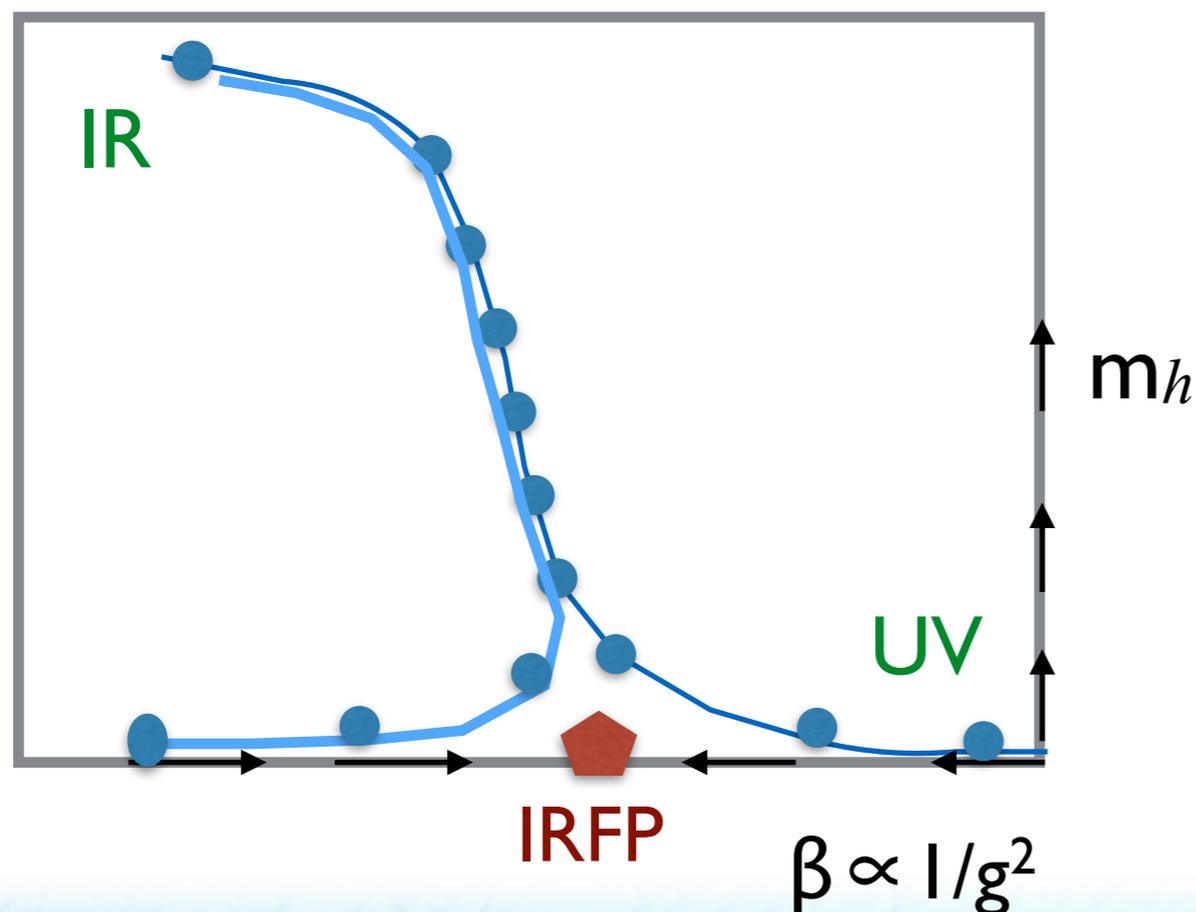


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What are the properties  
of these strongly coupled  
“walking” systems?

## $N_\ell + N_h$ (lattice) models

$N_\ell + N_h = 2 + 6$  if  $N_f = 8$  is the UV model

or

$N_\ell + N_h = 2 + 10$  for  $N_f = 12$  conformal behavior in the UV

Pilot study:

$N_\ell + N_h = 4 + 8$  : conformal in the UV,  $N_f = 4$  flavor in the IR

in collaboration with R. Brower, C. Rebbi, E. Weinberg, O. Witzel

arXiv:1411.3243

Why  $4+8$  ? We use staggered fermions:

4 and 8 flavors do not require rooting

(rooting is no-go in a conformal system near IRFP)

# $N_\ell + N_h = 4 + 8$ : The lattice action

Action: nHYP smeared staggered fermions,  
fundamental + adjoint gauge plaquette

This action was used in the Boulder 4, 8, and 12 flavor studies  
(1106.5293, 111.2317, 1404.0984)

It is the action used in the 8 flavor joint project with LSD  
(E. Rinaldi's talk, D. Schaich's poster)

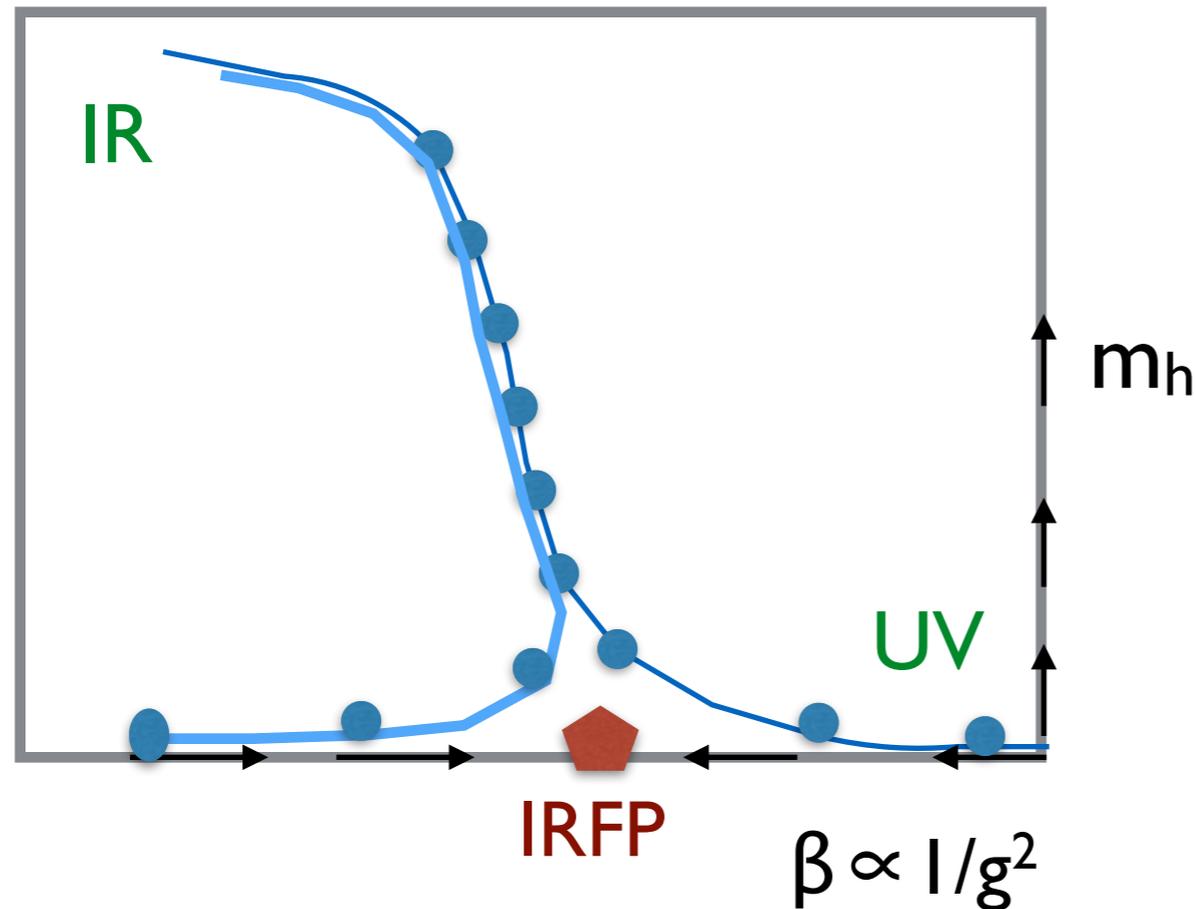
We understand this action well



# $N_\ell + N_h = 4 + 8$ : Parameter space

3 independent parameters:  $(g^2, m_\ell, m_h)$

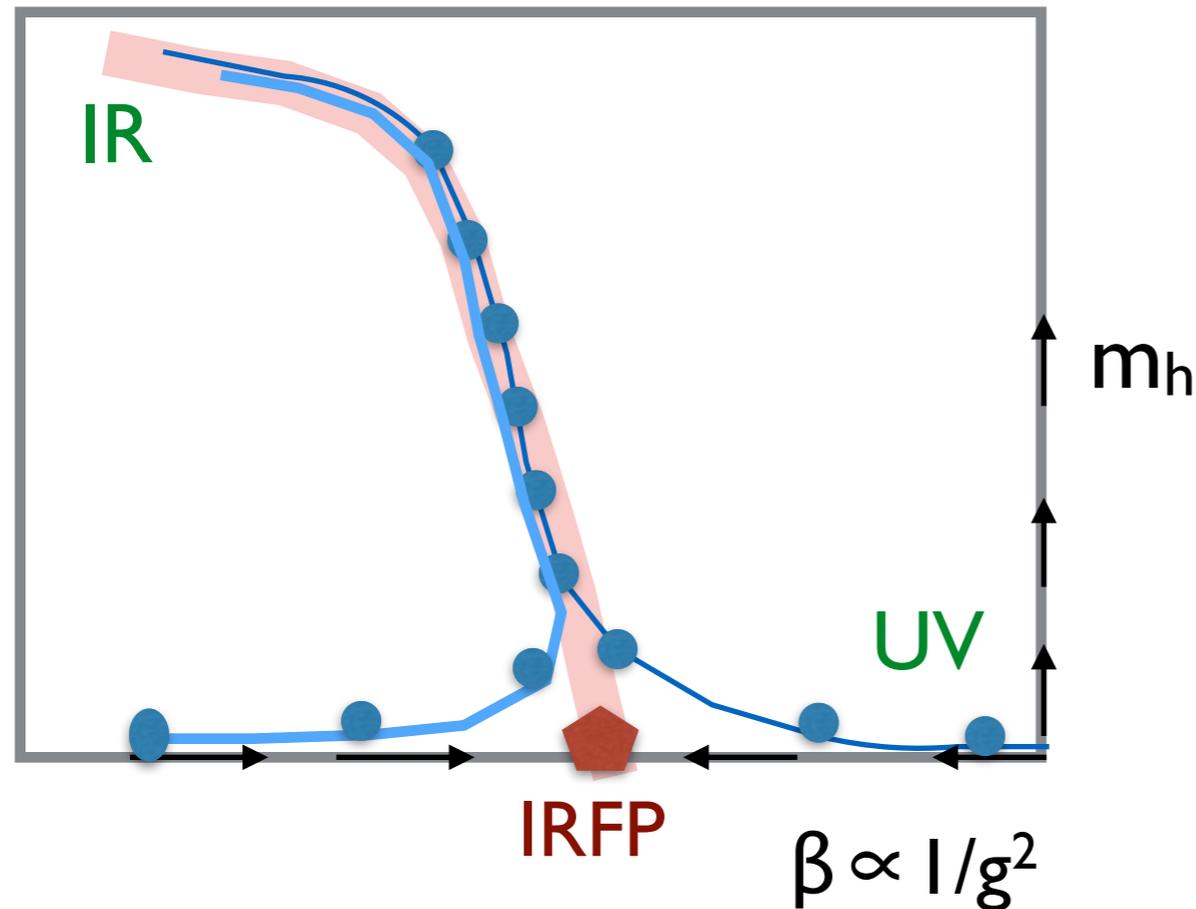
- $g^2$  does not matter once the flow has reached the RG trajectory
- sufficient to work at  $g^2 = \text{const}$ , vary  $m_h$  only



