

Relativistic Hydrodynamics in High-Energy Heavy Ion Collisions



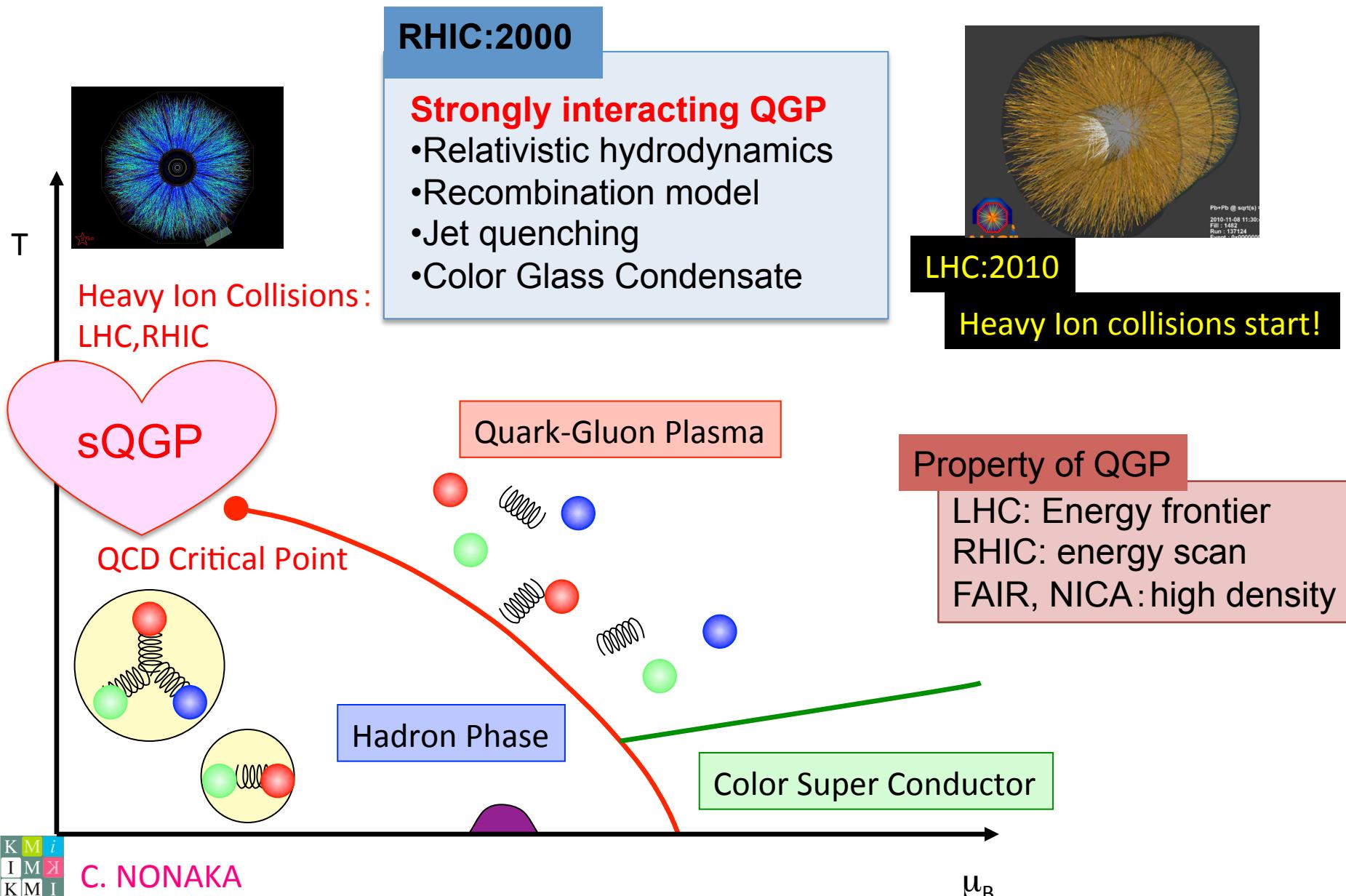
Kobayashi Maskawa Institute
Department of Physics, Nagoya University

Chiho NONAKA

Kobayashi-Maskawa Institute
for the Origin of Particles and the Universe

December 13, 2013 @ KMI 2013, Nagoya

Relativistic Heavy Ion Collisions



Dynamics of Heavy Ion Collisions

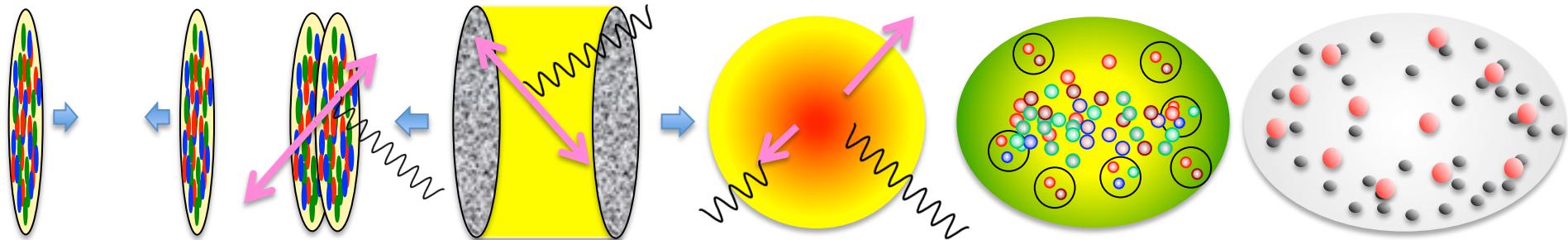
collisions

thermalization

hydro

hadronization

freezeout



Observables: a lot of experimental data at RHIC and LHC

photons/leptons

bulk property

Jets

heavy quarkonia

Phenomenological model

sQGP

Initial condition →

Hydrodynamic model

→ Freezeout process

?

