# Reaching the chiral limit in many flavor systems

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In collaboration with A. Cheng, G. Petropoulos and D. Schaich ArXiv:1111:2317,1207.7162,1207.7164

## 4<sup>th</sup> of July Independence Day Fireworks



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Discovery of a Higgs-like state at 125GeV







Discovery of a Higgs-like state at 125GeV

This is not what we expected, but we have to deal with it.

Is there room for a composite (strongly coupled) Higgs?



## Composite Higgs in strongly coupled systems:

Still an attractive idea:

 $SU(N_{color} \ge 2)$  gauge fields +  $N_{flavor}$  fermions in some representation



 $\mathsf{N}_{\mathsf{color}}$ 



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## Which model? What representation, N<sub>c</sub>, N<sub>f</sub>? What property? What method?

In Colorado we developed several methods to study conformal and near-conformal systems:

Phase diagram at zero and finite temperature

ArXiv:1111:2317,1207.7162

- Dirac eigenmodes & the mass anomalous dimension ArXiv:1207.7164
- Monte Carlo renormalization group matching ArXiv:1212.xxxx

We tested with N=4, 8 and 12 fundamental fermions with SU(3) gauge Found some surprising results





Bulk transition: lattice artifact but a real phase transition IRFP: its location is scheme dependent, not physically observable



## Finite temperature and bulk phase transitions



## In a conformal system

- finite temperature transitions run into a bulk (T=0) transition
- β<sub>bulk</sub> separates strong coupling (confining) and weak coupling (conformal) phases

Phase diagram in  $\beta$ -m space for  $(N_f=12)$ 

Intermediate phase bordered by **bulk** 1<sup>st</sup> order transitions The chiral bulk transition fissioned into two (This has been observed by Deuzeman et al, LHC collab. as well)



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A new symmetry breaking pattern

Single-site shift symmetry (S<sup>4</sup>):  $x_{\mu} \rightarrow x_{\mu} + \mu$ 

is exact symmetry of the action but broken in the IM phase

→ plaquette expectation value is "striped"



A new symmetry breaking pattern

Order parameters: Plaquette difference: Link difference:

## $\Delta P_{\mu} = \langle \operatorname{Re} \operatorname{Tr} \Box_{n} - \operatorname{Re} \operatorname{Tr} \Box_{n+\mu} \rangle_{n_{\mu} \operatorname{even}}$ $\Delta L_{\mu} = \langle \alpha_{\mu}(n) \overline{\chi}(n) U_{\mu}(n) \chi(n+\mu)$ $- \alpha_{\mu}(n+\mu) \overline{\chi}(n+\mu) U_{\mu}(n+\mu) \chi(n+2\mu) \rangle_{n_{\mu} \operatorname{ev}}$

n+2µ

n+µ

n



 $\beta$  = 2.6 IM phase  $\beta$ =2.7 weak coupling phase

## S<sup>4</sup>b symmetry breaking pattern

- Single-site shift symmetry is exact in the action, S<sup>4</sup>b phase has to be bordered by a "real" phase transition
- Exist with 8 & 12 flavors, not with 4
- S<sup>4</sup>b phase
  - Could signal a special taste breaking
  - Confining (static potential, Polyakov loop)
  - Chirally symmetric (meson spectrum, Dirac eigenvalue spectrum)

Such phase does not exist in the continuum limit

Must be pure lattice artifact



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## in gauge-fermion systems

Must be pure lattice artifact **within gauge fermion systems** Could become physical with some other interaction



Phase diagram in  $\beta$ -m space for  $N_f$ =12

What is the relation between bulk and finite T transitions? Finite T =  $1/(N_ta)$  simulations with  $N_t$ =8,12,16,20





Phase diagram in  $\beta$ -m space for  $(N_f=12)$ 

Finite T transitions are stuck to the S<sup>4</sup> phase boundary No confining phase at weak coupling:

transition from S<sup>4</sup> b $\rightarrow$  chirally symmetric



Phase diagram in  $\beta$ -m space for  $N_f$ =8

 $N_f$ =8 is expected to be chirally broken –  $S^4b$  phase ... must be an irrelevant lattice artifact ?





N<sub>t</sub> = 8,12,16 looks OK at m≥0.01.

- Weak coupling side shows both confining and deconfined phases
- Consistent with 2-loop PT



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At m=0.005 no confining phase on  $N_t \le 16$ the  $N_t = 12-16$  looses scaling ??



At m=0.005 no confining phase on  $N_t \le 16$ Let's try  $N_t = 20$ : looks OK.



We can check this in the chiral limit with direct m=0 simulations!