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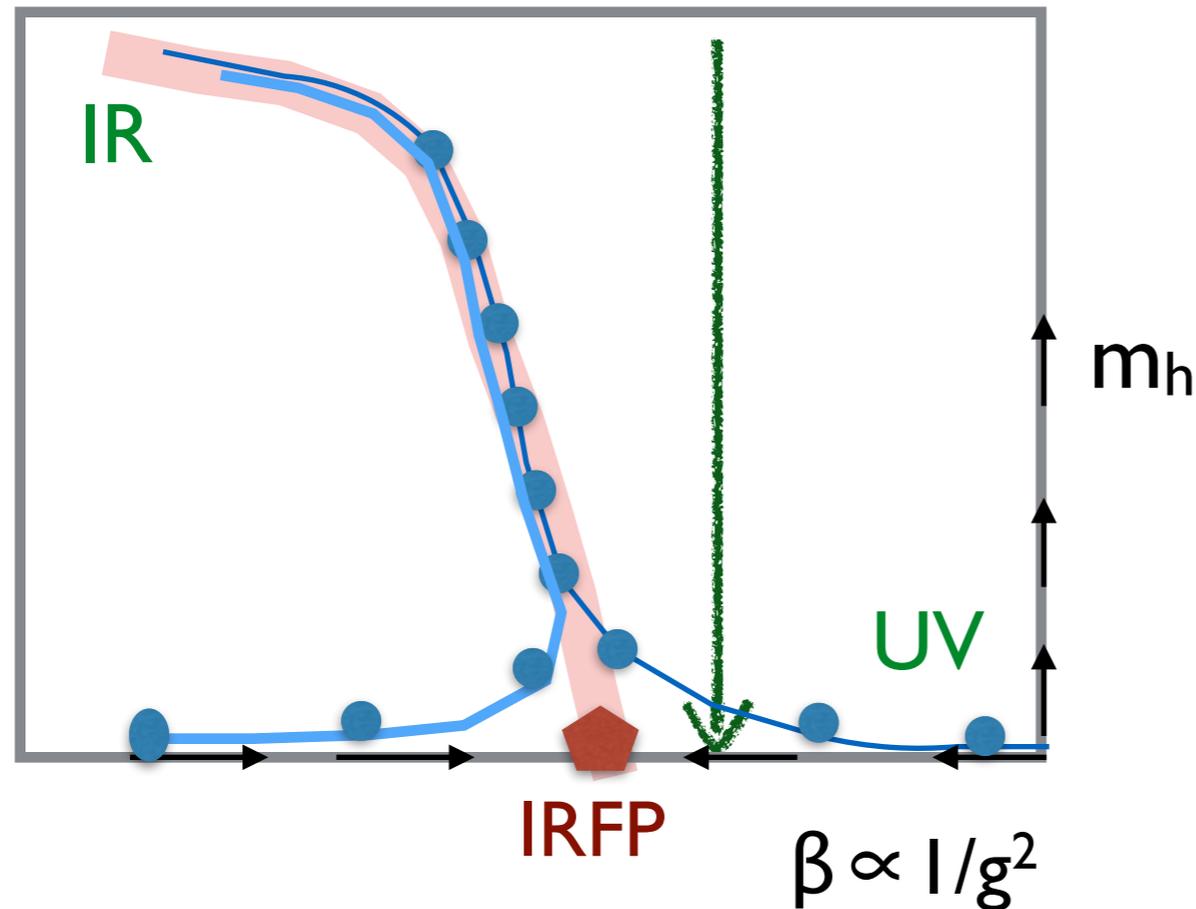
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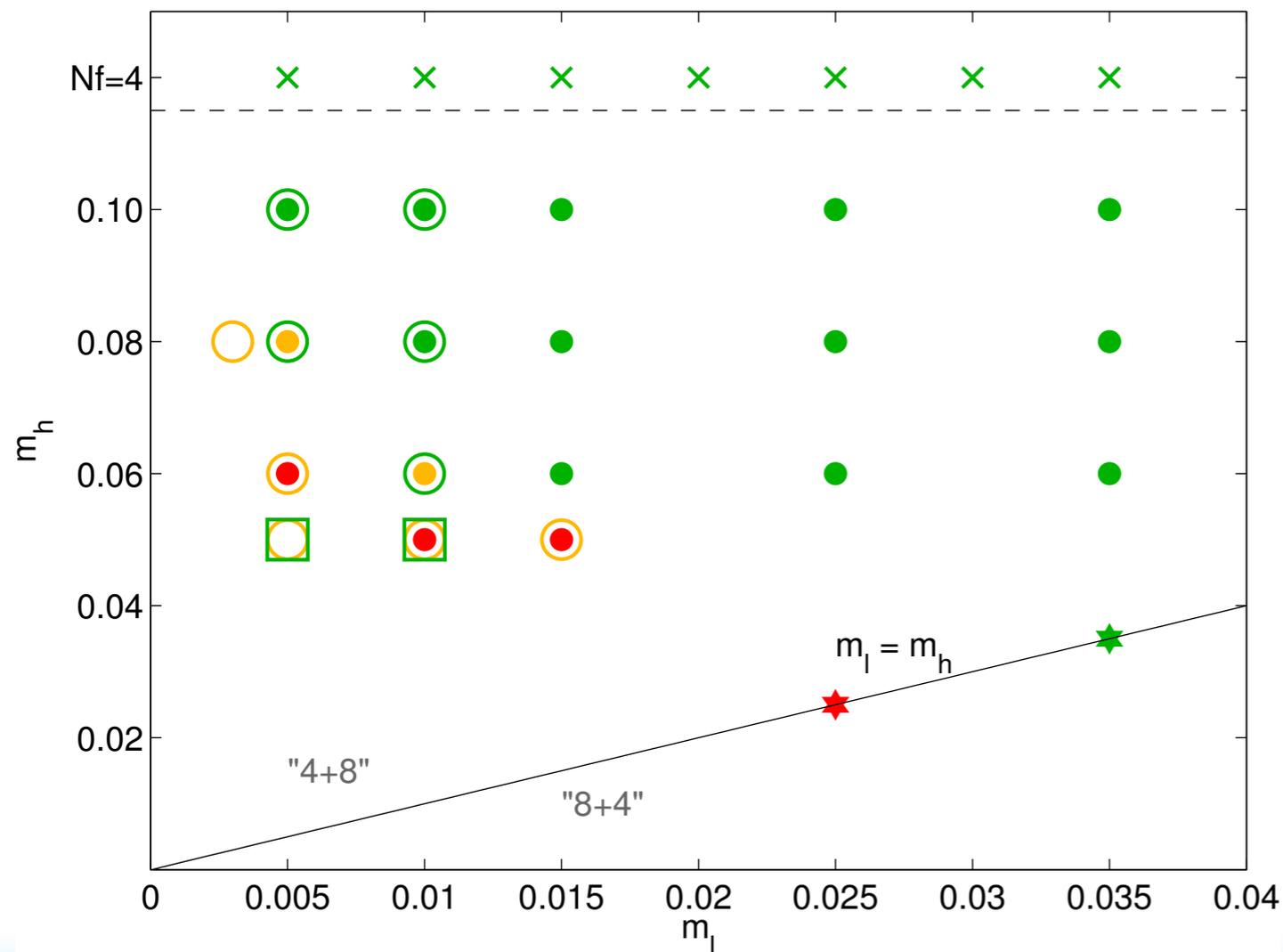
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# $N_\ell + N_h = 4 + 8$ : Parameter space

- $\beta = 4.0$  (close to the 12-flavor IRFP)
- $m_h = 0.10, 0.08, 0.06, 0.05$
- $m_\ell = 0.003, 0.005, 0.010, 0.015, 0.025, 0.035$



## Volumes :

$24^3 \times 48$ , (dots)

$32^3 \times 64$  (circle),

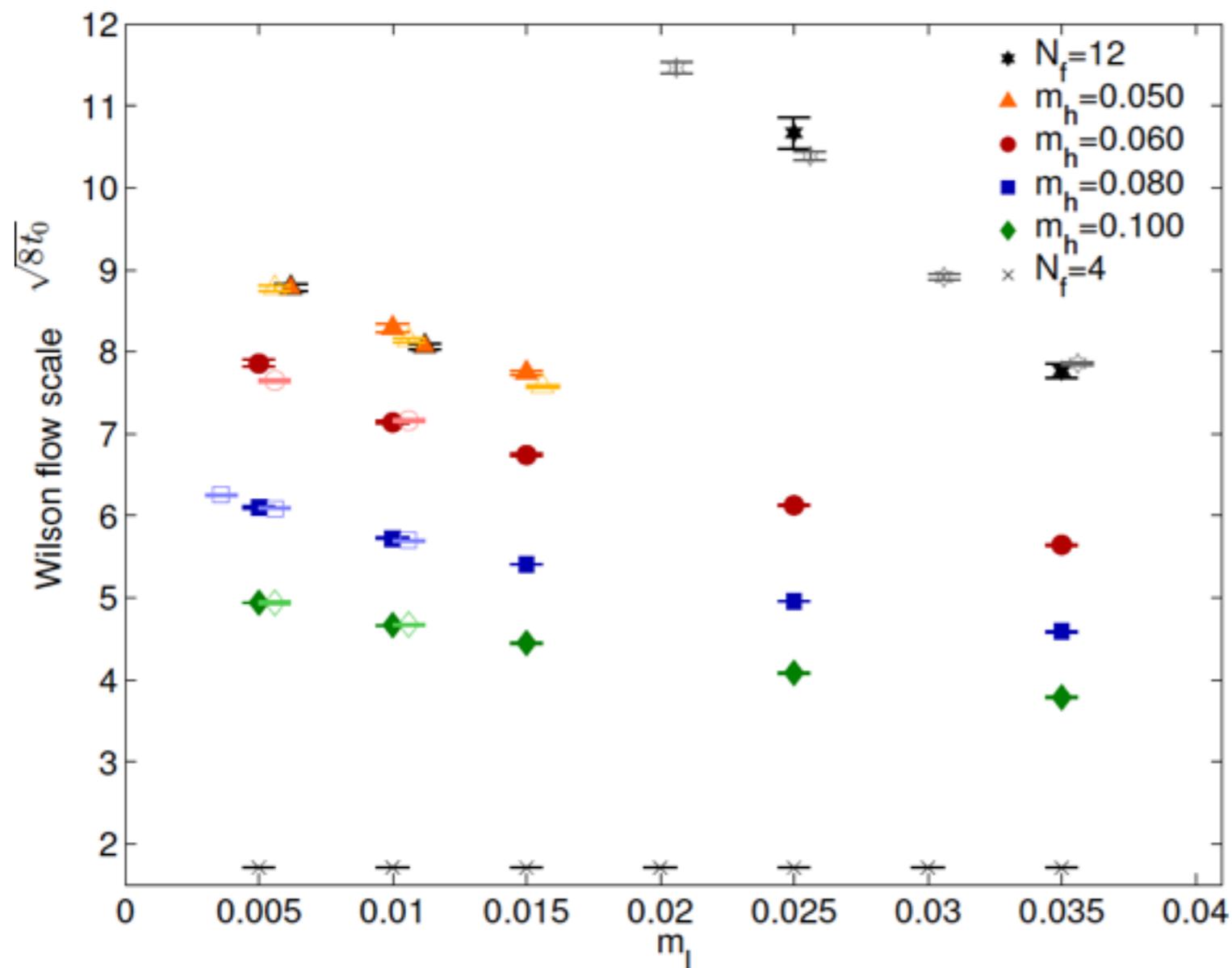
$48^3 \times 96$  (square)

Color: volume OK / marginal / squeezed

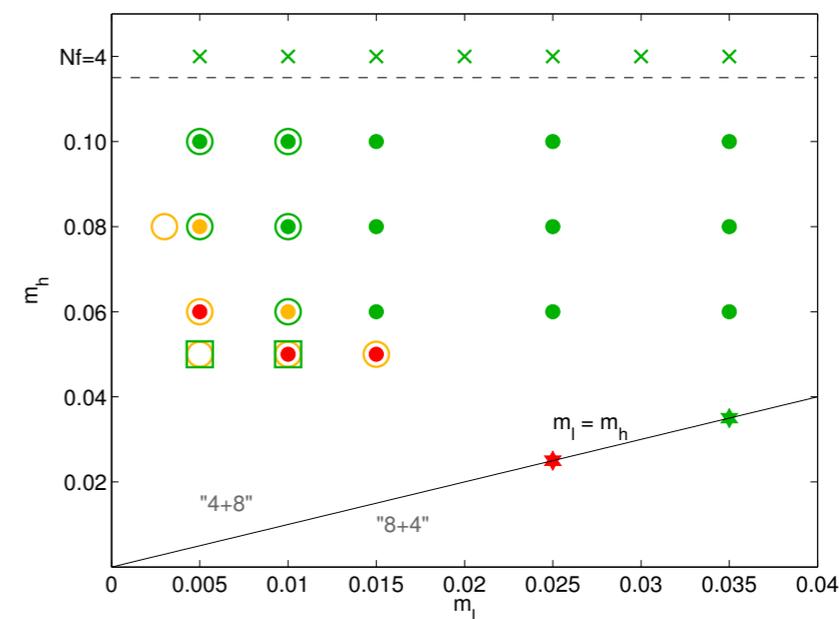
20,000 MDTU, most still in progress

# Lattice scale

Use Wilson flow to estimate the lattice scale  $\sqrt{8t_0}$

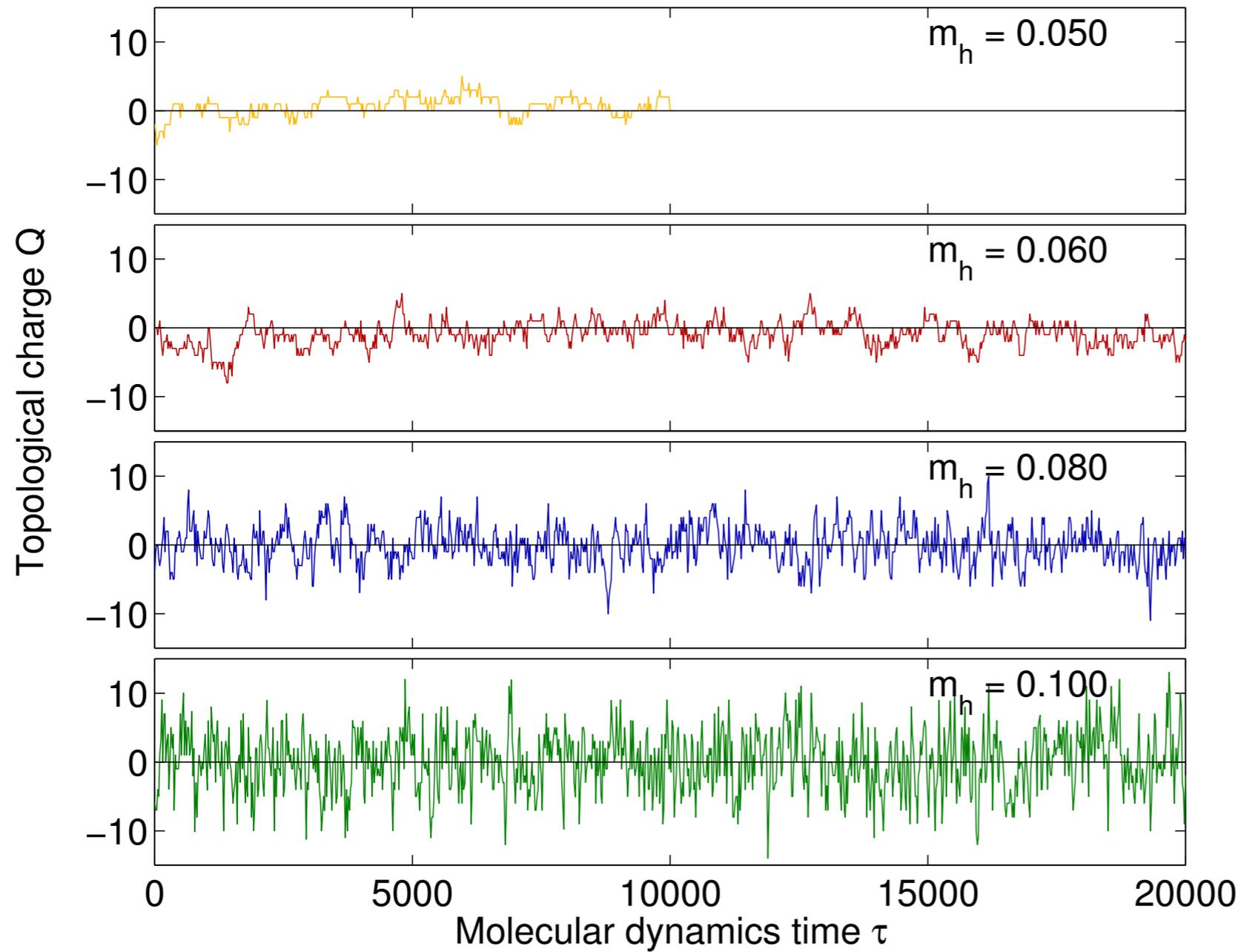


$\sqrt{8t_0} \lesssim L/5$   
 is usually sufficient  
 → color coding



# Topology evolution

Topology is moving well even with the lightest mass



$m_\ell = 0.010$ ,  
 $24^3 \times 48$  volume

# Running coupling

Gradient flow transformation defines a renormalized coupling

arXiv:1006.4518

$$g_{GF}^2(\mu = \frac{1}{\sqrt{8t}}) = \frac{1}{\mathcal{N}} t^2 \langle E(t) \rangle$$

t: flow time;  
E(t):energy density

$g_{GF}^2$  is used for scale setting as

$$g_{GF}^2(t = t_0) = \frac{0.3}{\mathcal{N}}$$

Is it appropriate for renormalized running coupling?