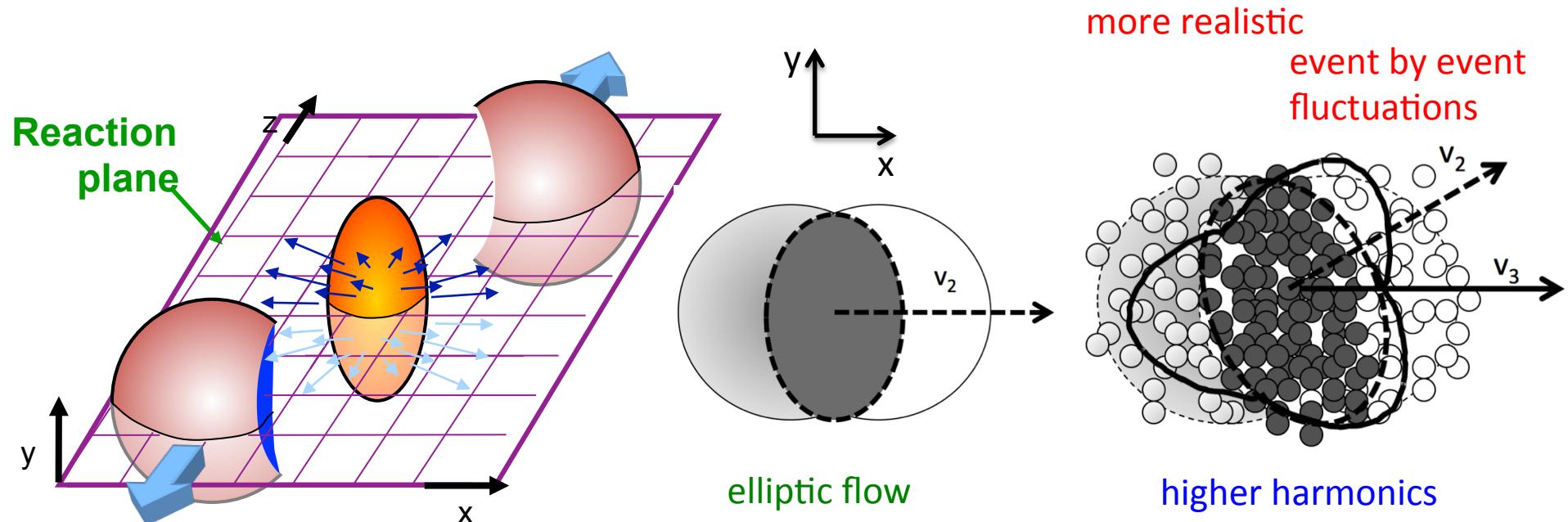
Experimental data

higher harmonics

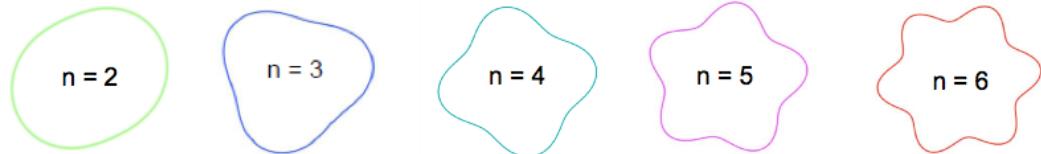


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Higher Harmonics

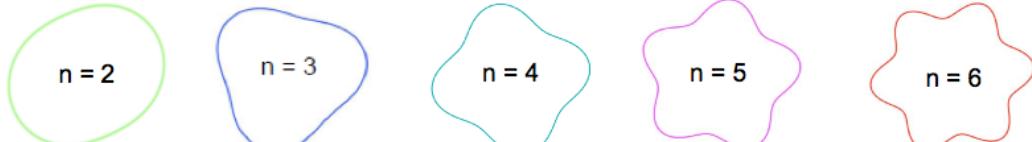


$$\frac{dN}{dyd\phi} \propto 1 + 2v_1 \cos(\phi - \Theta_1) + 2v_2 \cos 2(\phi - \Theta_2) + 2v_3 \cos 3(\phi - \Theta_3) + 2v_4 \cos 4(\phi - \Theta_4) + \dots$$

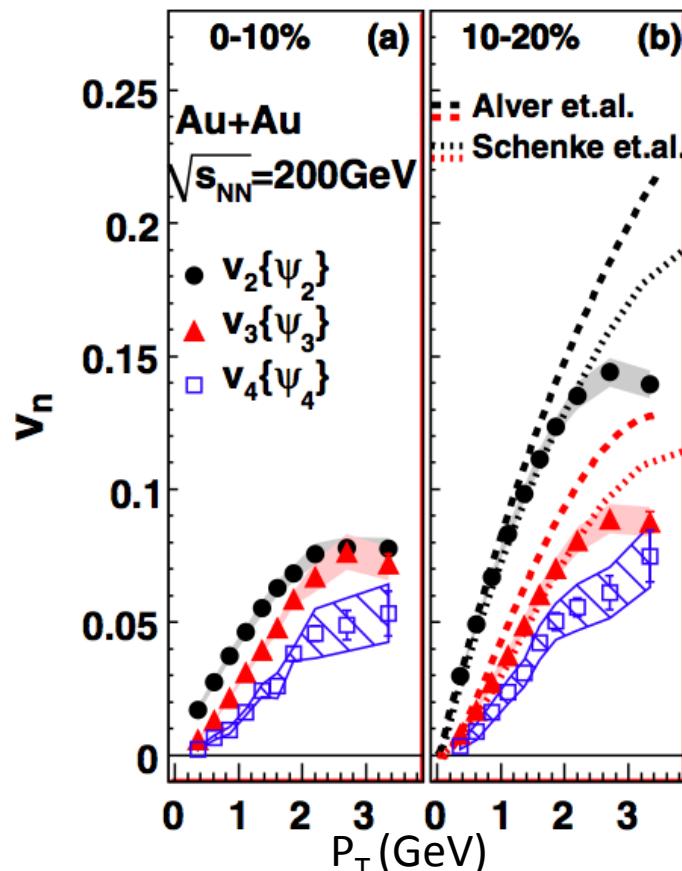


Higher Harmonics @ RHIC & LHC

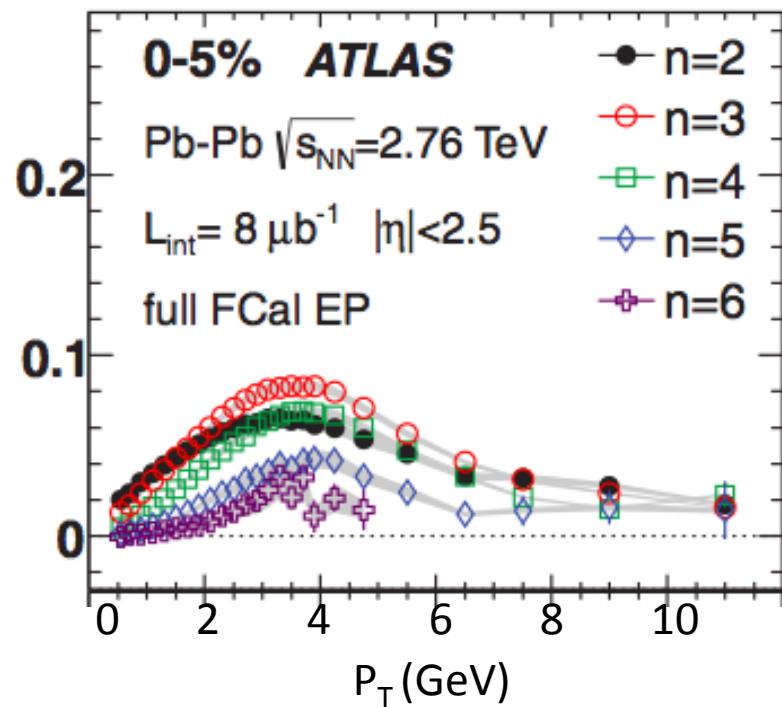
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PHENIX@RHIC, PRL107,252301(2011)

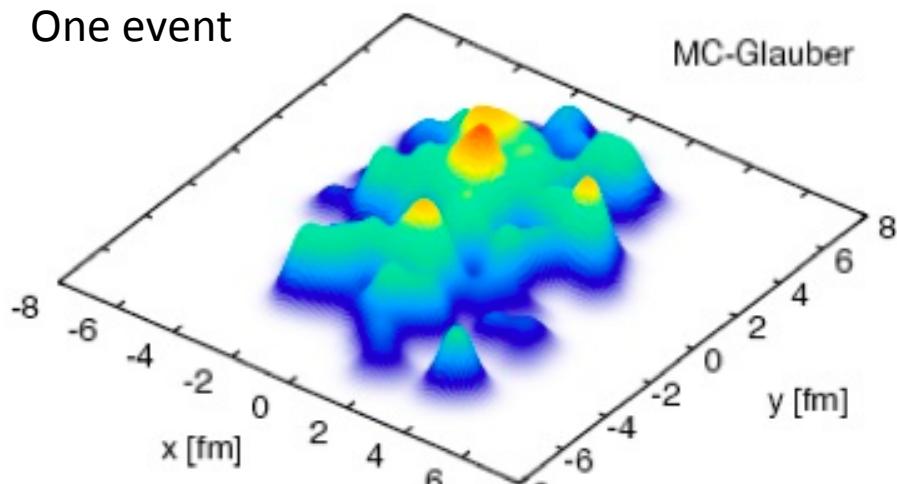


ATLAS@LHC, PRC86,014907(2012)

