

The Gradient Flow and the Running Coupling

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- The Gradient Flow
 - New tool for lattice calculations
 - Renormalized quantities
 - Scale setting, chiral condensate
- The Coupling Scheme
 - Several possible schemes
 - Correction terms
 - SF results

The Gradient Flow

Gradient Flow

- Flow along the gradient of the an action ¹
e.g. the gauge action:

$$\partial_t B_{\mu,t} = D_{\mu,t} B_{\mu\nu,t},$$

$$B_{\mu,0} = A_\mu$$

$$B_{\mu\nu,t} = \partial_\mu B_{\nu,t} - \partial_\nu B_{\mu,t} + [B_{\mu,t}, B_{\nu,t}]$$

- Diffusive flow with scale $\sqrt{8t}$
- Produces a finite, smoothed field
- Correlators of the flow field renormalized

¹M. Luscher and P. Weisz, JHEP **1102** (2011) 051 [arXiv:1101.0963].

Gradient Flow

- On the lattice, Wilson flow

$$G_{\mu,t+\epsilon} = e^{-\epsilon Z_{\mu,t}(x)} G_{\mu,t}(x),$$

$$Z_{\mu,t}(x) = (1/2) \left(M - M^\dagger \right) - (1/6) \text{Tr} \left(M - M^\dagger \right)$$

$$M = \sum_{\nu \neq \mu} \left(\begin{array}{c} \square \leftarrow \\ \square \uparrow \end{array} + \begin{array}{c} \square \rightarrow \\ \square \downarrow \end{array} \right)$$

- Stout smearing steps, continuous flow
- Cooling, flow generated by the gauge force

Gradient Flow

- For example, the field strength:²

$$\begin{aligned}\langle E_t \rangle &= \frac{1}{4} \langle G_{\mu\nu,t} G_{\mu\nu,t} \rangle \\ &= \frac{3(N_c^2 - 1)g_{MS}^2}{128\pi^2 t^2} \left[1 + k_1 g_{MS}^2 + \mathcal{O}(g_{MS}^4) \right]\end{aligned}$$

$$\begin{aligned}k_1 &= 1/(4\pi) [N(11/3\gamma_E + 52/3 - 3\ln 3) \\ &\quad - N_f(2/3\gamma_E + 4/9 - 4/3\ln 2)]\end{aligned}$$

$$\mu = (8t)^{-1/2}$$

- Renormalization scale set by t

²M. Lüscher, JHEP **1008** (2010) 071 [arXiv:1006.4518].

Scale Setting

- One of the first applications, scale setting
- Find t_0 such that

$$t_0^2 \langle E_{t_0} \rangle = 0.3$$

- A physical scale given by $\ell = \sqrt{8t_0}$
- Another scale w_0 ³

$$W(t) = t \frac{d}{dt} [t^2 \langle E(t) \rangle]$$
$$w_0 = \sqrt{t}|_{W(t)=0.3}$$

³S. Borsanyi *et al.* JHEP **1209** (2012) 010 [arXiv:1203.4469].

Extension to Fermions

- The flow equations can be extended to fermionic fields: ⁴

$$\begin{aligned}\partial_t \chi_t &= \Delta_t \chi_t, & \partial_t \bar{\chi}_t &= \bar{\chi}_t \overset{\leftarrow}{\Delta}_t, \\ \chi_0 &= \psi & \bar{\chi}_0 &= \bar{\psi} \\ \Delta_t &= D_{\mu,t} D_{\mu,t}, & D_{\mu,t} &= \partial_\mu + B_{\mu,t}\end{aligned}$$

- Not automatically renormalized, but

$$\chi_{R,t} = Z_x^{1/2} \chi_t \quad \bar{\chi}_{R,t} = \bar{\chi}_t Z_x^{1/2}$$

- Independent of t

⁴M. Lüscher, JHEP **1304** (2013) 123 [arXiv:1302.5246]

Extension to Fermions

- Non-zero flow time chiral condensate:

$$\Sigma_t^{rr} = -\langle \bar{\chi}_{r,t} \chi_{r,t} \rangle, \quad \Sigma_{R,t}^{rr} = Z_\chi \Sigma_t^{rr}$$