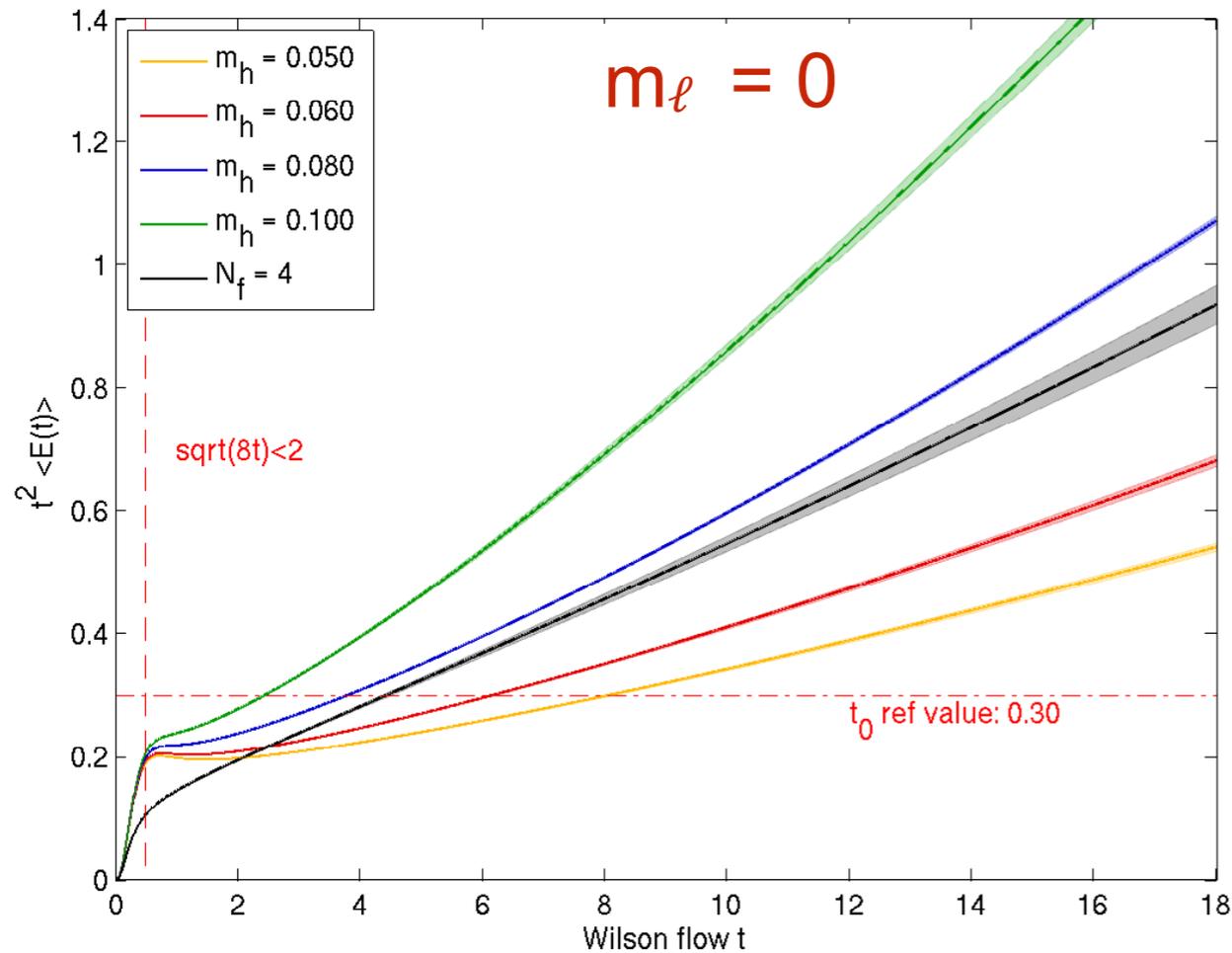
Yes,

- on large enough volumes
- at large enough flow time
- in the continuum limit

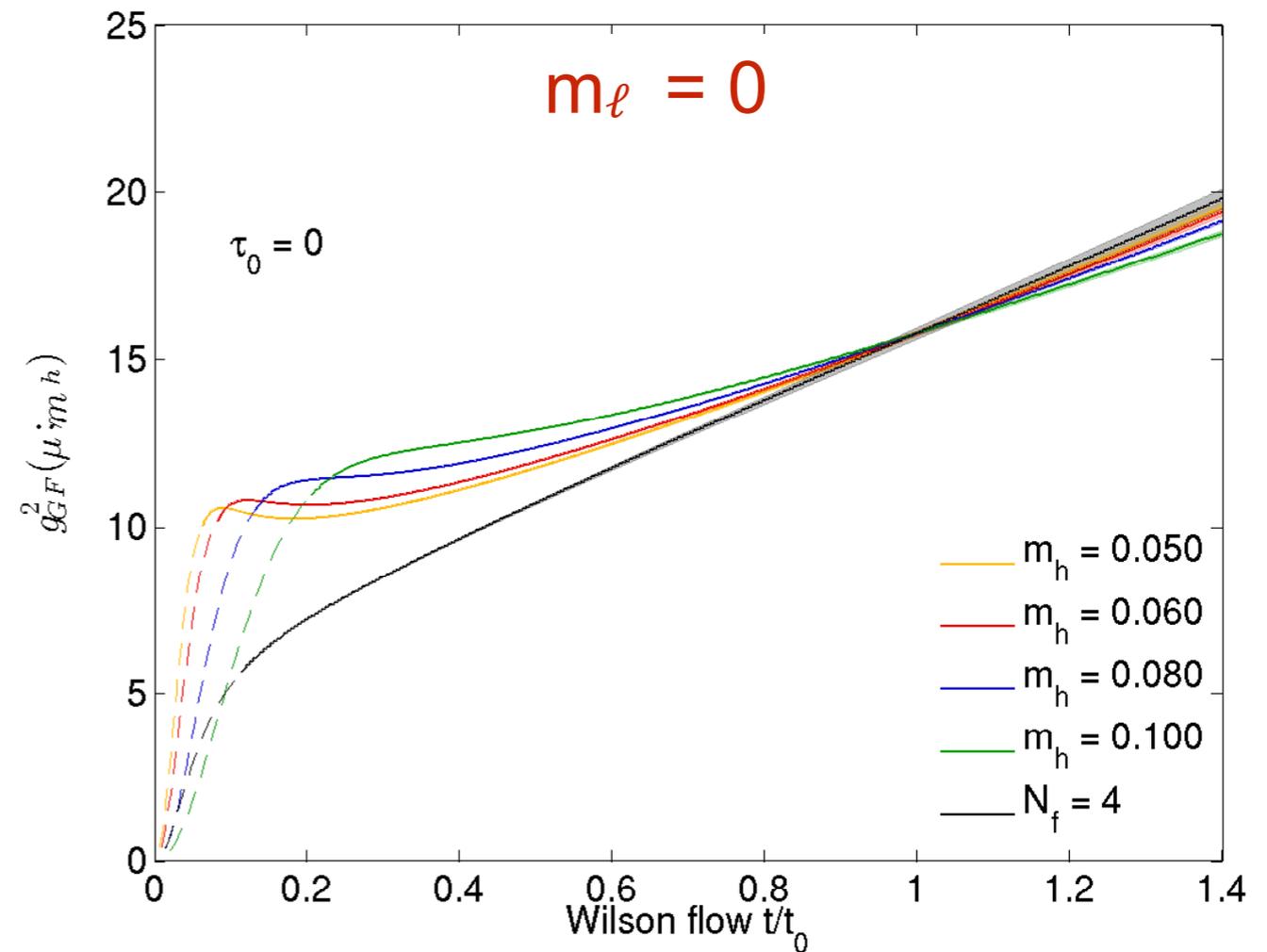


# Running coupling

$t^2 \langle E(t) \rangle$  in the chiral limit  
at various  $m_h$  values



$g_{GF}^2(t/t_0)$  rescaled by  $t_0$   
at various  $m_h$  values



Rescaling forces the renormalized couplings to agree at  $t_0$   
Fan-out before and after are due to cut-off lattice artifacts

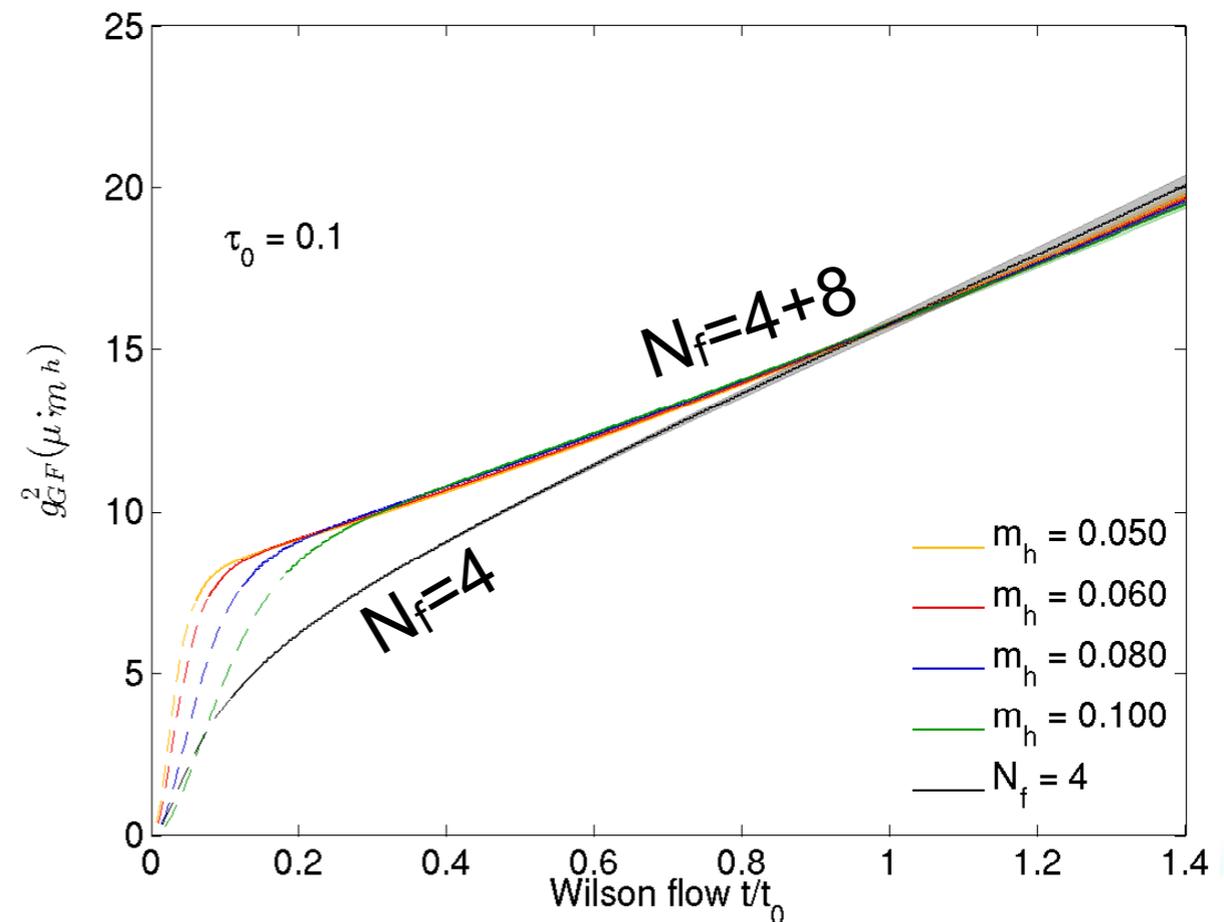
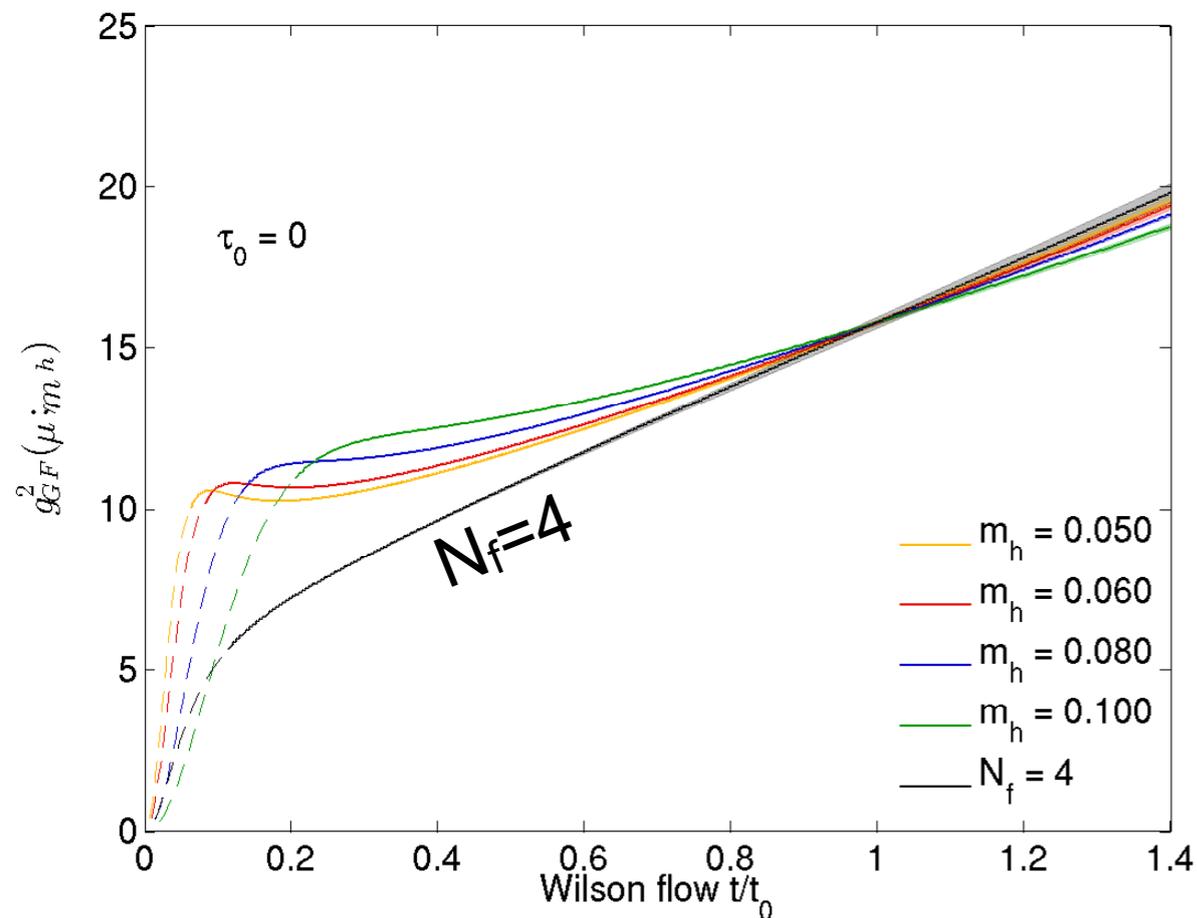
# Improved running coupling