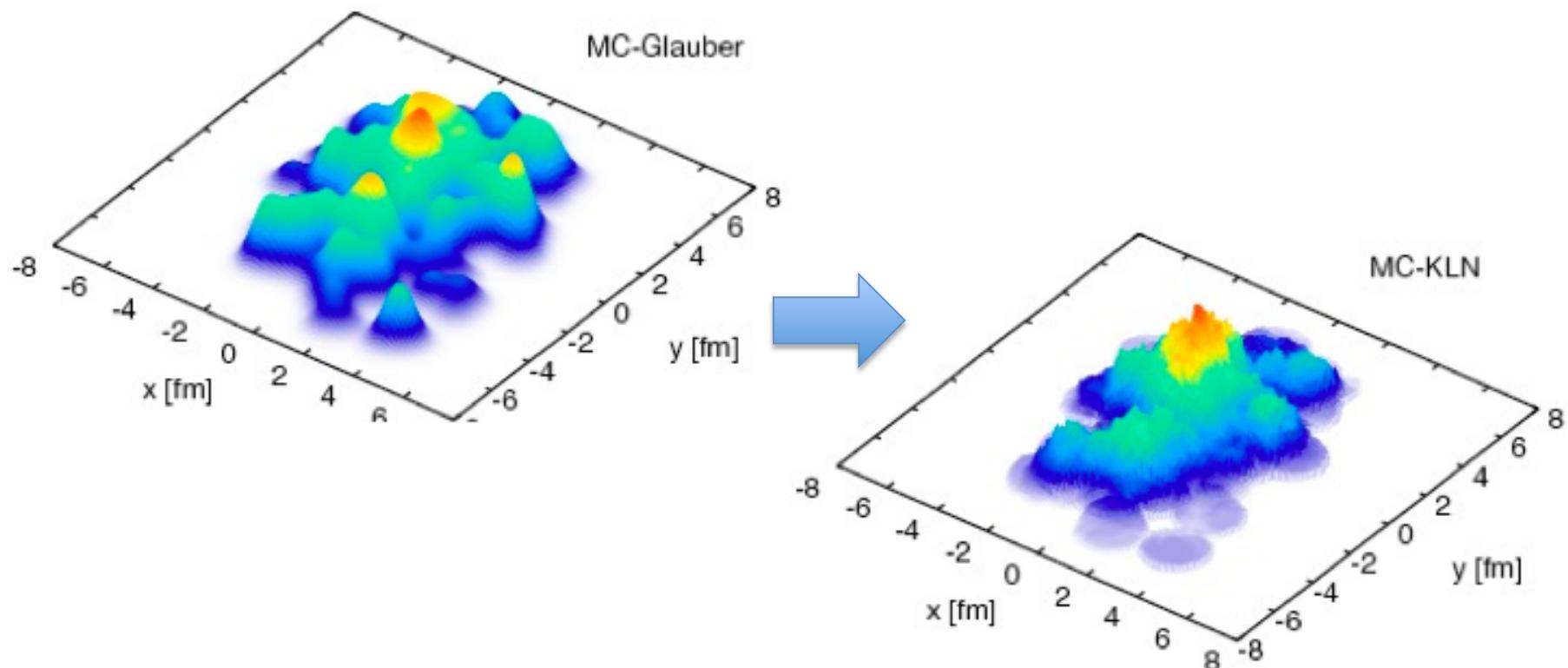


Initial Conditions

One event

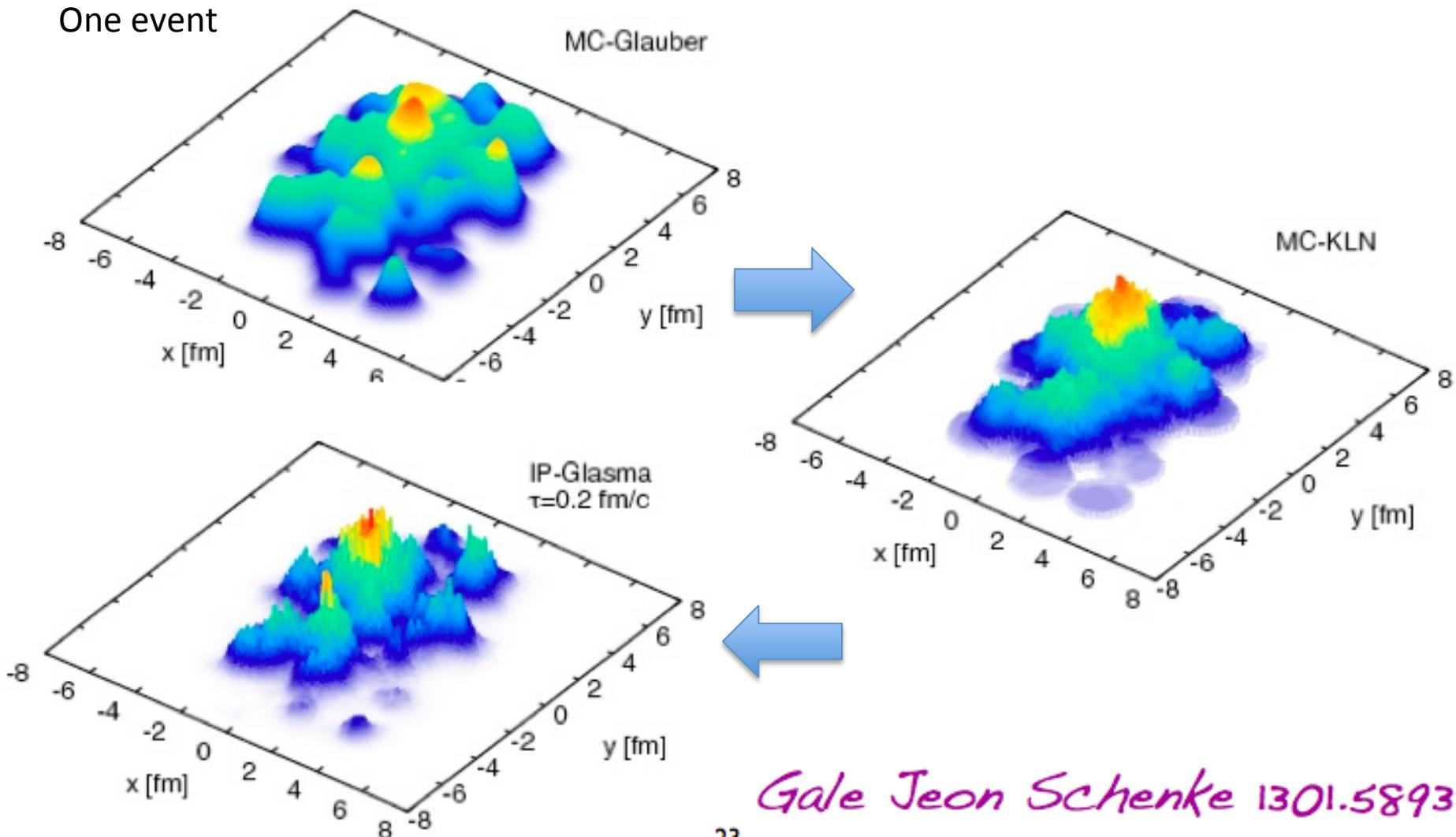


Initial Conditions



Initial Conditions

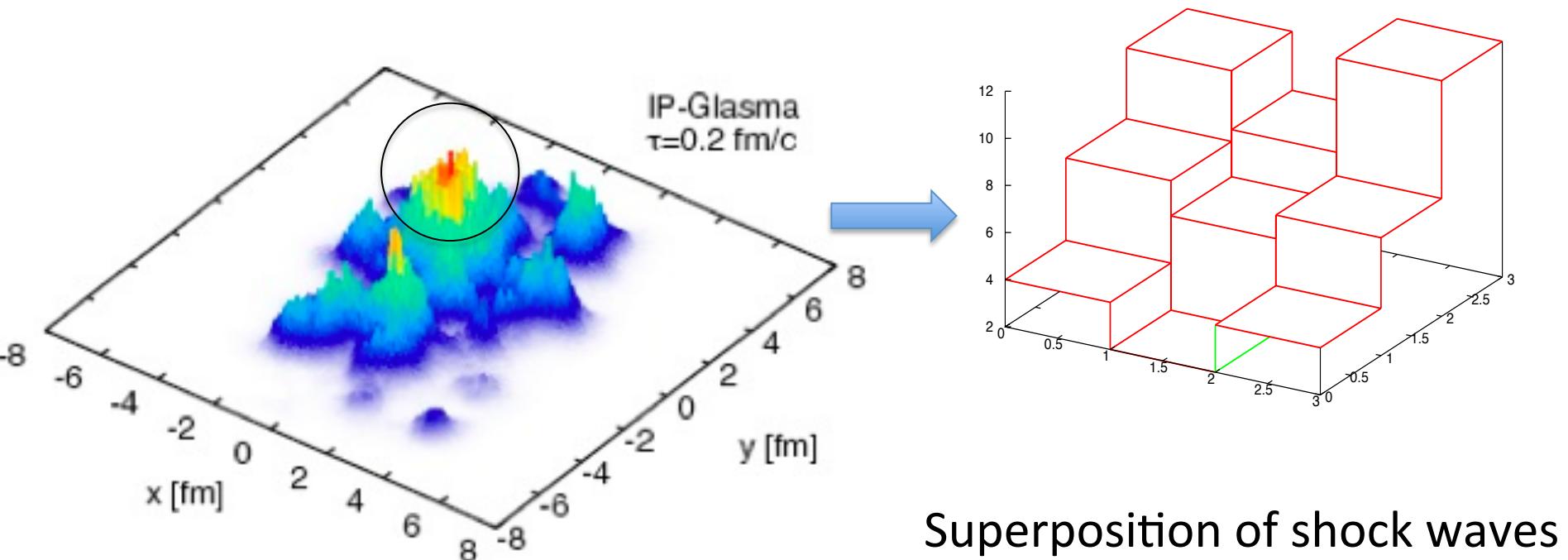
One event



23

Gale Jeon Schenke 1301.5893

Numerical Scheme



Numerical algorithm for hydrodynamic evolution

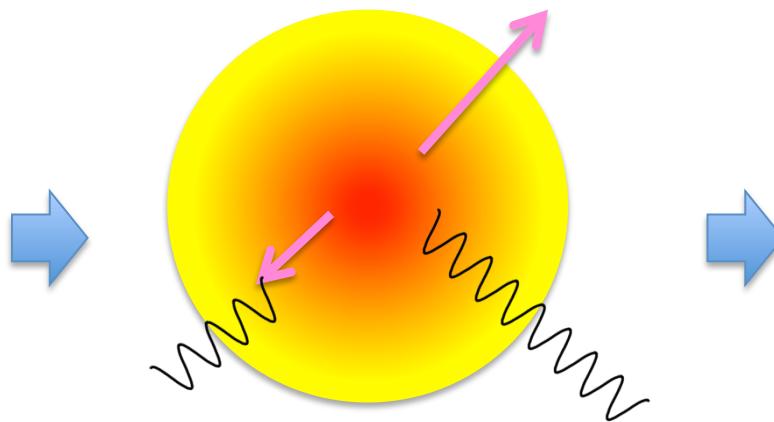


- ✓ shock-wave capturing scheme
- ✓ stable
- ✓ less numerical viscosity

Hydrodynamic Expansion



fluctuating
initial conditions



v_n

Bulk property
transport coefficients

importance of numerical algorithm !

Akamatsu, Inutsuka, CN, Takamoto:
arXiv:1302.1665, J. Comp. Phys. (2014) 34

HYDRODYNAMIC MODEL

Viscous Hydrodynamic Model

- Relativistic viscous hydrodynamic equation

$$\partial_\mu T^{\mu\nu} = 0$$

- First order in gradient: acausality
- Second order in gradient:

- Israel-Stewart, Ottinger and Grmela, AdS/CFT,

- Grad's 14-momentum expansion, Renormalization group

- Numerical scheme

- Shock-wave capturing schemes: Riemann problem

- Godunov scheme: analytical solution of Riemann problem

- SHASTA: the first version of Flux Corrected Transport algorithm, Song, Heinz, Pang, Victor...

- Kurganov-Tadmor (KT) scheme, McGill

Our Approach

- Israel-Stewart Theory