- To the first order

$$\Sigma_{R,t}^{rr} = \frac{2Nm_{R,r}}{(4\pi)^2 t} \int_0^\infty dv \frac{e^{-vz}}{(1+v)^2}, \quad z = 2tm_{R,r}^2$$

- Can also be related to the chiral condensate

Other Applications

- Anisotropy tuning ⁵
- Energy-momentum tensor ⁶
- Matching improved actions with new observables:

$$E(t) = \frac{1}{4} G_{\mu\nu}(t) G_{\mu\nu}(t),$$

$$S_t^{rs} = -\bar{\chi}_{r,t} \chi_{s,t}, \quad P_t^{rs} = -\bar{\chi}_{r,t} \gamma_5 \chi_{s,t}$$

⁵S. Borsanyi *et al.* arXiv:1205.0781.

⁶H. Suzuki, PTEP **2013** (2013) 8, 083B03 [arXiv:1304.0533].

L. Del Debbio *et al.* JHEP **1311** (2013) 212 [arXiv:1306.1173].

Coupling from the Gradient flow

Coupling from the Gradient flow

- $E(t)$ is related to the coupling:⁷

$$g_{GF}^2 = \frac{t^2 \langle E(t) \rangle}{N}$$

- To study scale dependence, need a single scale
- Now we have L and t , so fix

$$t = \frac{(\textcolor{blue}{c} L)^2}{8}$$

⁷Z. Fodor *et al.* JHEP **1211** (2012) 007 [arXiv:1208.1051].

Z. Fodor *et al.* PoS LATTICE **2012** (2012) 050 [arXiv:1211.3247].

Boundary Conditions

- Periodic
- Schrödinger Functional boundary conditions⁸
 - Unique global minimum: simpler perturbation theory
 - Breaks the translation symmetry
- Twisted boundary conditions ⁹
 - Milder than the SF boundary
 - Restricts the number of fermions in the fundamental representations to $N_f = 2nN_c$

⁸P. Fritzsch and A. Ramos, JHEP **1310** (2013) 008 [arXiv:1301.4388].

P. Fritzsch and A. Ramos, arXiv:1308.4559.

J. Rantaharju, arXiv:1311.3719.

⁹A. Ramos, arXiv:1308.4558.

Boundary Conditions

- Schrödinger Functional boundary conditions:
- Periodic spatial boundary conditions, trivial time boundaries

$$U_k(x) = 1, \text{ when } x_0 = 0, L$$

$$U_\mu(x + L\hat{k}) = U_\mu(x)$$

- Fermion fields disappear at time boundaries, periodic spatial boundary conditions

$$\psi(x) = 0, \text{ when } x_0 = 0, L$$

$$\psi(x + L\hat{k}) = \psi(x).$$

The Model

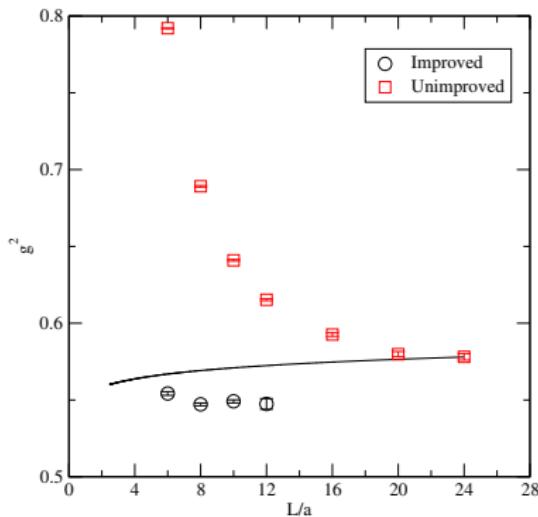
- The Model:

$$S = (1 - c_g)S_G(U) + c_g S_G(V) + S_F(V) + c_{SW} \delta S_{SW}(V)$$

- Smeared Wilson fermion action $S_F(V)$
- and gauge action $S_G(V)$
- HEX Smearing
- Bulk correction term $\delta S_{SW}(V)$
- Here, $c_g = 0.5$ and $c_{SW} = 1$
- $N(a/L)$ in $g^2 = t^2 \langle E \rangle / N(a/L)$ measured at $\beta = 80$

Lattice Results

- Large discretization errors in previous results, corrections:



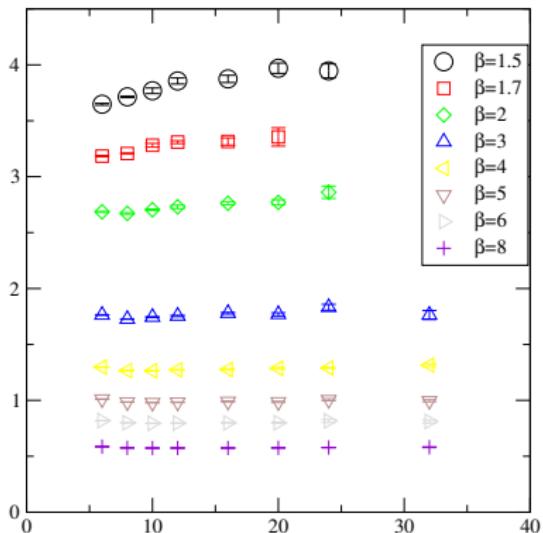
- Gradient flow and E from the Symanzik action

$$\square \rightarrow \frac{5}{3} \square - \frac{1}{10} \square$$

- Correction to the flow time

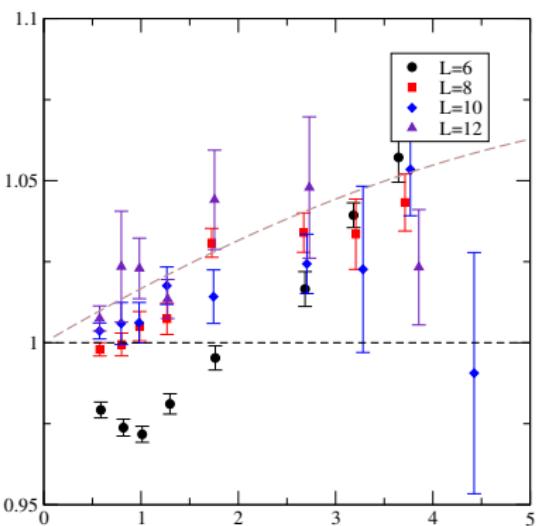
$$t = c^2 L^2 / 8 - d a^2,$$
$$c = 0.6, d = 0.5$$

SU(2) with SF boundaries



- Preliminary
- SU(2) with 8 flavors
- Relatively stable results
- Allows for large lattice sizes
- Measurements still at low coupling

SU(2) with SF boundaries



- Preliminary
- Scale dependence described by the step scaling function

$$\Sigma(u, a/L) = \frac{g_{GF}^2(g_0, 2L/a)|_{g_{GF}^2(g_0, L/a)=u}}{u}$$

$$\sigma(u) = \lim_{a \rightarrow 0} \Sigma(u, a/L)$$

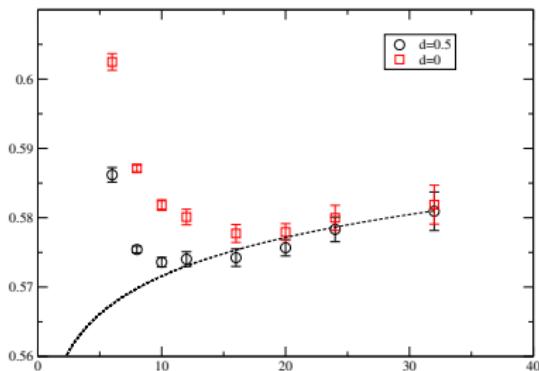
- The Gradient Flow

- Renormalized quantities
- Scale setting
- Chiral condensate
- Space-time symmetries
- Improved actions