t-shift improved running coupling

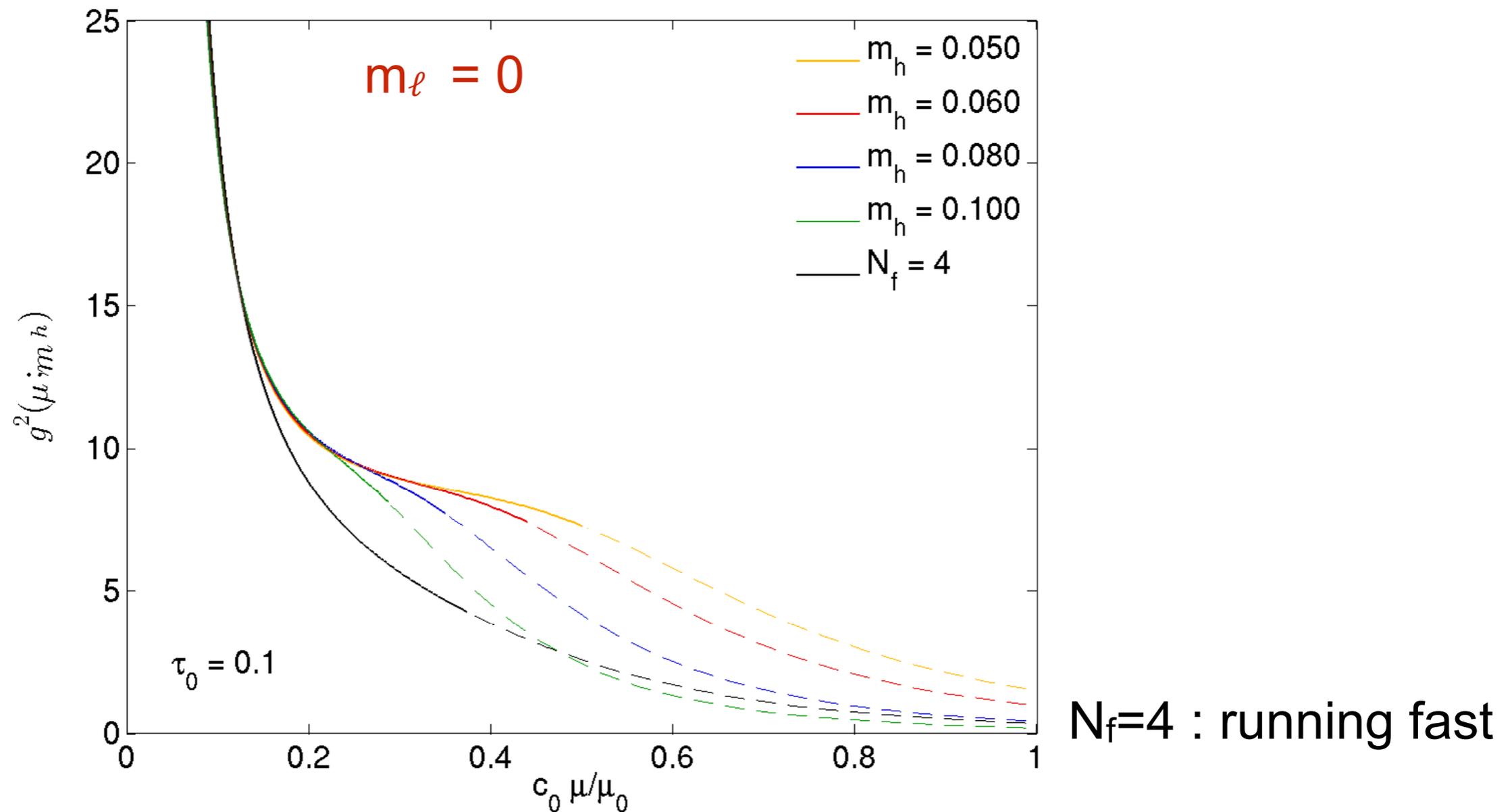
$$\tilde{g}_{GF}^2(\mu = \frac{1}{\sqrt{8t}}) = \frac{1}{\mathcal{N}} t^2 \langle E(t + \tau_0) \rangle$$

by adjusting  $\tau_0$  most cut-off effects can be removed

(1404.0984, 1501.07848)



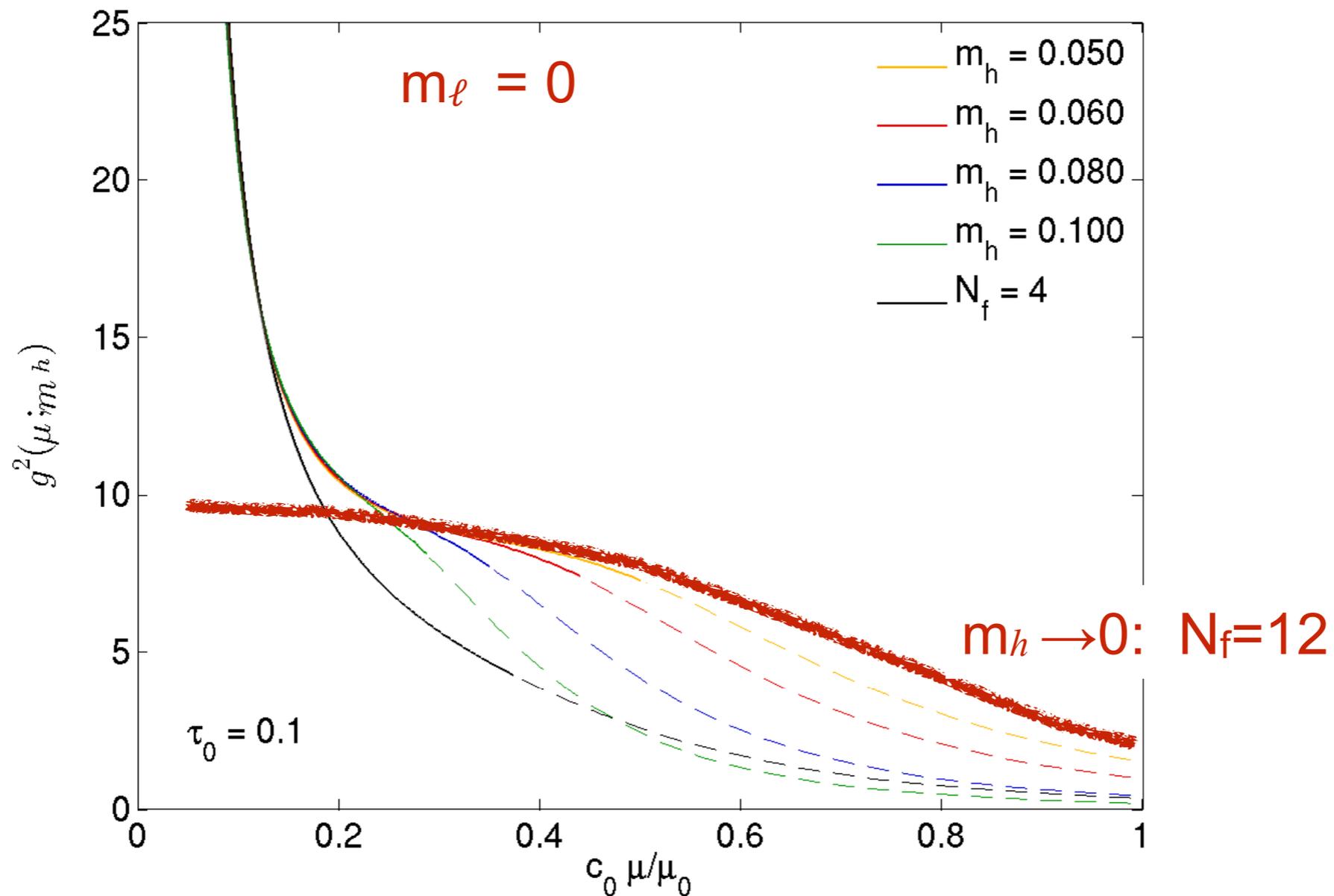
# Improved running coupling : 4+8 flavors



$g_{GF}^2(\mu)$  develops a “shoulder” as  $m_h \rightarrow 0$  : this is walking !

Walking range can be tuned arbitrarily with  $m_h$

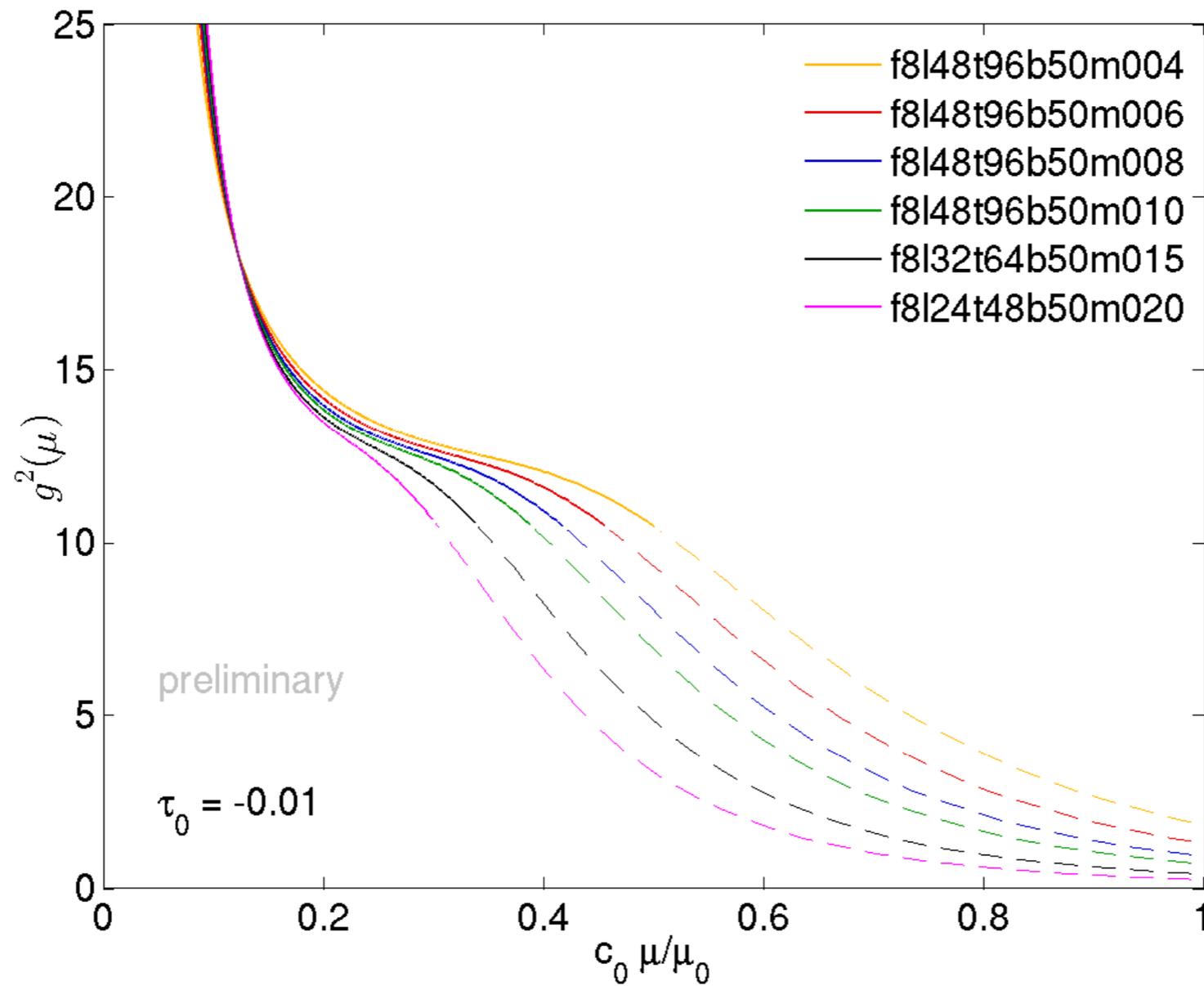
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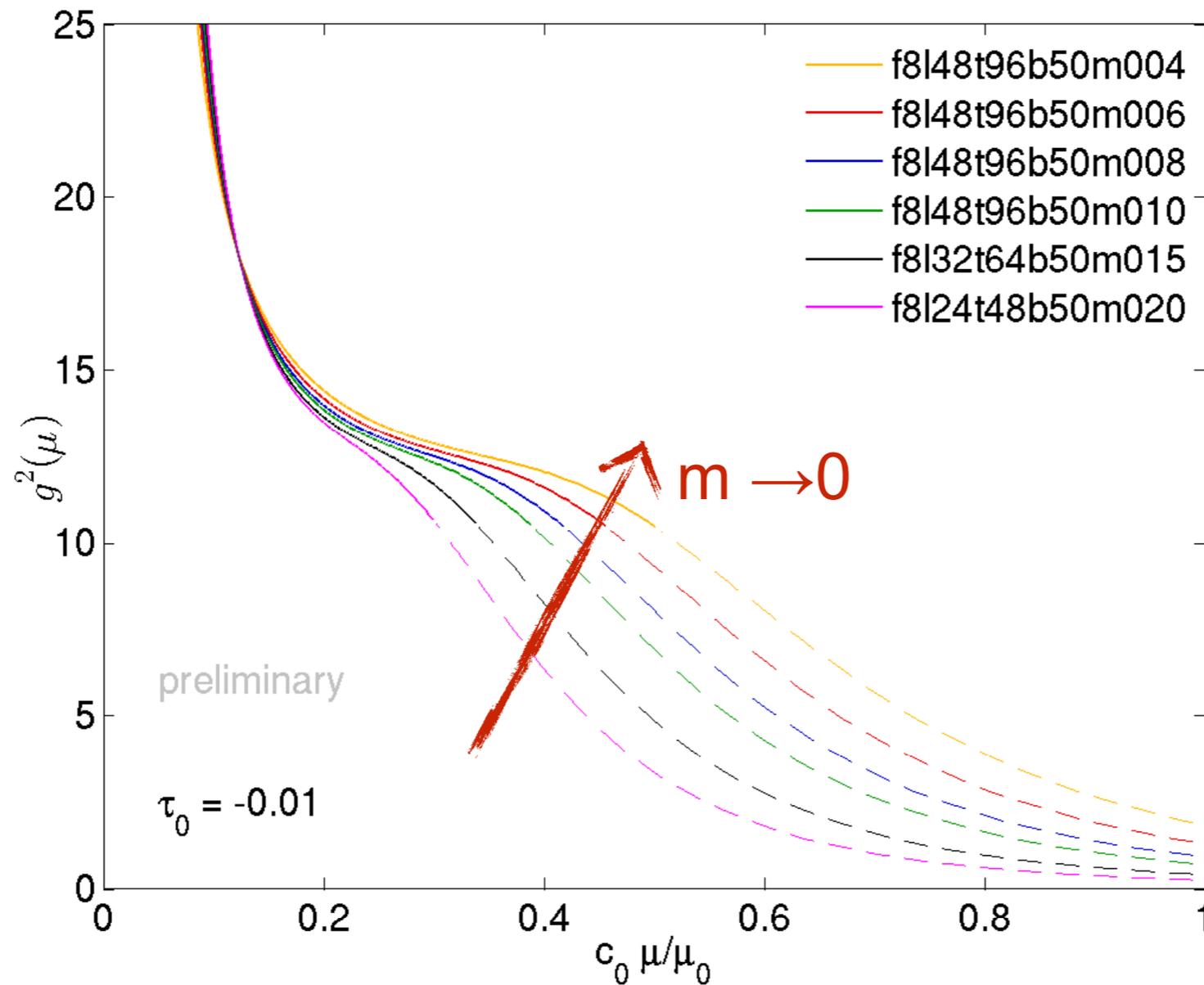
# Improved running coupling : 8 flavors



Is this walking?



