Twisted reduction in large N QCD with adjoint Wilson fermions

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We study the large N QCD with adjoint fermions using the twisted space-time reduced model.

For two flavor theory ($N_f=2$) with N=289, string tension is calculated which seems to vanish at $m_q=0$, in a way consistent with the theory governed by an infrared fixed point with $\gamma_*=0.8\sim1.2$.

For one flavor theory ($N_f = 1$), string tension remains finite at $m_q = 0$, indicating a confining theory .

Plan of the talk

- Twisted Eguchi-Kawai model for pure SU(N) gauge theory
- large N QCD with two adjoint fermions
- large N QCD with one adjoint fermion

• Eguchi-Kawai model

Eguchi-Kawai model is obtained from the usual SU(N) lattice gauge theory

$$Z_W = \int \prod_{x,\mu} dU_{x,\mu} \exp\left\{bN\sum_{x} \sum_{\mu\neq\nu=1}^d Tr\left(U_{x,\mu}U_{x+\mu,\nu}U_{x+\nu,\mu}^{\dagger}U_{x,\nu}^{\dagger}\right)\right\}$$

by neglecting the space-time dependence of the link variables

$$U_{x,\mu} \to U_{\mu} \longrightarrow$$

$$Z_{EK} = \int \prod_{\mu} dU_{\mu} \exp\left\{bN \sum_{\mu \neq \nu=1}^{d} Tr\left(U_{\mu}U_{\nu}U_{\mu}^{\dagger}U_{\nu}^{\dagger}\right)\right\}, \qquad b = \frac{1}{g^{2}N}$$

In the same way, Wilson loop is defined by

$$W_W(C) = \left\langle Tr \left(U_{x,\mu} U_{x+\mu,\nu} \cdots U_{x-\rho,\rho} \right) \right\rangle$$
$$W_{EK}(C) = \left\langle Tr \left(U_{\mu} U_{\nu} \cdots U_{\rho} \right) \right\rangle$$

Eguchi and Kawai show that in the large N limit, the Schwinger-Dyson eqs. satisfied by Wilson loops are identical in both theories provided that the $(Z(N))^d$ symmetry

$$U_{\mu} \rightarrow e^{i\theta_{\mu}}U_{\mu}, \quad e^{i\theta_{\mu}} \in Z(N)$$

of the EK model is not spontaneously broken.

Bhanot, Heller and Neuberger found, however, that this symmetry is broken spontaneously in the weak coupling region.

• Twisted Eguchi-Kawai model



EK model can be viewed as the usual Wilson gauge theory having only one site with the periodic boundary conditions.

We introduce twisted boundary conditions in SU(N), $N = L^2$ theory

$$S_{TEK} = bN \sum_{\mu \neq \nu=1}^{d} Tr \left(Z_{\mu\nu} U_{\mu} U_{\nu} U_{\mu}^{\dagger} U_{\nu}^{\dagger} \right)$$
$$Z_{\mu\nu} = \exp\left(k \frac{2\pi i}{L}\right), \quad Z_{\nu\mu} = -Z_{\mu\nu}, \qquad \mu > \nu$$

k, *L*: co-prime, *k*/*L* fixed as we go $N = L^2 \rightarrow \infty$

Our model is related to ordinary SU(N) lattice theory on $V = L^4$ space-time volume up to $O(1/N^2)$ corrections. The number of degree of freedom of SU(N) matrix is $N^2 = L^4$ If the TEK model is correct nonperturbatively, we should be able to calculate the string tension.

In our reduced model, the Wilson loop W(R,T) is defined by

$$W(R,T) = Z_{\mu\nu}^{RT} \left\langle Tr \left(U_{\mu}^{R} U_{\nu}^{T} U_{\mu}^{\dagger R} U_{\nu}^{\dagger T} \right) \right\rangle$$

Then the string tension σ is obtained from Creutz ratio as

$$\chi(R',T') = -\log \frac{W(R'+0.5,T'+0.5)W(R'-0.5,T'-0.5)}{W(R'+0.5,T'-0.5)W(R'-0.5,T'+0.5)}$$

$$\chi(R',R') = \sigma + \frac{2\gamma}{R'^2} + \frac{\eta}{R'^4}$$

with half-integer R', T'.



We calculate the continuum string tension by extrapolating the TEK data with $N = 841 = 29^2$, k = 9 at 6 values of b

b = 0.36, 0.365, 0.37, 0.375, 0.38, 0.385

Our system should be related to the lattice theory with $V = 29^4$

For comparison, we also calculate the continuum string tension using ordinary SU(N) lattice gauge theory with N = 3, 4, 5, 6, 8 on a $V = 32^4$ lattice

Comparison of the continuum string tension $\Lambda_{\overline{MS}}$ / $\sqrt{\sigma}$

TEK model with $N = 841 = 29^2$ and LGT with N = 3, 4, 5, 6, 8



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Motivation for $N_f = 2$ adjoint fermions

SU(N) LGT with two adjoint fermions is thought to be conformal or nearly conformal for any value of N since the first two coefficient of beta functions expressed in term of 't Hooft coupling is independent of N.

$$b_{0} = \frac{4N_{f} - 11}{24\pi^{2}}, \quad b_{1} = \frac{16N_{f} - 17}{192\pi^{4}}$$

$$b_{0} < 0 \rightarrow N_{f} < \frac{11}{4} = 2.75 \quad \text{asymptotic free}$$

$$b_{1} > 0 \rightarrow N_{f} > \frac{17}{16} = 1.08 \quad \text{infrared fixed point}$$