- The Coupling Scheme

- Several possible schemes
- Correction terms
- Promising SF results

The effect of d



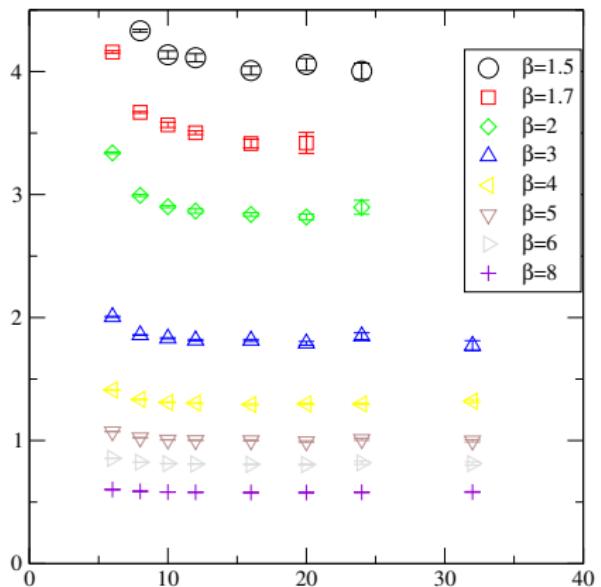
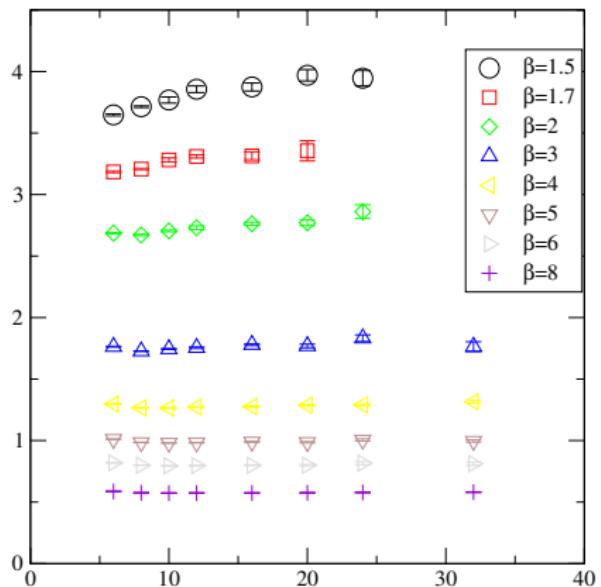
- The effect of d in the flow time
- SU(2) with 2 adjoint fermions,
 $\beta = 8$

$$t = c^2 L^2 / 8 - da^2,$$

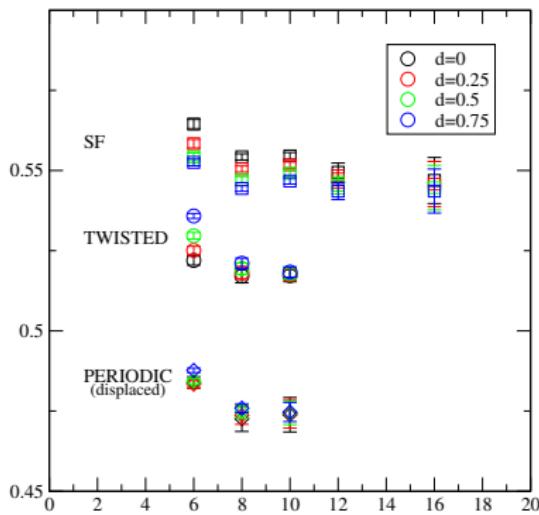
twisted: $d = 0$

SF,periodic: $d = 0.5$

Compared to $d = 0$



The effect of d



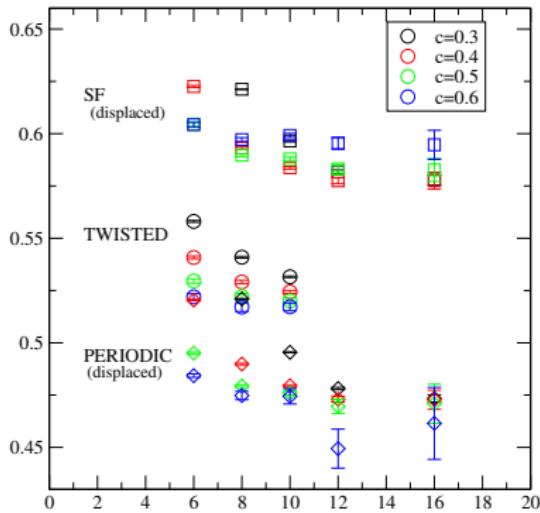
- The effect of d in the flow time
- SU(2) with 2 adjoint fermions,
 $\beta = 8$

$$t = c^2 L^2 / 8 - da^2,$$

twisted: $d = 0$

SF,periodic: $d = 0.5$

The effect of c



- The effect of c in the flow time
- Labels different schemes in the continuum

$$t = c^2 L^2 / 8 - da^2,$$

twisted: $d = 0$

SF,periodic: $d = 0.5$