# Improved running coupling : 8 flavors



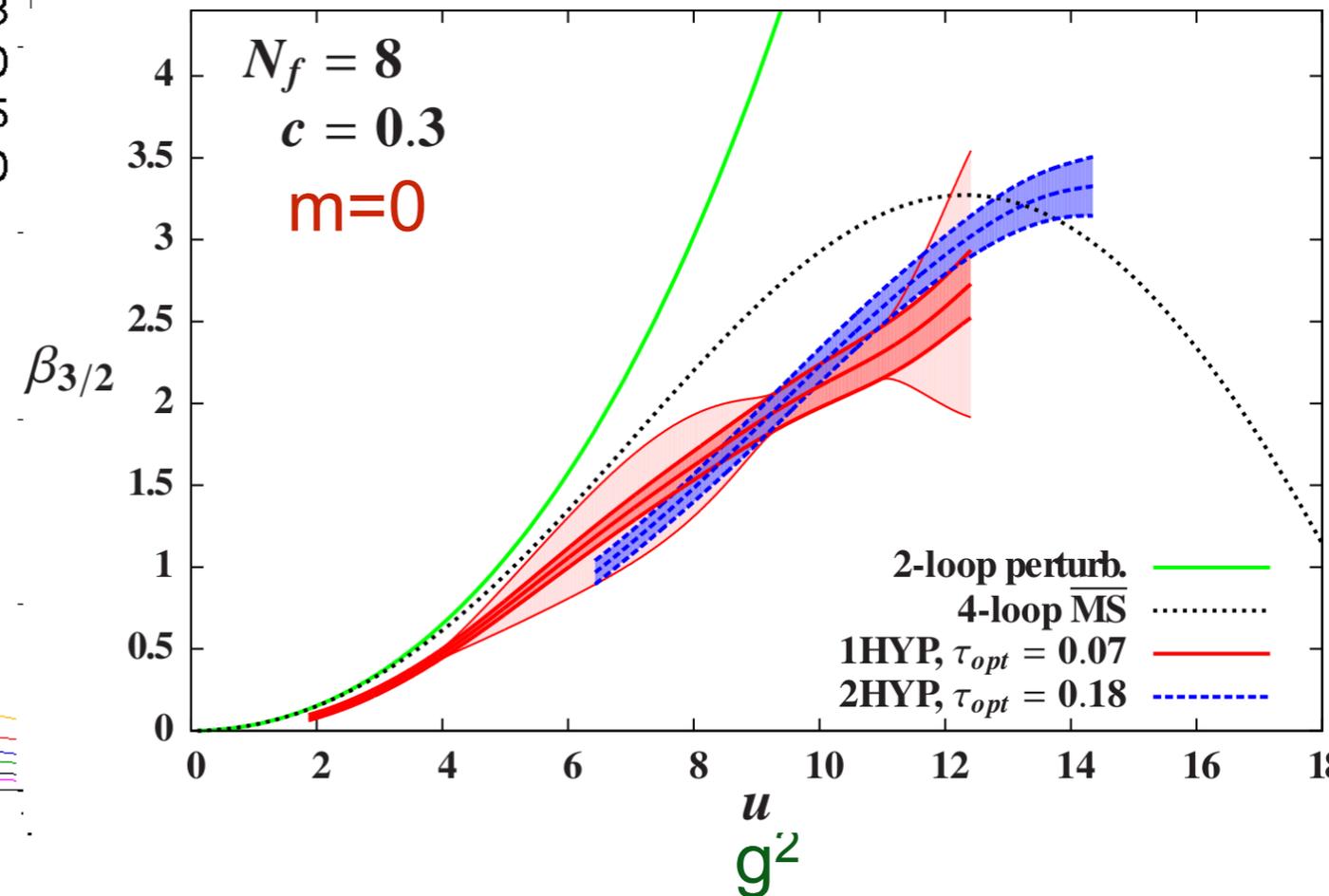
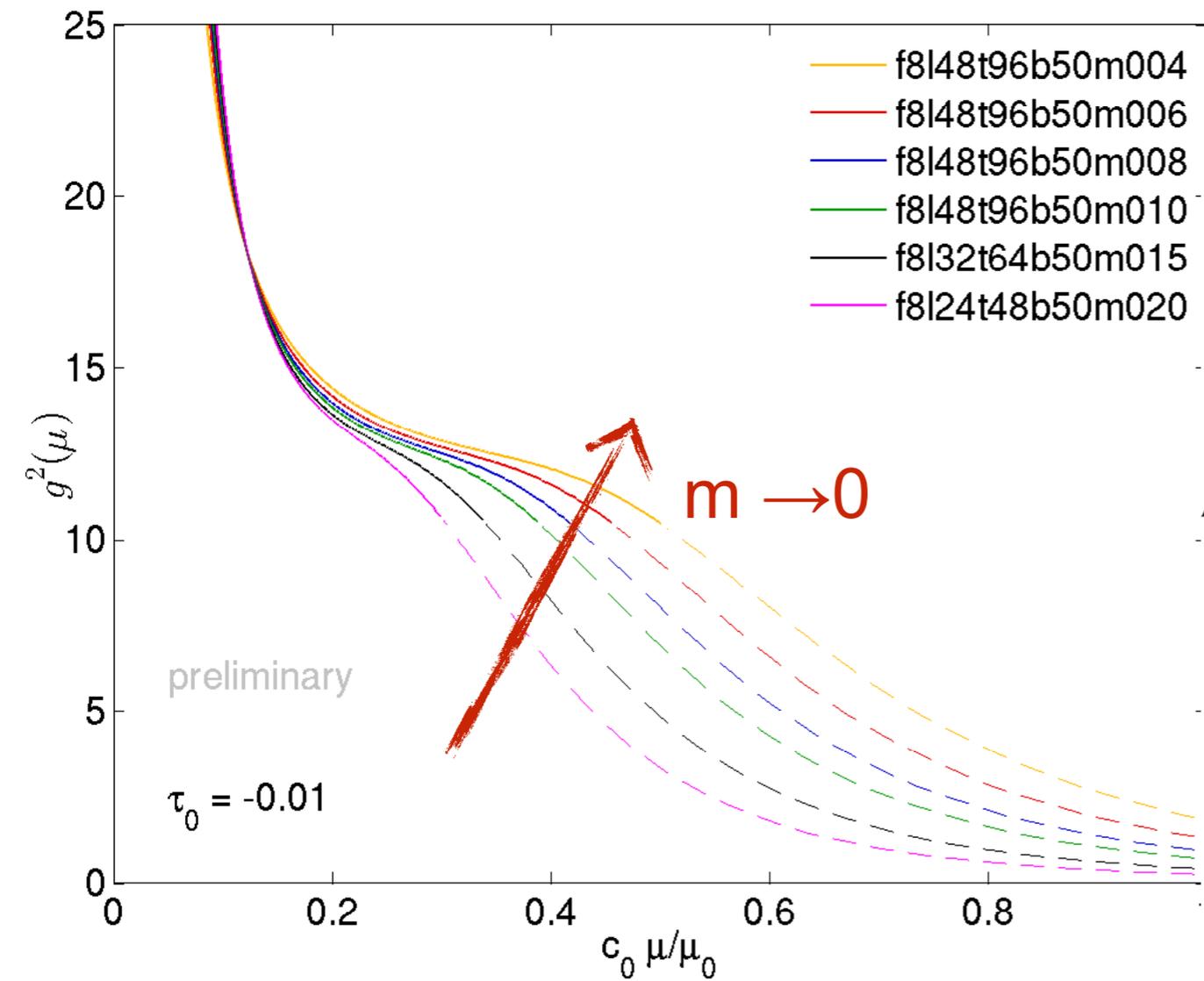
Is this walking?

The “shoulder” is the gauge dynamics : slow evolution



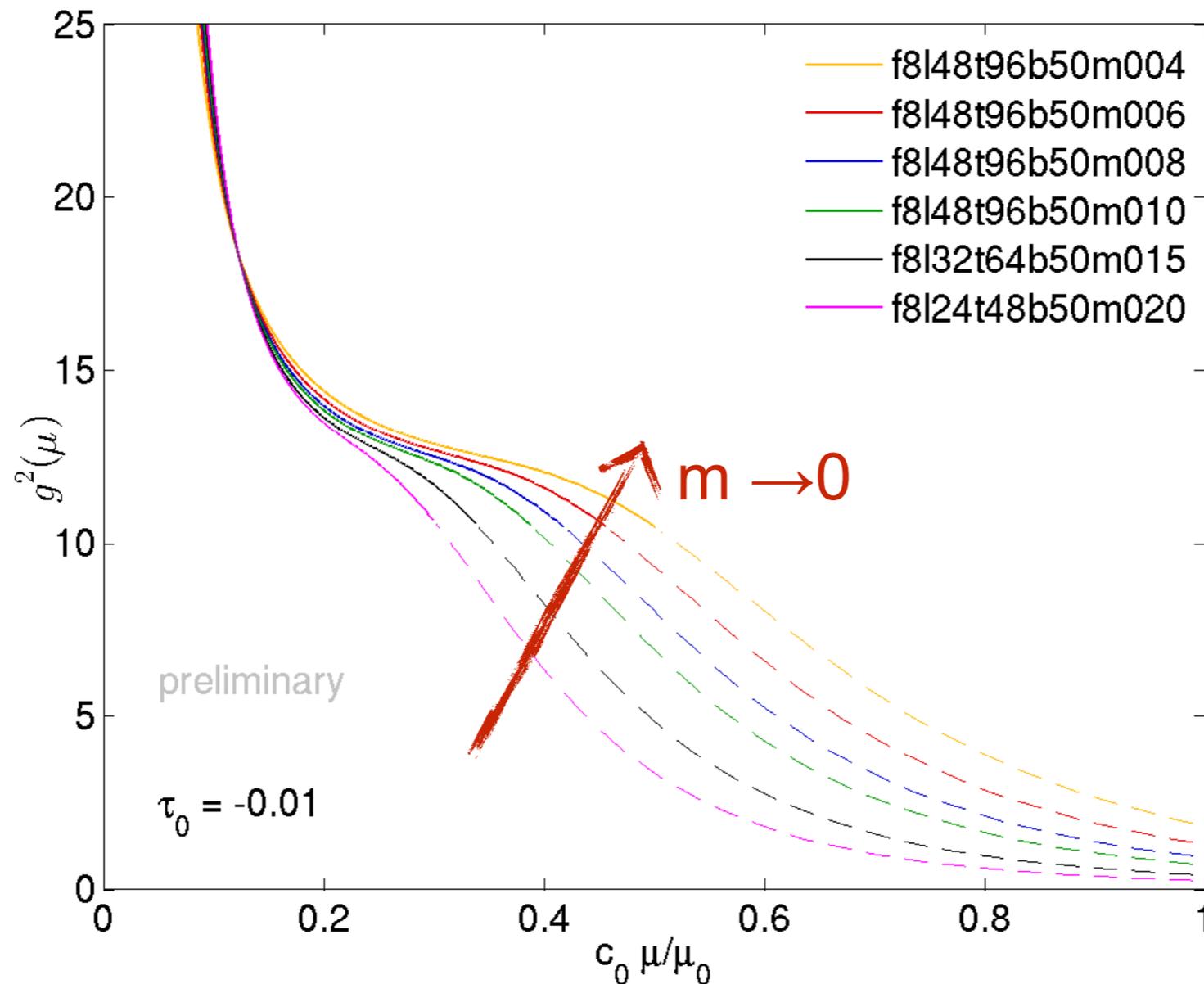
# Improved running coupling : 8 flavors

arXiv:1410.5886



The "shoulder" is the gauge dynamics : slow evolution

# Improved running coupling : 8 flavors

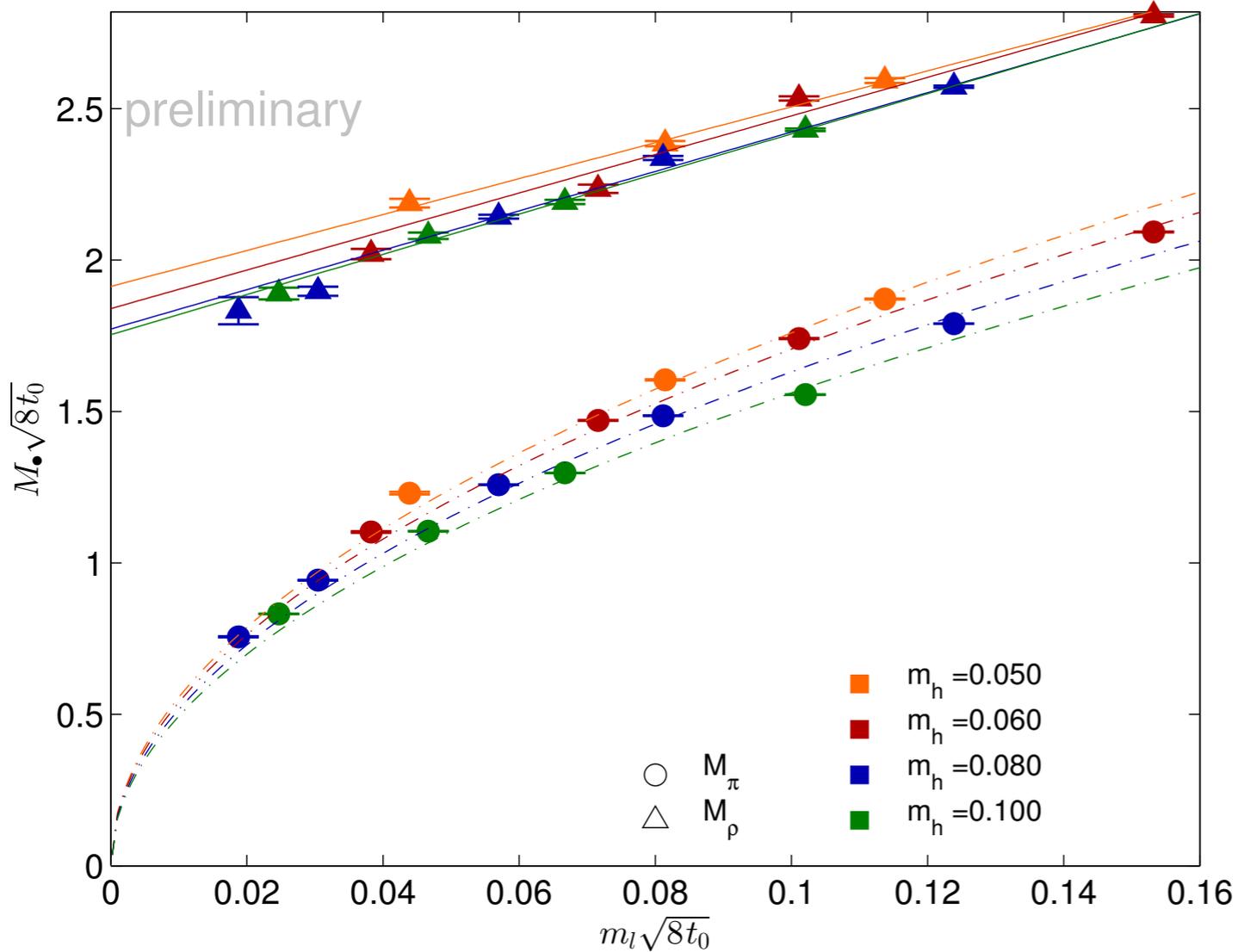


The “shoulder” is the gauge dynamics : slow evolution

The fast rise is due to the fermion mass running

What is the consequence of the two separate regimes?