 $\rightarrow$  lost the confining phase in the chiral limit even on N<sub>t</sub>=20



## Dirac eigenvalue spectrum

Eigenvalues at small  $\lambda$  are related to IR physics

In conformal systems the eigenvalue density  $\rho$  scales as  $\rho(\lambda) \propto \lambda^{\alpha}$ 

The mode number  $v(\lambda) = V \int_{-\lambda}^{\lambda} \rho(\omega) d\omega \propto V \lambda^{\alpha+1}$  is RG invariant (Giusti,Luscher)

 $\boldsymbol{\rightarrow} \alpha$  is related to the anomalous dimension

$$\frac{4}{1+\alpha} = y_m = 1 + \gamma_m$$

(Zwicky,DelDebbio;Patella)



## The energy dependence of $\gamma_m$

 $\gamma_m$  depends on the energy scale : this is manifest as  $\lambda$  dependence of the eigenmode scaling



**IR** – small  $\lambda$  region:

 $\gamma_m(\lambda \to 0) \to \gamma^*$ 

predicts the universal anomalous dimension at the IRFP

**UV** – large  $\lambda$  =O(1) region: Governed by the UVFP (asymptotically free perturbative FP)  $\gamma_m(\lambda) \rightarrow 0$ 

**In between:** Energy dependent γ<sub>m</sub>

## The energy dependence of $\gamma_m$ :Chirally broken systems

The picture is still valid in the UV and moderate energy range





## Volume dependence

The scaling form is valid in  $V \rightarrow \infty$  only!

- Increase the volume until volume dependence vanishes
- **OR** combine different volumes & use the finite volume as advantage



## Extracting $\gamma_m$

- Fit:  $\log(v(\lambda))=c+(\alpha+1)\log(\lambda)$
- Volume dependence:
  - Ignore small  $\lambda$  /volume transient
  - Look for overall "envelope"



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We know what to expect:



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We know what to expect:

broken chiral symmetry in IR, asymptotic freedom in UV



Most of these data were obtained on deconfined (small) volumes at m=0!

Every test we have done in /near the chiral limit suggests IR conformality but the system is still controversial



#### β=3.0, 4.0, 5.0, 6.0

- There is no sign of asymptotic freedom behavior for β<6.0,</li>
  γ<sub>m</sub> grows towards UV
- Not possible to rescale different β's

Looks as if there were an IRFP around  $\beta$ =5.0



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- Not possible to rescale different β's

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Extrapolate to  $\lambda = 0$ :  $\gamma_m(\lambda \to 0) \to \gamma^* \approx 0.30(3)$ 

## The mode number

A few lessons on  $\gamma_m$  and the mode number

- Volume dependence is important, especially deep in the weak coupling
- $\gamma_m$  depends on  $\lambda$ , a constant fit will not work
- $\gamma_m$  shows strong  $\beta$  dependence :  $\lambda \rightarrow 0$  extrapolation is tricky





The finite temperature structure shows strange behavior. Eigenmodes are also closer to 12 than 4 flavors:



No asymptotic free scaling No rescaleability of different couplings When  $\gamma_m \sim 2$  in the UV, the S<sup>4</sup>b phase develops

If N<sub>f</sub>=8 is not conformal, it must be slowly walking.



## Conclusion & summary

Even after the 4<sup>th</sup> of July fireworks, strongly coupled systems are worth investigating:

- Lattice regularized models can show unexpected phases : S<sup>4</sup>b phase
- Finite temperature studies are reliable to study the phase structure only in the chiral limit (or very small bare mass)
- Dirac eigenmodes predict the energy dependent anomalous dimension but careful control of finite volume and  $\lambda \rightarrow 0$  extrapolation is needed

#### SU(3) gauge + fundamental fermions:

- N<sub>f</sub>=12 system looks conformal
- N<sub>f</sub>=8 system is unexpected: if not conformal, it must be slowly walking



## EXTRA SLIDES



## The finite temperature phase structure of $N_f=12$

were among the first BSM studies :

–Finite T transition with  $N_f \ge 4$  flavors is expected to be first order

- First results were as expected (2008) (Deuzeman, Lombardo, Pallante)
- Second generation studies found 2 first order transitions



(both Deuzeman et al and LHC)



## The phase structure of $N_f=12$

2 jumps in the fermion condensate on T=0 lattices (at finite T as well)



These are bulk transitions, present at T=0 and independent of the volume.



## Dirac eigenvalue spectrum

Much less is known about chirally symmetric systems:

- $\rho(0) = 0$  suggests the scaling form  $\rho(\lambda) \propto \lambda^{\alpha}$  $\lambda_0$  is a "soft edge", in conformal systems  $\lambda_0 = 0$
- The exponent  $\alpha$  is related to the mass anomalous dimension

(Luscher&Giusti,Zwicky& DelDebbio)

The mode number

$$v(\lambda) = V \int_{-\lambda}^{\lambda} \rho(\omega) d\omega \propto V \lambda^{1+\alpha} = (L \lambda^{(1+\alpha)/4})^4$$

is RG invariant  $\rightarrow$ 

$$\frac{1+\alpha}{4} = y_m = 1 + \gamma_m$$



## Extracting $\gamma_m$

- Configurations: 20-50 independent,  $12^3x24 \rightarrow 32^3x64$  volumes
- mass: 0.0025 → 0
  no observable mass effect (but m=0.01 would be too large!)
- Calculate eigenmodes: ~1000 per configuration
  Different volumes cover different λ range



- Volume dependence: The scaling form is valid in V→∞ only!
   Increase the volume until volume dependence vanishes
  - Combine different volumes & use the finite volume as advantage