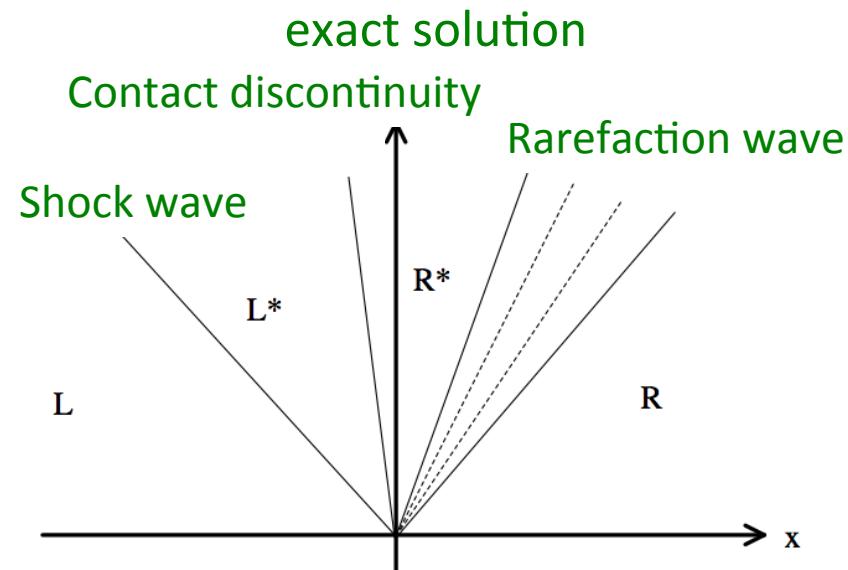
1. dissipative fluid dynamics = advection + dissipation



Takamoto and Inutsuka, arXiv:1106.1732

Akamatsu, Inutsuka, CN, Takamoto, arXiv:1302.1665

(ideal hydro)



Riemann solver: Godunov method

Two shock approximation

Mignone, Plewa and Bodo, *Astrophys. J.* S160, 199 (2005)

Rarefaction wave → shock wave

2. relaxation equation = advection + stiff equation

Numerical Scheme

- Israel-Stewart Theory

Takamoto and Inutsuka, arXiv:1106.1732

1. Dissipative fluid equation

$$\partial_\mu T^{\mu\nu} = 0$$

$$T^{\mu\nu} = (\epsilon + p)u^\mu u^\nu - pg^{\mu\nu} + q^\mu u^\nu + q^\nu u^\mu + \tau^{\mu\nu}$$
$$= T_{\text{ideal}} + T_{\text{dissip}}$$

$$\partial_t U + \nabla \cdot F(U) = 0 \quad U = U_{\text{ideal}} + U_{\text{dissip}}$$

2. Relaxation equation

$$\hat{D}\Pi = \frac{1}{\tau_\Pi}(\Pi_{NS} - \Pi) - I_\Pi, \quad \Rightarrow \quad \left(\frac{\partial}{\partial t} + v^j \frac{\partial}{\partial x^j} \right) \Pi = -\frac{I_\Pi}{\gamma}, \quad + \quad \frac{\partial}{\partial t} \Pi = \frac{1}{\gamma \tau_\Pi}(\Pi_{NS} - \Pi),$$

$$\hat{D}\pi^{\mu\nu} = \frac{1}{\tau_\pi}(\pi_{NS}^{\mu\nu} - \pi^{\mu\nu}) - I_\pi^{\mu\nu}, \quad \text{advection} \quad \Delta t < \tau_{\text{relax}} \ll \tau_{\text{fluid}}$$