Twisted reduced model of large N QCD with two adjoint Wilson fermions

We consider gauge group SU(N), $N = L^2$

$$\begin{split} S &= bN \sum_{\mu \neq \nu=1}^{d} Tr \left(Z_{\mu\nu} U_{\mu} U_{\nu} U_{\mu}^{\dagger} U_{\nu}^{\dagger} \right) + \sum_{j=1}^{N_{f}} \overline{\psi}_{j} D_{W} \psi_{j} \\ Z_{\mu\nu} &= \exp \left(k \frac{2\pi i}{L} \right), \quad Z_{\nu\mu} = Z_{\mu\nu}^{*}, \qquad \mu > \nu \\ D_{W} &= 1 - \kappa \sum_{\mu=1}^{4} \left[(1 - \gamma_{\mu}) U_{\mu}^{adj} + (1 + \gamma_{\mu}) U_{\mu}^{\dagger adj} \right], \quad U_{\mu}^{adj} \psi_{j} = U_{\mu} \psi_{j} U_{\mu}^{\dagger} \end{split}$$

k, L: co-prime

k = 0 corresponds to periodic boundary condition $k \neq 0$ corresponds to twisted boundary condition



К

We calculate the string tension with $N = 289 = 17^2$, k = 5at 2 values of b = 0.35, 0.36 for various values of κ

Our system should be related to the lattice theory with $V = 17^4$

•For $N_f = 2$, we use the Hybrid Monte Carlo method.

Simulations have been done on Hitachi SR16000 at KEK One node: 32 cores power 7, peak speed 980 GFlops 256 GB shared memory

Sustained speed of our code in one node is 600 Gflops at N=289

We thank to Hitachi system engineers !



$$m_q = 1 / \kappa - 1 / \kappa_c$$

So far, we have not calculated any hadronic spectrum.

However, it is straightforward to calculate the lowest eigenvalue of positive hermitian Wilson Dirac operator $Q^2 = (D_W \gamma_5)^2$, which should be related to the physical quark mass square.

$$Q^{2} = (D_{W}\gamma_{5})^{2}, \quad D_{W} = 1 - \kappa \sum_{\mu=1}^{4} \left[(1 - \gamma_{\mu})U_{\mu}^{adj} + (1 + \gamma_{\mu})U_{\mu}^{\dagger adj} \right]$$





We can fit both the string tension and the lowest eigenvalue of Q^2 with the following fitting form

$$a(1/\kappa-1/\kappa_c)^b$$

Requiring that both the string tension and the lowest eigenvalue of Q^2 vanish at the same critical κ_c , we have $\kappa_c = 0.1694(7)$ and

String tension

$$b = 1.10(4) = \frac{2}{1 + \gamma_*}$$
 $\therefore \gamma_* = 0.81(8)$

The lowest eigenvalue of Q^2

$$b = 1.28(1) \stackrel{?}{=} f(\gamma_*)$$

We repeat the same analysis at b=0.36

$$a(1/\kappa-1/\kappa_c)^b$$

String tension

$$b = 0.92(9) = \frac{2}{1 + \gamma_*}$$
 $\therefore \gamma_* = 1.17(21)$

The lowest eigenvalue of Q^2

$$b = 1.33(4)$$

with common $\kappa_c = 0.168(1)$



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• We can simulate $N_f \neq 2$ theory by using the Rational Hybrid Monte Carlo method

$$\left(Q^2\right)^{-s} = \sum_i \frac{\beta_i}{Q^2 - \alpha_i}$$

 $N_f = 1$

In the large N limit, equivalent to $N_f = 2$ fundamental fermion in rank two anti-symmetric rep. For N=3, it is just two flavor QCD

 $N_f = 1/2$

supersymmetric gauge theory (one majorana fermion)?

 $N_f = 3/2$

conformal or confining ?

 $N_f = 5 / 2$

conformal or confining ?

Motivation for $N_f = 1$ adjoint fermion

In the large N limit, $N_f = 1$ adjoint fermion is equivalent to $N_f = 2$ fundamental fermion in rank two anti-symmetric rep. (orientifold equivalence, Kovtun, Unsal, Yaffe)

For N=3, the latter theory is just two flavor QCD and our model corresponds to Corrigan-Ramond large-N limit.

We then expect the reduced model of $N_f = 1$ adjoint fermion as

- confining theory
- there might exist either Aoki phase or Sharpe-Singleton phase transition since we use Wilson fermion.





Conclusion

We have demonstrated that the twisted reduced model of large N QCD with adjoint Wilson fermions works quite well.

 $N_{f} = 2$

String tension is calculated at N=289, which clearly decreases as we increase kappa and seems to vanish around kappa ~ 0.17 in a way consistent with the theory governed by an infrared fixed point with $\gamma_* = 0.8 \sim 1.2$.

$$N_{f} = 1$$

String tension is calculated at N=289, which clearly remains finite around kappa ~ 0.16 strongly suggesting this is the confining theory. We also find that cg iteration does not converge for kappa > 0.16. Aoki phase or Sharpe-Singleton phase transition? Remaining important problems

- We need to understand the finite N (finite volume) effects make simulation with larger $N = 529 = 23^2$
- calculate hadronic correlators

On going project

Wilson fermions with $N_f = 1/2$, 1, 2

overlap fermions with $N_f = 2$ with Garcia-Perez, Gonzalez-Arroyo, Ishikawa, Keegan