# Connected spectrum, 4+8 flavors

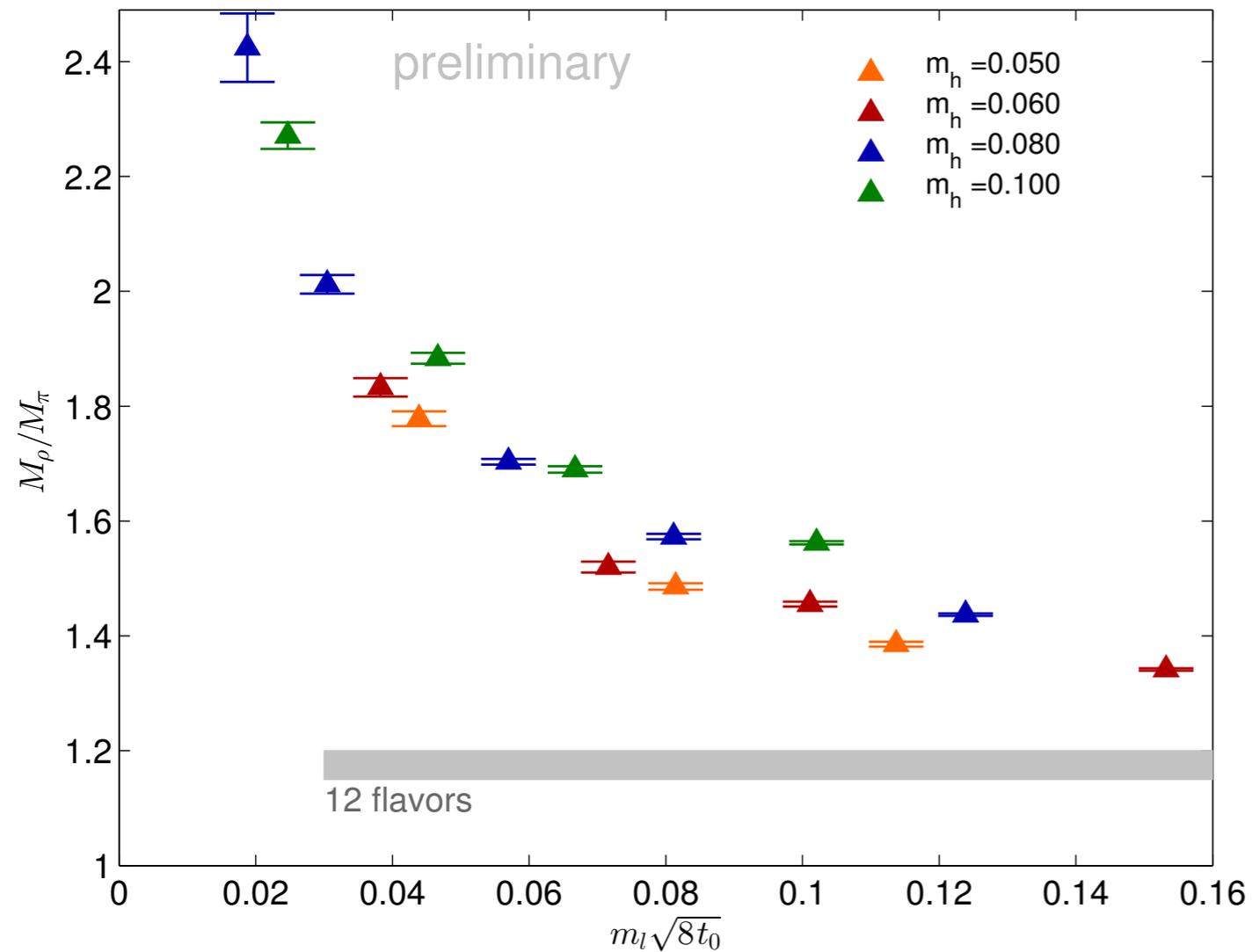


>  $M_\pi, M_\rho$  vs  $m_\ell$   
(rescaled by the gradient flow  
scale  $\sqrt{8t_0}$  )

– little variation with  $m_h$

# Chiral limit ?

$M_\rho/M_\pi$  shows that we approach the chiral regime

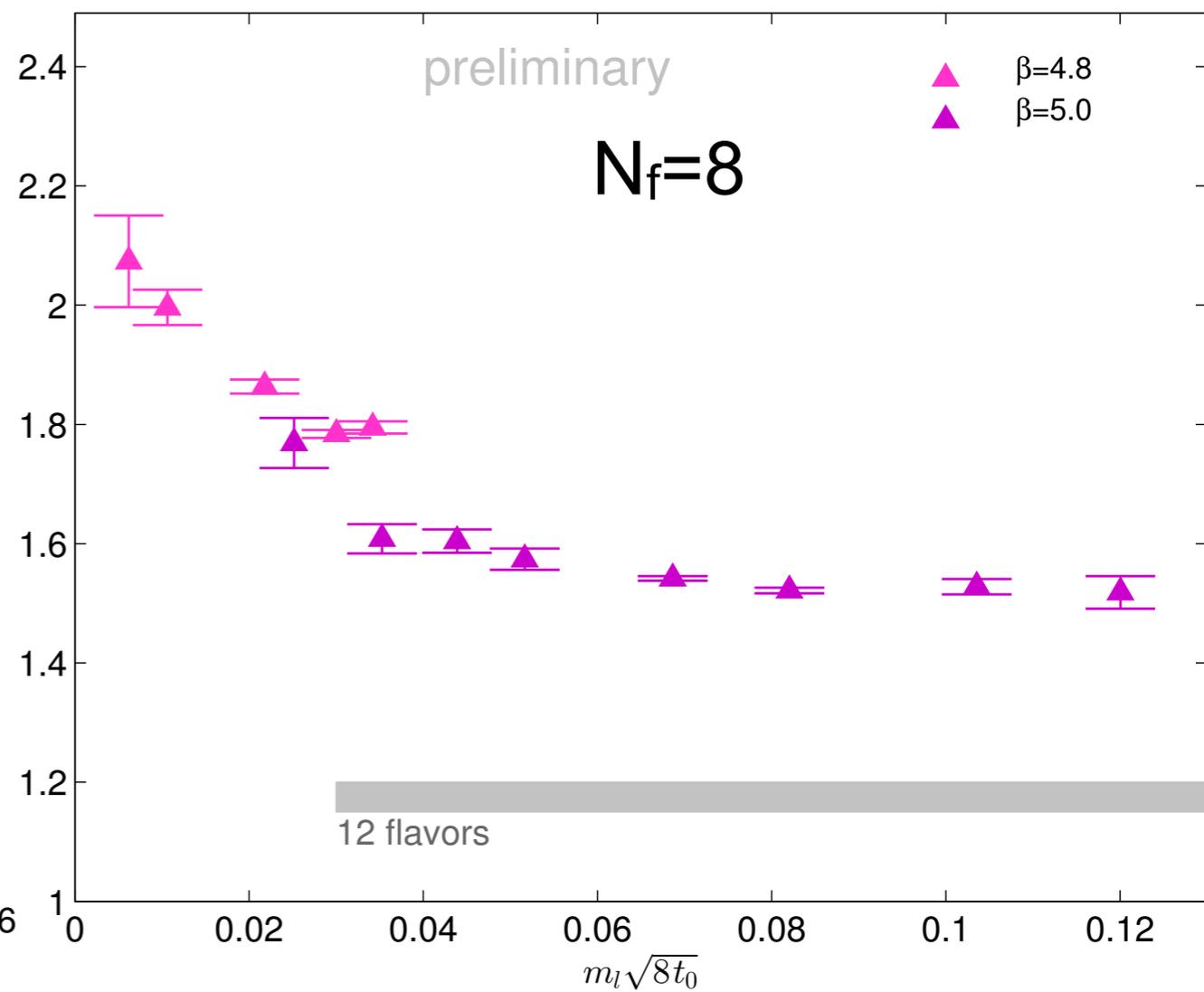
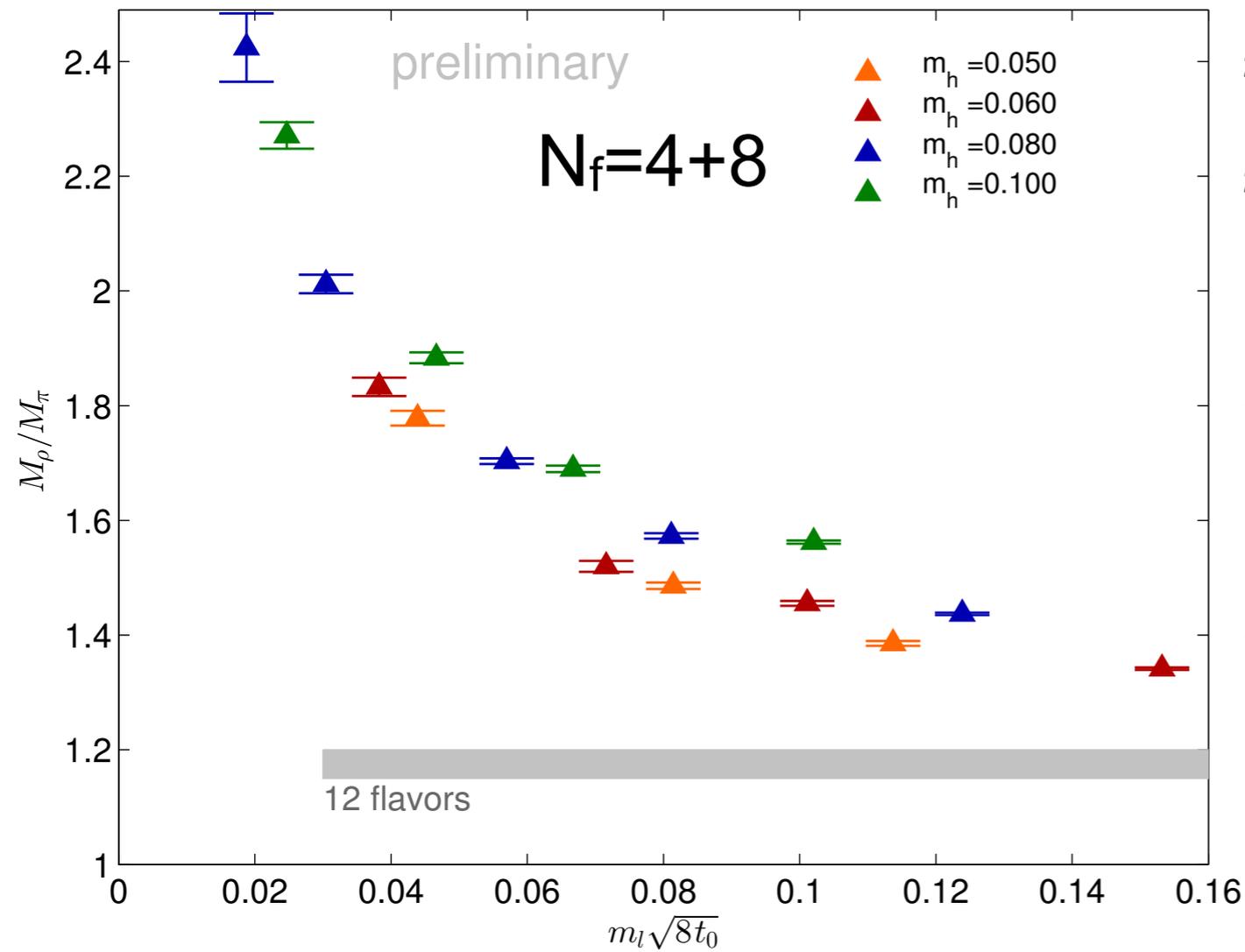


$< N_f=12$  predicts an almost constant ratio (as should be in a conformal system)

(arXiv:1401.0195)

# Chiral limit ?

$M_\rho/M_\pi$  : compare to 8 flavors



## Finally : the $0^{++}$ scalar state

We use the same method to construct and fit the correlators as with  $N_f = 8$  joint LSD project:

- Disconnected correlators:
  - 6 U(1) sources
  - diluted on each timeslice, color, even/odd spatial
  - variance reduced  $\langle \bar{\psi}\psi \rangle$
- Fit:
  - correlated fits to both parity (staggered) states
  - the **vacuum subtraction** introduces very large uncertainties
    - it is advantageous to add a (free) constant to the fit

$$C(t) = c_{0^{++}} \cosh\left(M_{0^{++}}\left(N_T/2 - t\right)\right) + c_{\pi_{\overline{sc}}} (-1)^t \cosh\left(M_{\pi_{\overline{sc}}}\left(N_T/2 - t\right)\right) + v$$

–this is equivalent to fitting the finite difference of the correlator

$$C(t+1) - C(t)$$

# Mixing in the $0^{++}$ channel

There is one major difference between  $N_f = 4 + 8$  and  $8$  :

- with non-degenerate masses the  $0^{++}$  splits to light and heavy states
- there is mixing the heavy and light species

This is similar to  $\eta - \eta'$  mixing in QCD

→ need to diagonalize the correlator matrix

$$C(t) = \begin{pmatrix} D_{ll}(t) - C_{ll}(t) & \sqrt{2}D_{lh}(t) \\ \sqrt{2}D_{hl}(t) & 2D_{hh}(t) - C_{hh}(t) \end{pmatrix}$$

Normalization: even though we describe 4 and 8 flavors, on the lattice they correspond to 1 and 2 staggered species

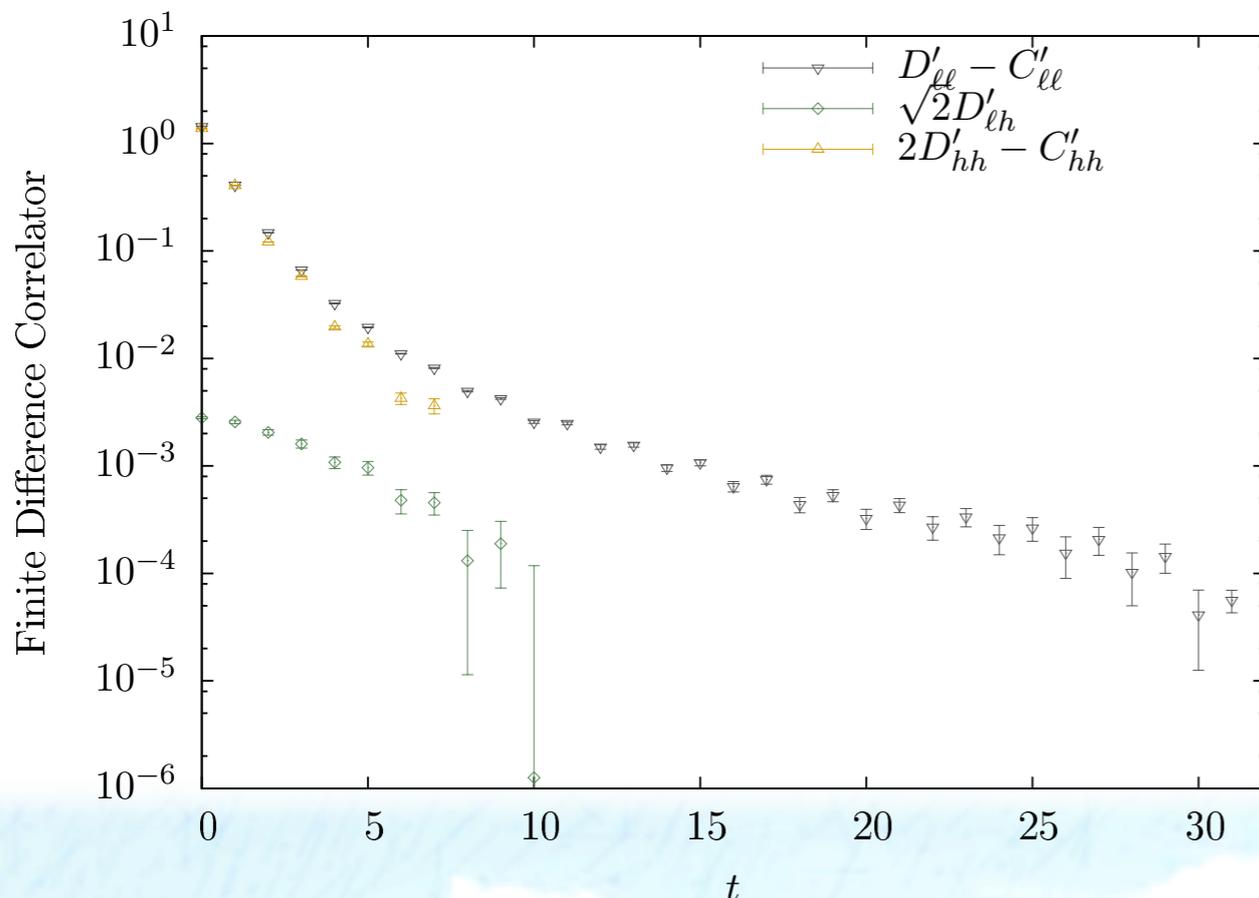


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Diagonalizing  $C(t)$  could lead to very large statistical errors.