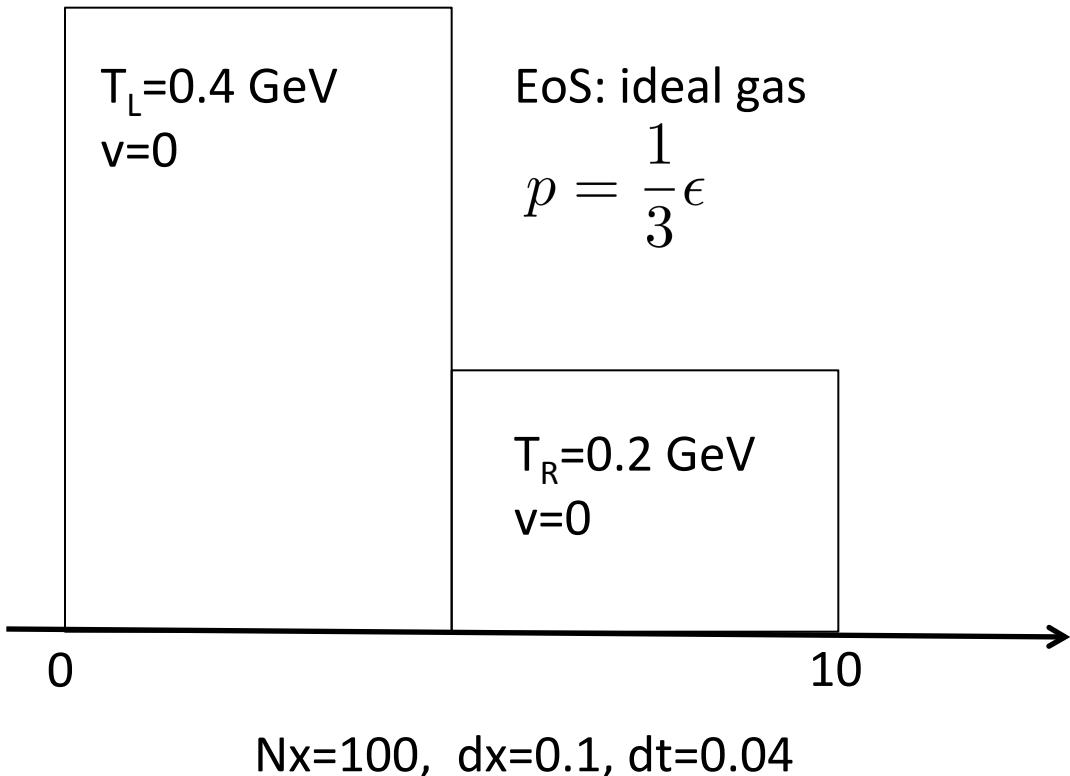
$$\hat{D}q^\mu = \frac{1}{\tau_q}(q_{NS}^\mu - q^\mu) - I_q^\mu,$$

$\hat{D} = u^\mu \partial_u$ l: second order terms

$$\tau^{\mu\nu} = \Pi \Delta^{\mu\nu} + \pi^{\mu\nu}$$

Comparison

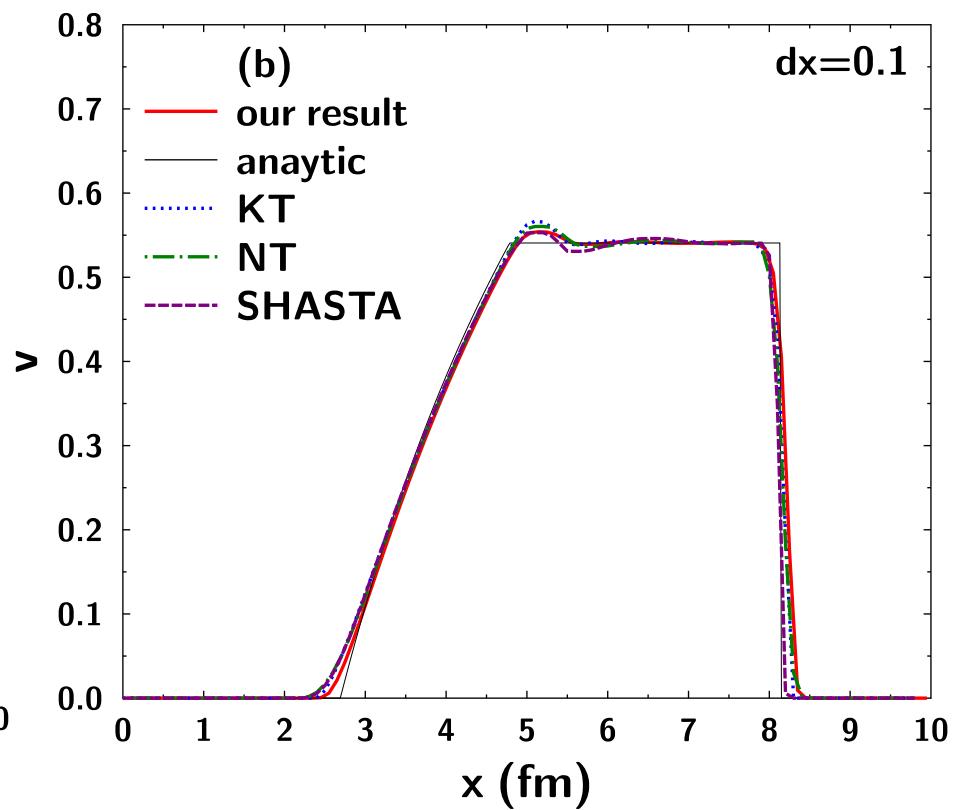
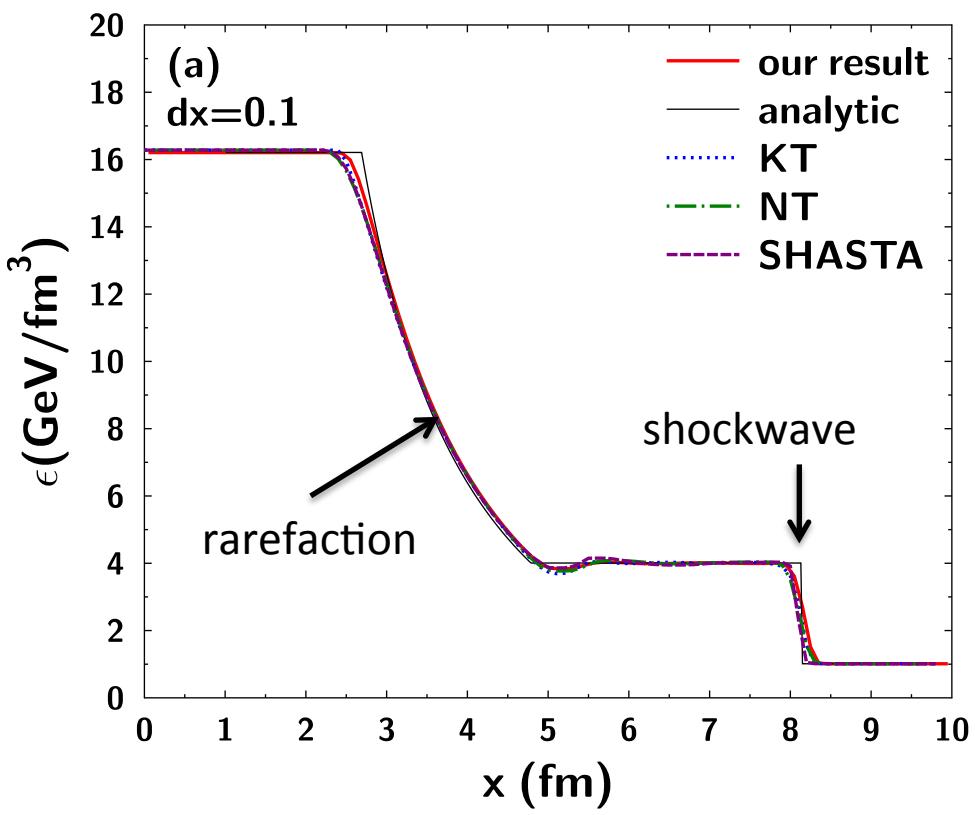
- Shock Tube Test : *Molnar, Niemi, Rischke*, Eur.Phys.J.C65,615(2010)



- Analytical solution
 - Numerical schemes
SHASTA, KT, NT
Our scheme

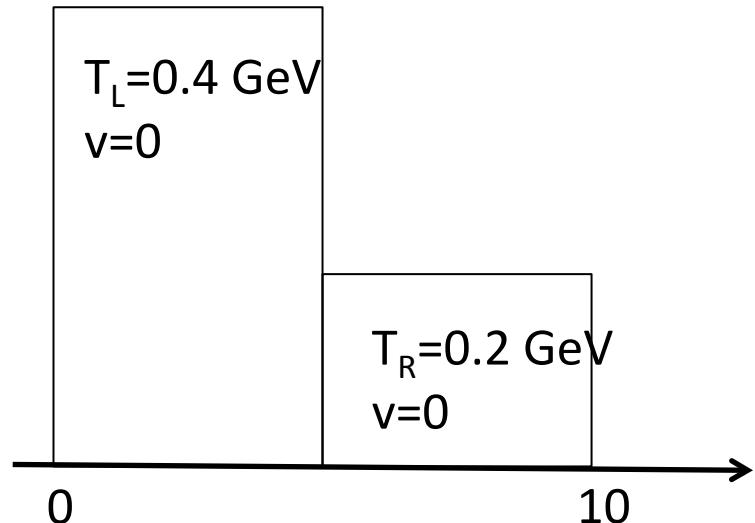
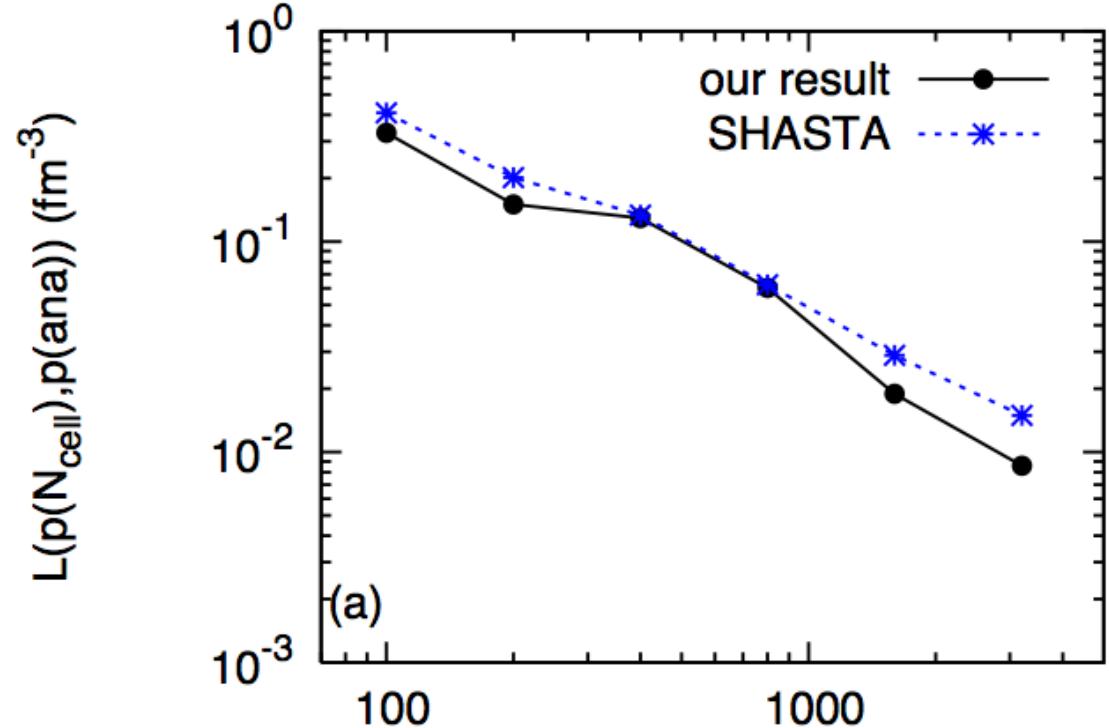
Shocktube problem

- Ideal case



L1 Norm

- Numerical dissipation: deviation from analytical solution



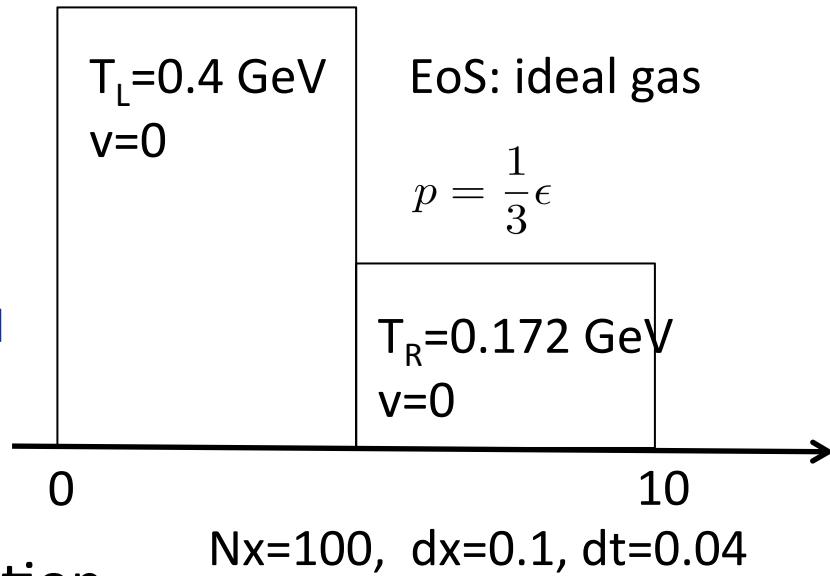
For analysis of heavy ion collisions

$$L(p(N_{\text{cell}}), p(\text{analytic})) = \sum_{i=1}^{N_{\text{cell}}} |p(N_{\text{cell}}) - p(\text{analytic})| \frac{\lambda}{N_{\text{cell}}}$$