**Fortunately:**  $D_{\ell h} \ll$  diagonal terms for almost all parameter values



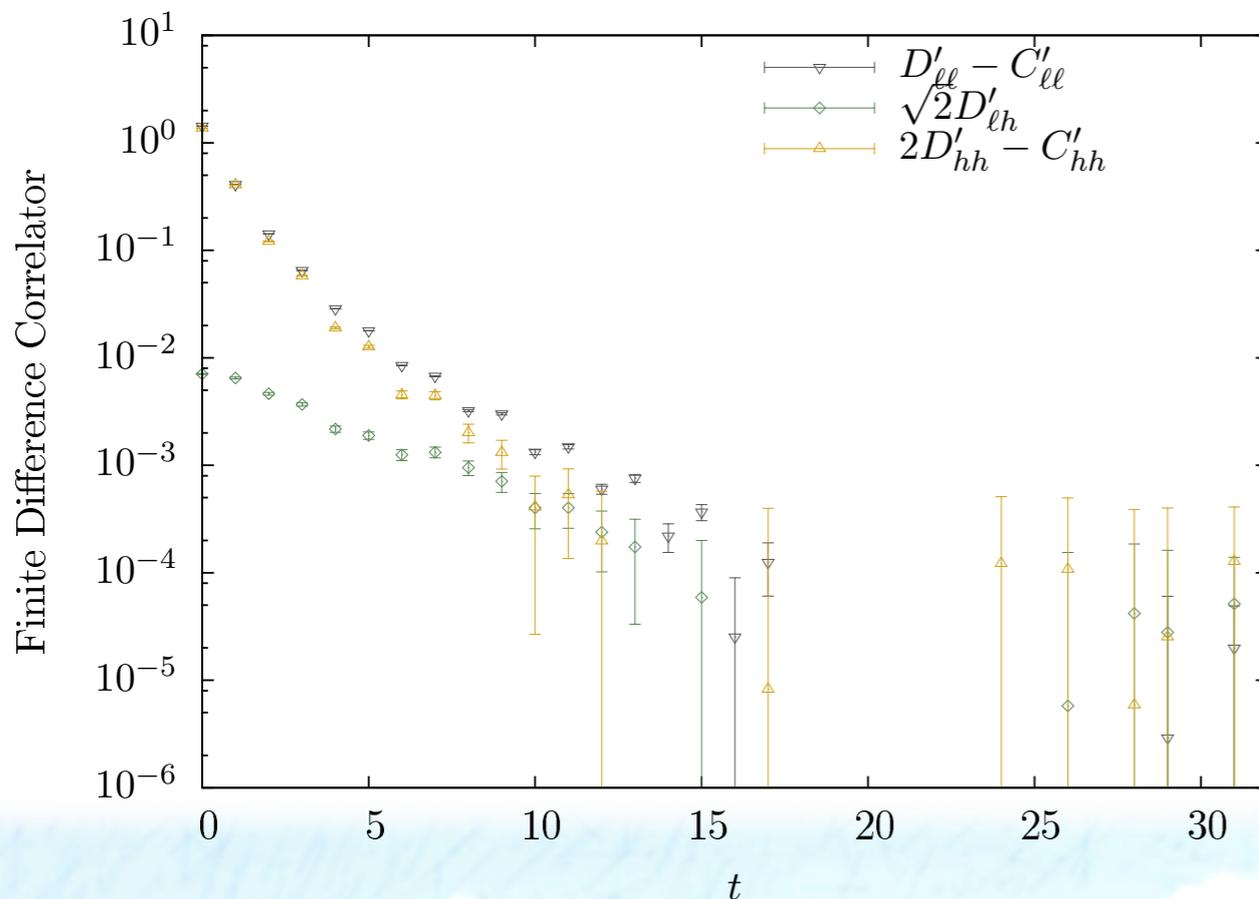
Finite difference correlators at  
 $m_h = 0.05$ ,  $m_\ell = 0.005$

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Diagonalizing  $C(t)$  could lead to very large statistical errors.

**Fortunately:**  $D_{\ell h} \ll$  diagonal terms for almost all parameter values  
**but not always!**



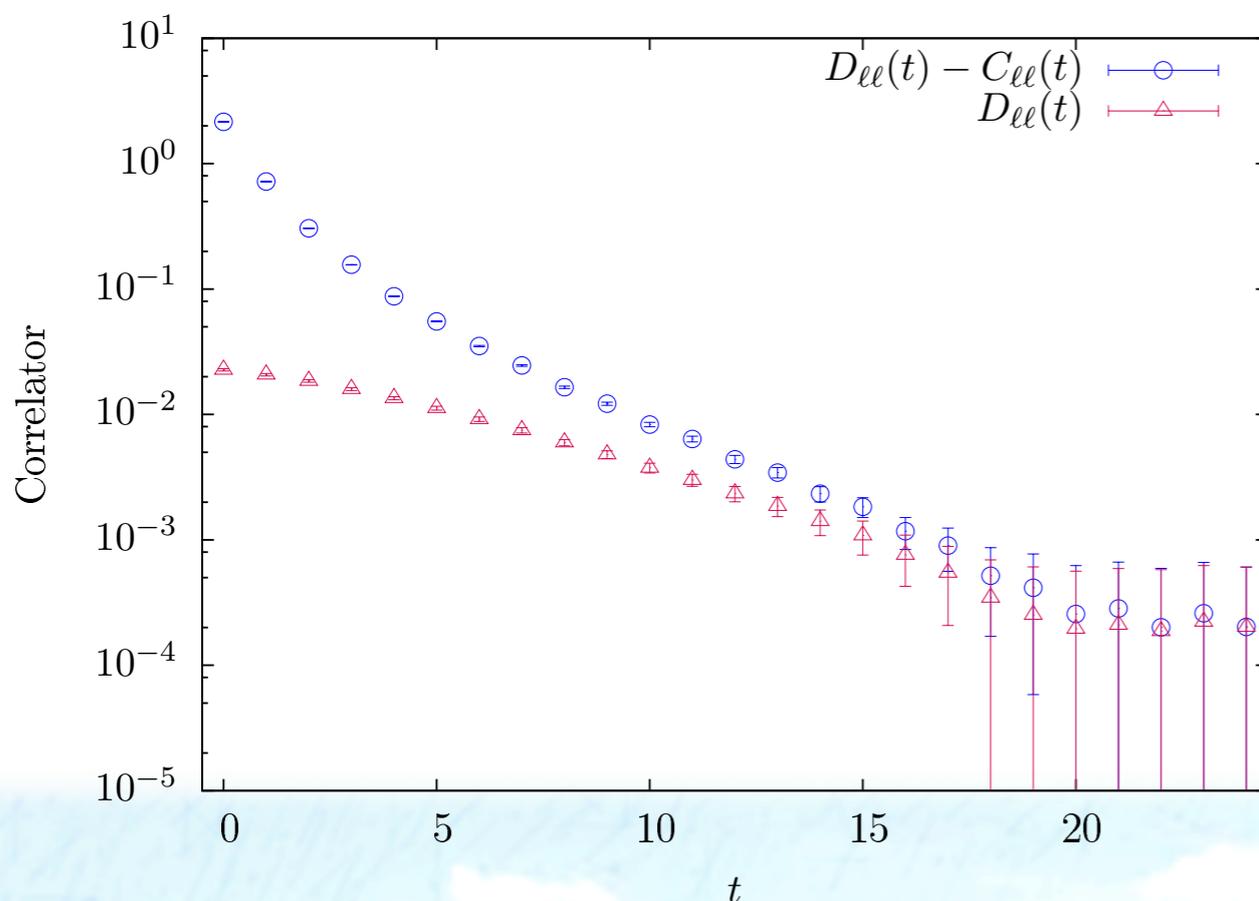
Derivative correlators at  
 $m_h = 0.05$ ,  $m_\ell = 0.015$

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Diagonalizing  $C(t)$  could lead to very large statistical errors.

**Fortunately:** the lightest excitation in  $D_{\ell\ell}$  (and  $D_{\ell h}$ ,  $D_{hh}$ ) is the  $0^{++}$



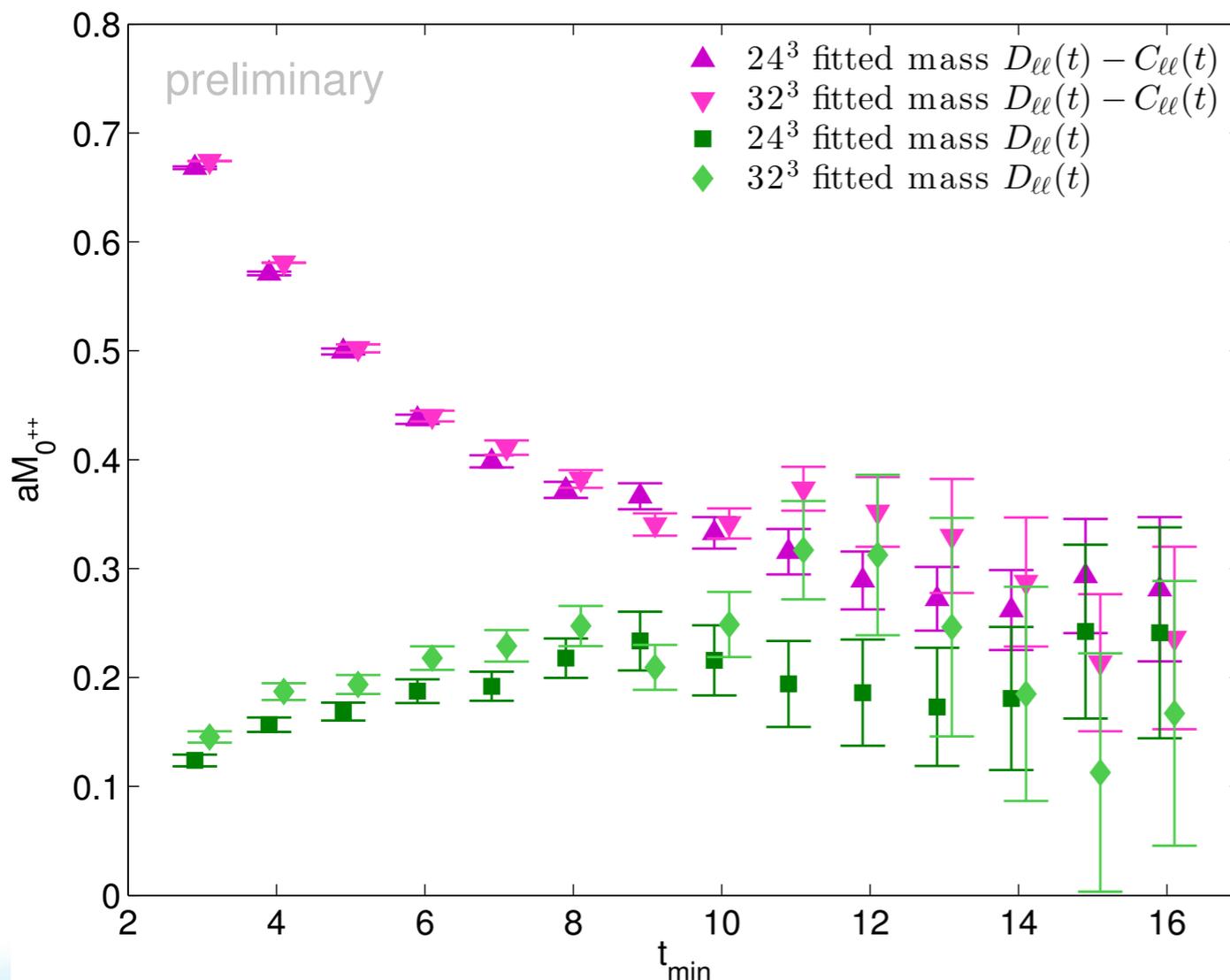
Derivative correlators at  
 $m_h = 0.06$ ,  $m_\ell = 0.010$ :

$D_{\ell\ell}$  and  $D_{\ell\ell} - C_{\ell\ell}$

# The $0^{++}$ mass

We strive to compare predictions from  $D_{\ell\ell}$  and  $D_{\ell\ell} - C_{\ell\ell}$  correlators  
– in the  $t \rightarrow \infty$  limit they should agree

$m_h = 0.06$  ,  $m_\ell = 0.010$ :