$\lambda = 10 \text{ fm}$

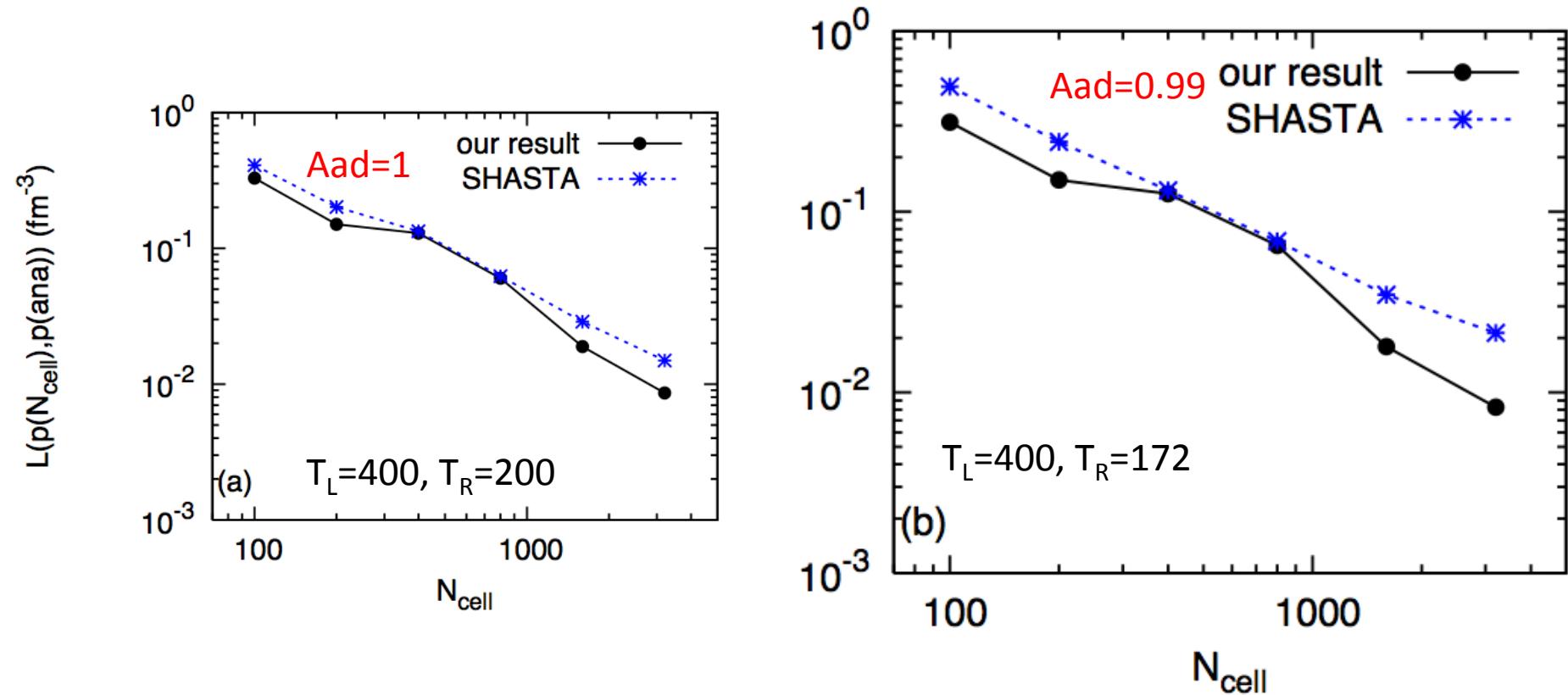
Large ΔT difference

- $T_L=0.4 \text{ GeV}$, $T_R=0.172 \text{ GeV}$
 - SHASTA becomes unstable.
 - Our algorithm is stable.
- SHASTA: anti diffusion term, A_{ad}
 - $A_{ad} = 1$: default value, unstable
 - $A_{ad} = 0.99$: stable,
more numerical dissipation



L1 norm

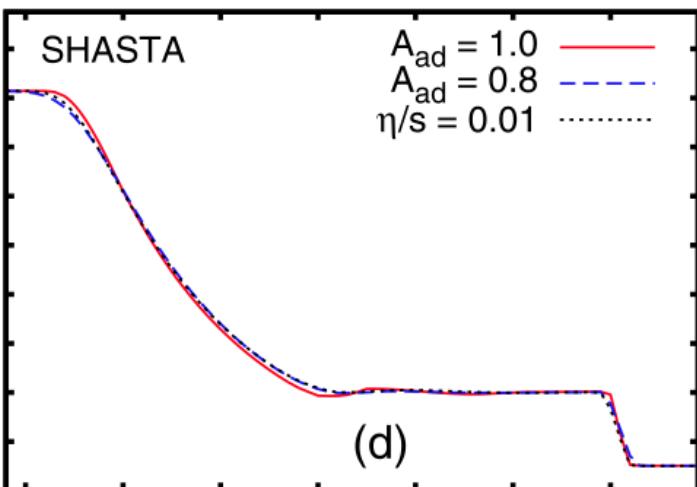
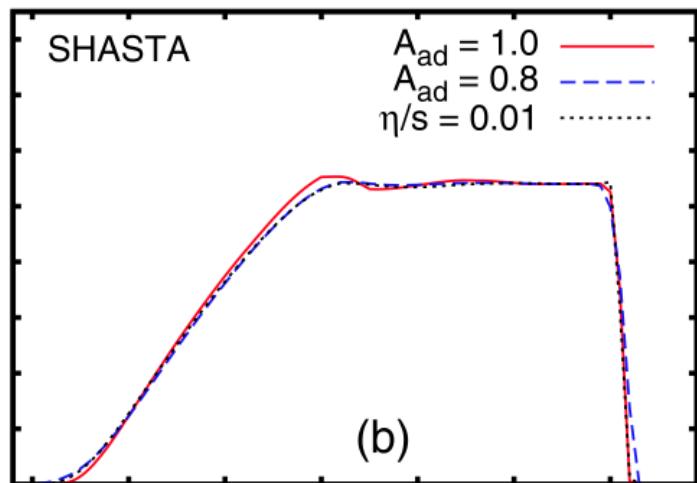
- SHASTA with small A_{ad} has large numerical dissipation



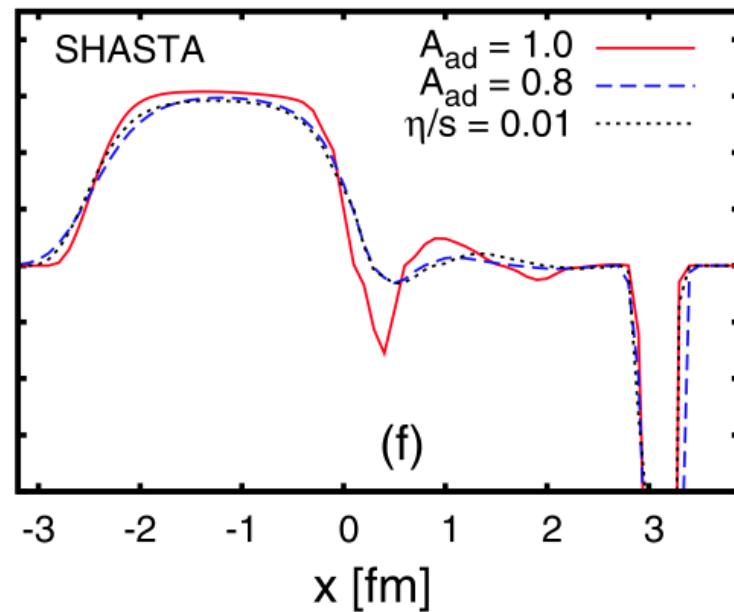
$$L(p(N_{cell}), p(\text{analytic})) = \sum_{i=1}^{N_{cell}} |p(N_{cell}) - p(\text{analytic})| \frac{\lambda}{N_{cell}}$$

$\lambda = 10 \text{ fm}$

Artificial and Physical Viscosities



Molnar, Niemi, Rischke, Eur.Phys.J.C65, 615(2010)

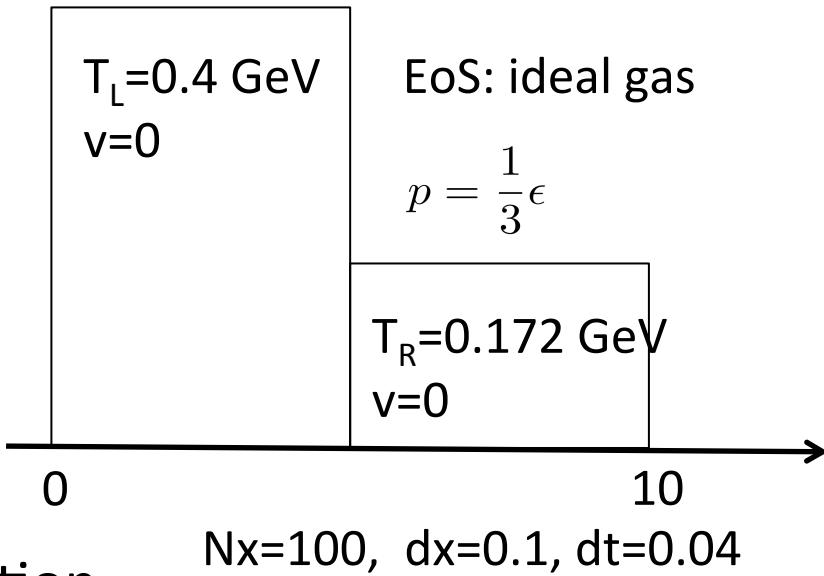


Antidiffusion terms : artificial viscosity stability

$$U_i^{n+1} = \tilde{U}_i - \tilde{A}_i + \tilde{A}_{i-1}$$
$$A_i = A_{ad} \tilde{\Delta}_i / 8$$