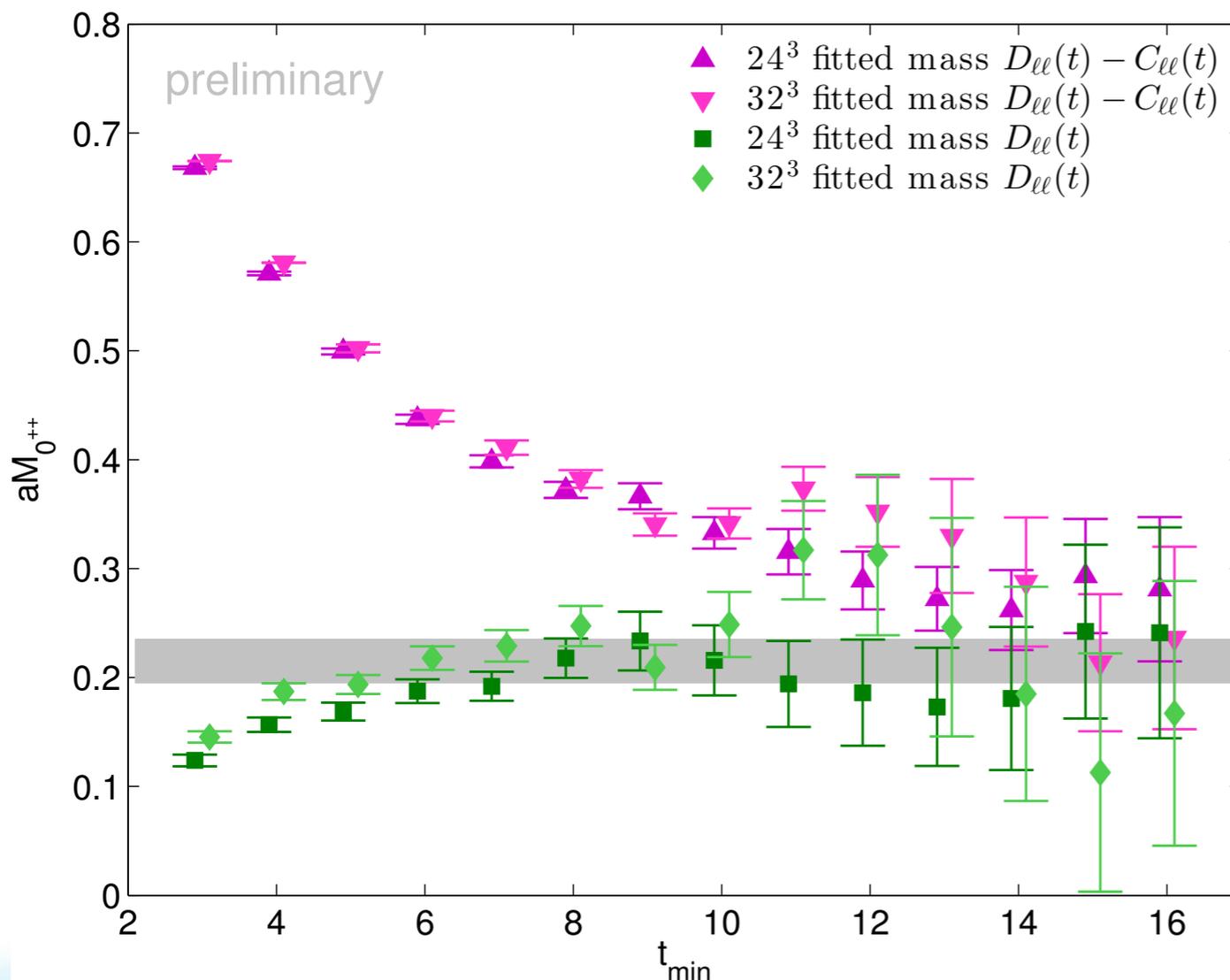


$M_{0^{++}}$  predicted from non-linear  
range fits ( $t_{\min} - N_T/2$ )

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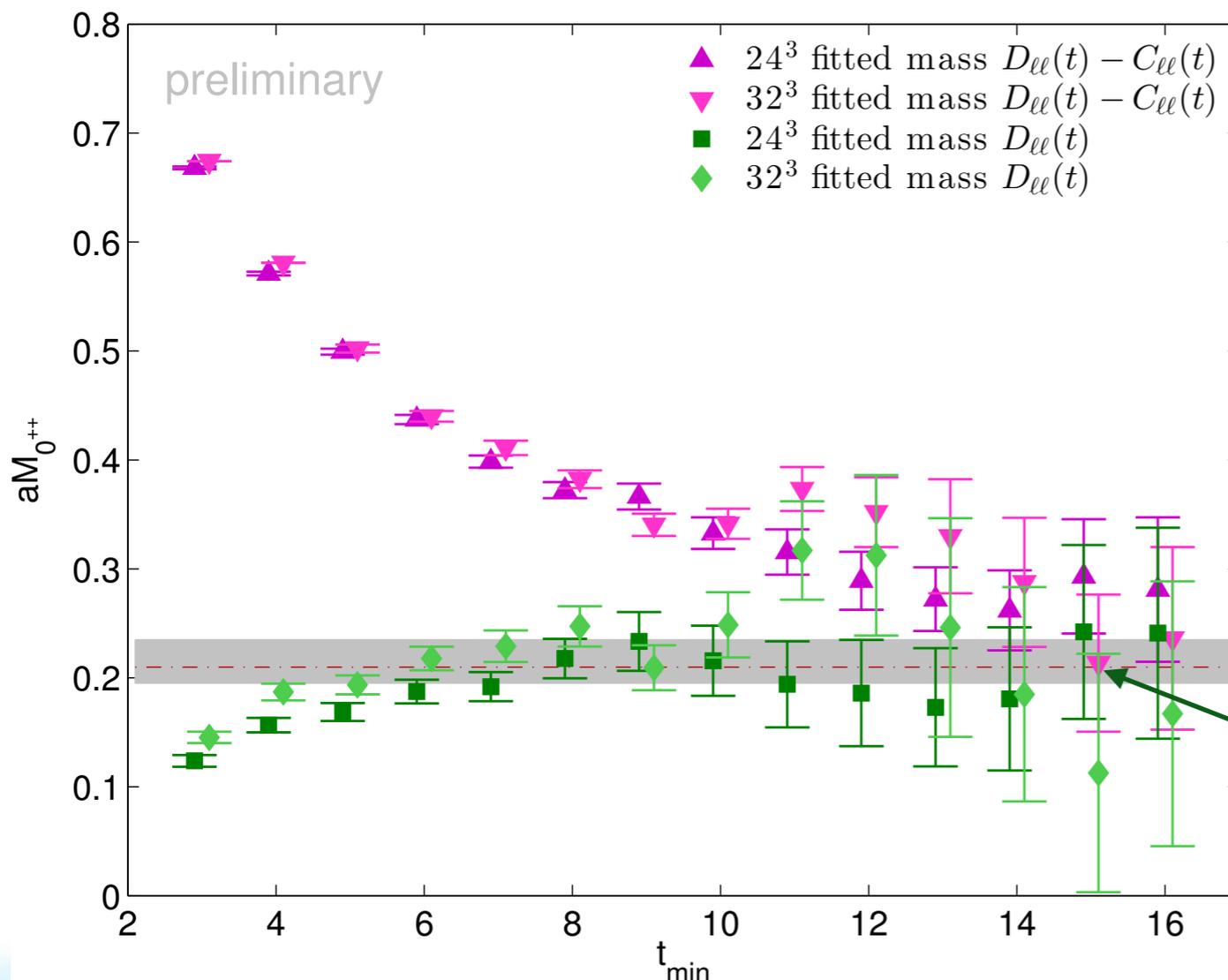
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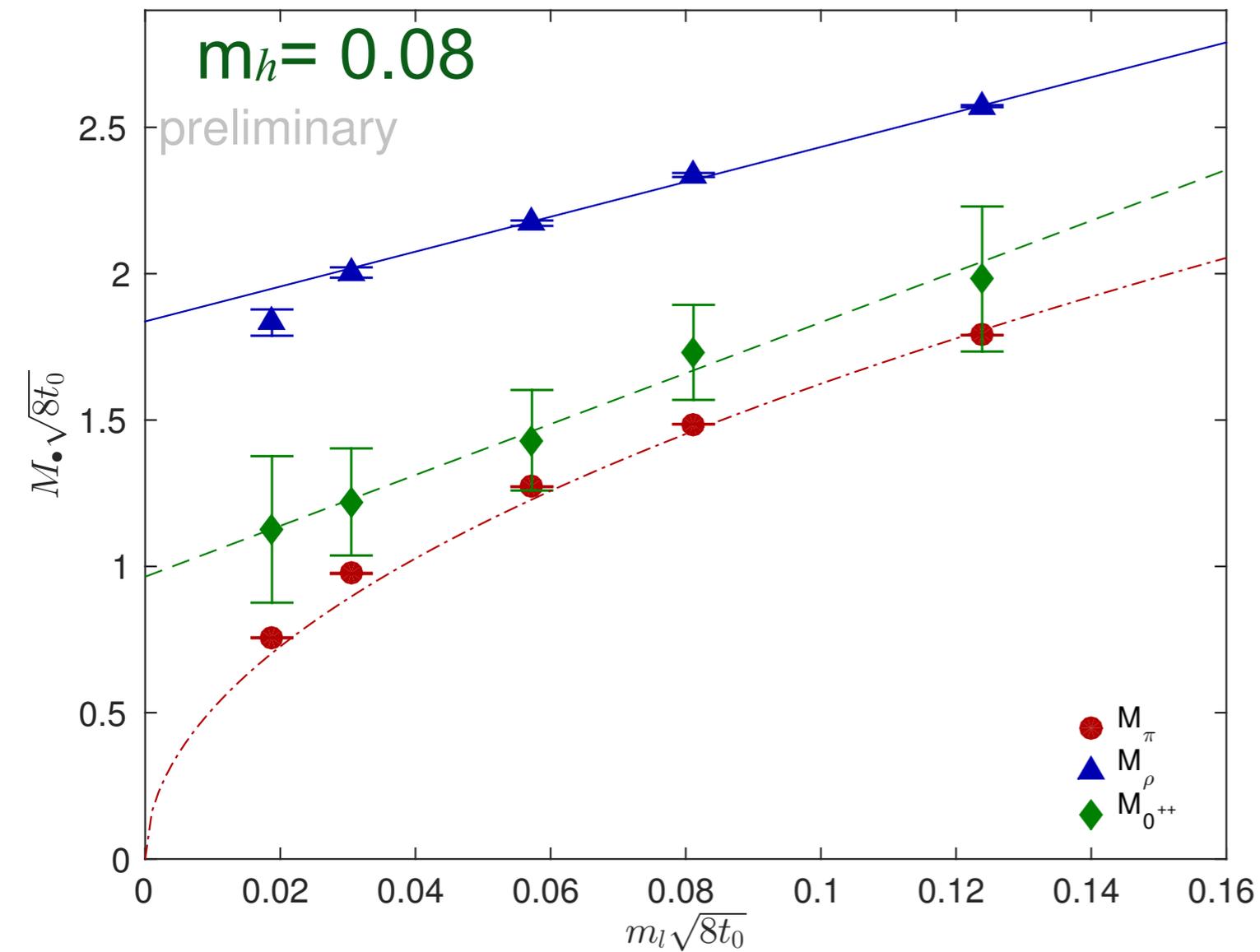
$M_{0^{++}}$  predicted from non-linear range fits ( $t_{\min} - N_T/2$ )

both volumes, both correlators predict a consistent value

pion

# Spectrum

Compare the pion, rho and  $0^{++}$  masses:

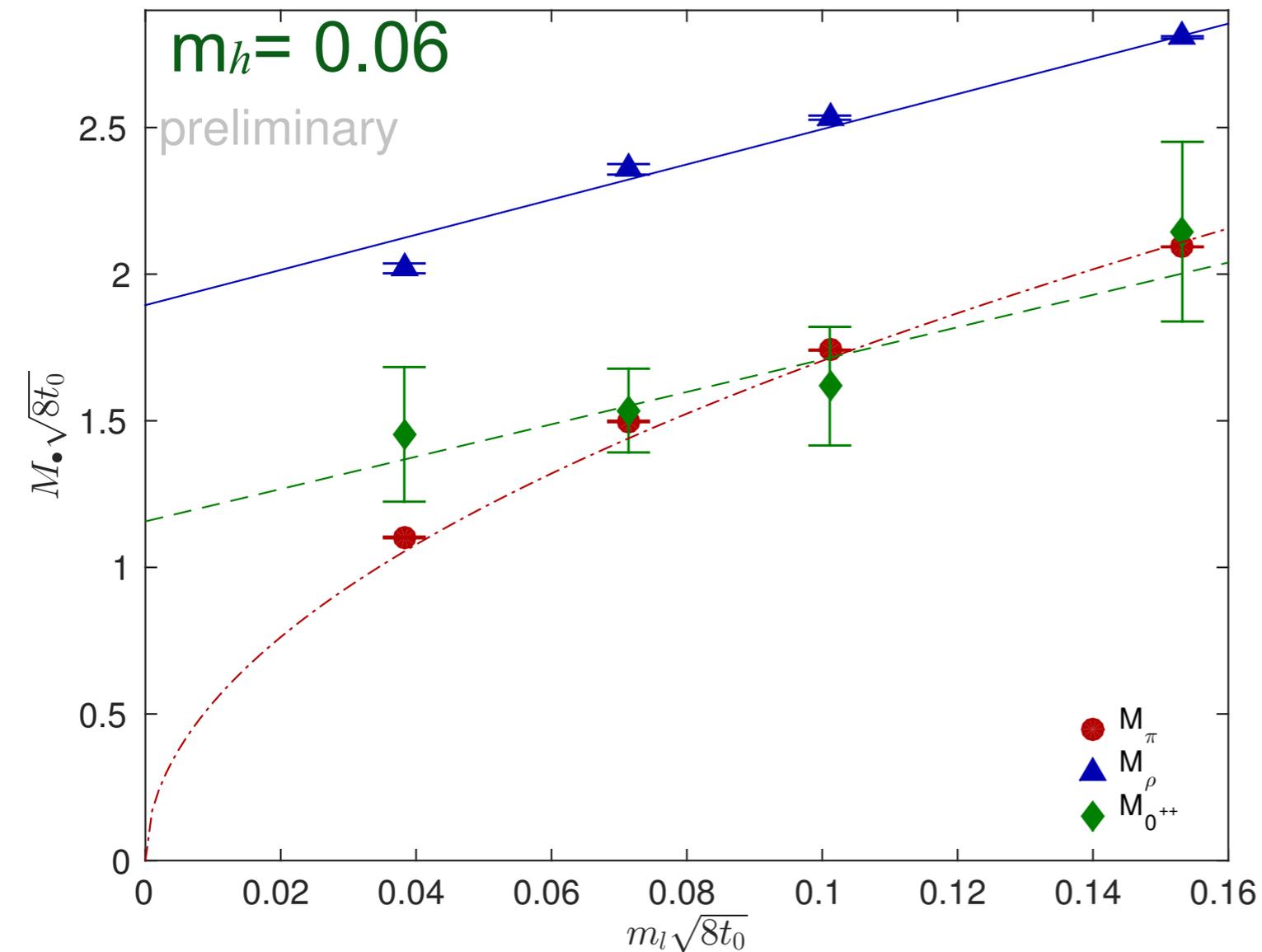


$m_h = 0.08$ : the  $0^{++}$

- is just above the pion,
- not Goldstone
- well below the rho

# Spectrum

Compare the pion, rho and  $0^{++}$  masses:



$m_h = 0.06$ : the  $0^{++}$

- is degenerate with pion at heavier  $m_\ell$
- need larger volumes, more statistics to resolve the small  $m_\ell$  region

# Conclusion & Summary

Lots of interesting possibilities ....

Lattice studies are needed to investigate strongly coupled systems

Even those without apparent phenomenological importance can teach us :

- understand universality
  - Wilson vs staggered vs rooted staggered vs domain wall fermions
- understand general properties of strongly coupled systems
  - walking near the conformal window
  - $0^{++}$  near the conformal window

Models with split fermion masses, like the 4+8 flavor model, help us navigate the landscape



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