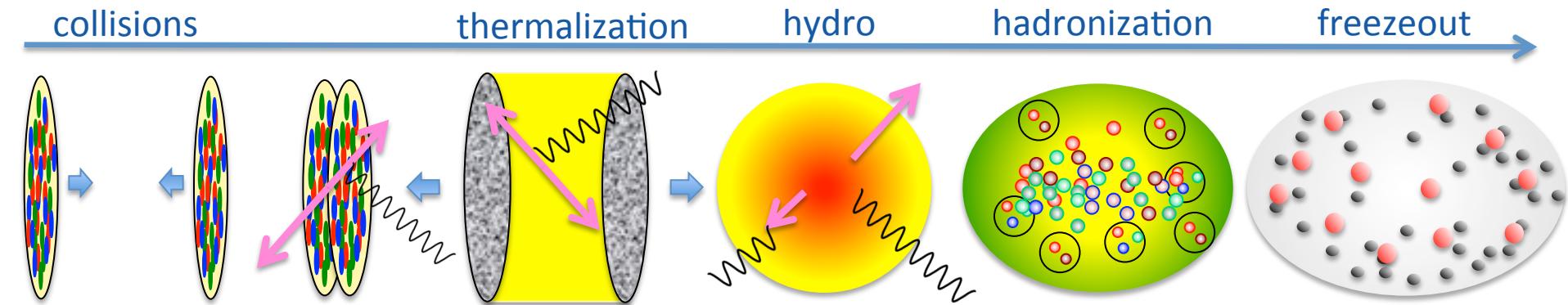
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 - Our algorithm is stable.
- SHASTA: anti diffusion term, A_{ad}
 - $A_{ad} = 1$: default value
 - $A_{ad} = 0.99$: stable,
more numerical dissipation
- Large fluctuation (ex initial conditions)
 - Our algorithm is stable even with small numerical dissipation.



DYNAMICAL MODEL

Our Dynamical Model



Fluctuating Initial conditions

Hydrodynamic expansion

Freezeout process

- From Hydro to particle
- Final state interactions

*Akamatsu, Inutsuka, CN, Takamoto,
arXiv:1302.1665, J. Comp. Phys. (2014) 34*

MC-KLN

hydrodynamic model

Cornelius

Oscar sampler

Nara

Freezeout hypersurface finder

Ohio group

<http://www.aiu.ac.jp/~ynara/mckln/>

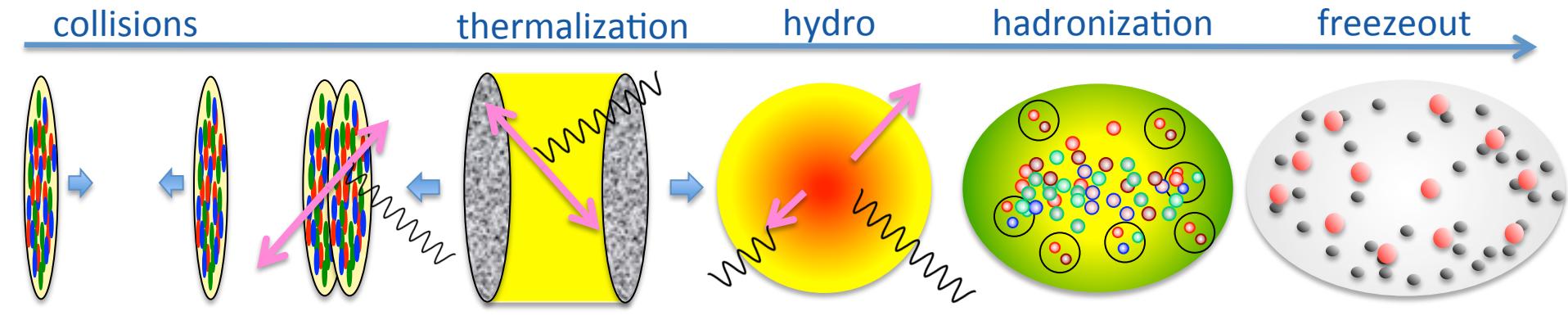
Huovinen, Petersen

UrQMD



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Simulation setups:

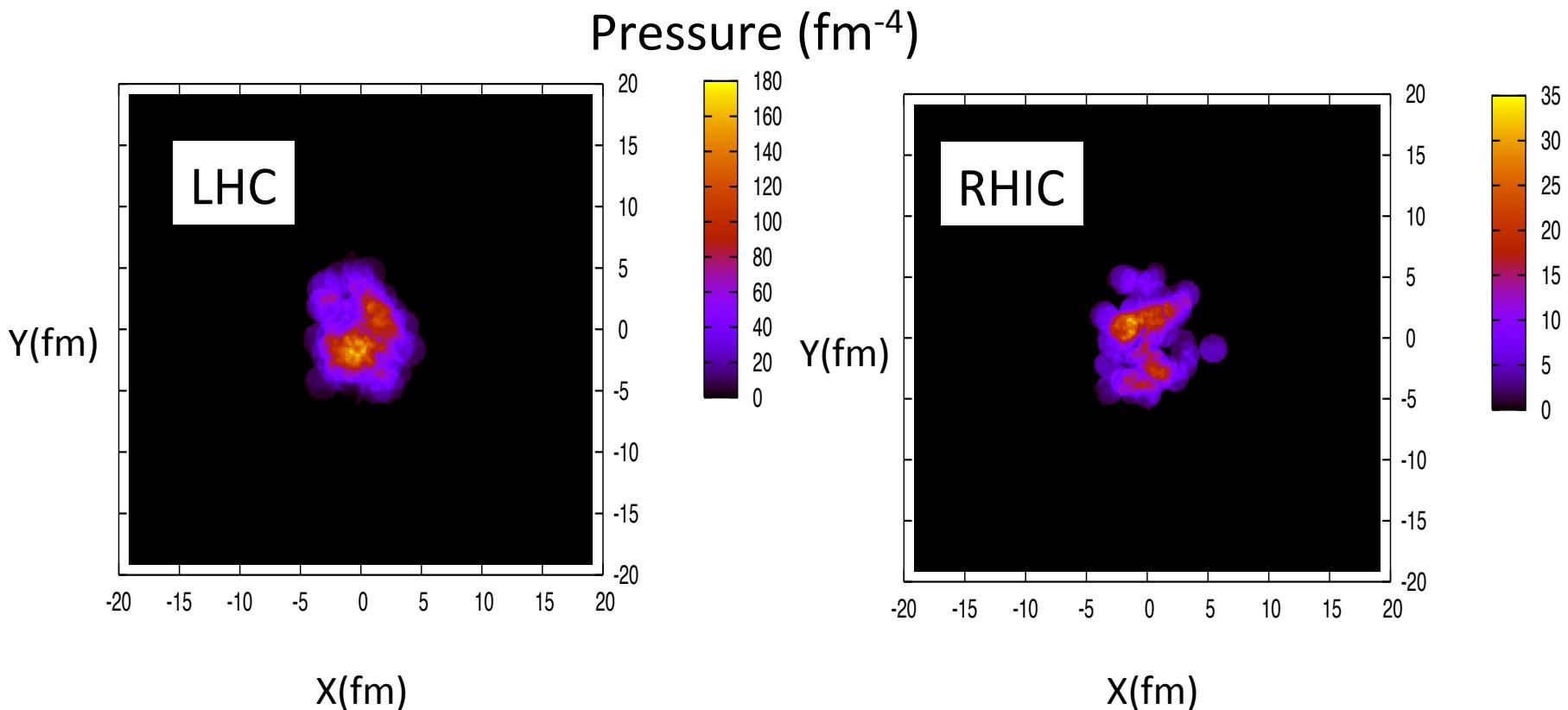
- Free gluon EoS
- Hydro in 2D boost invariant simulation
- UrQMD with $|y| < 0.5$



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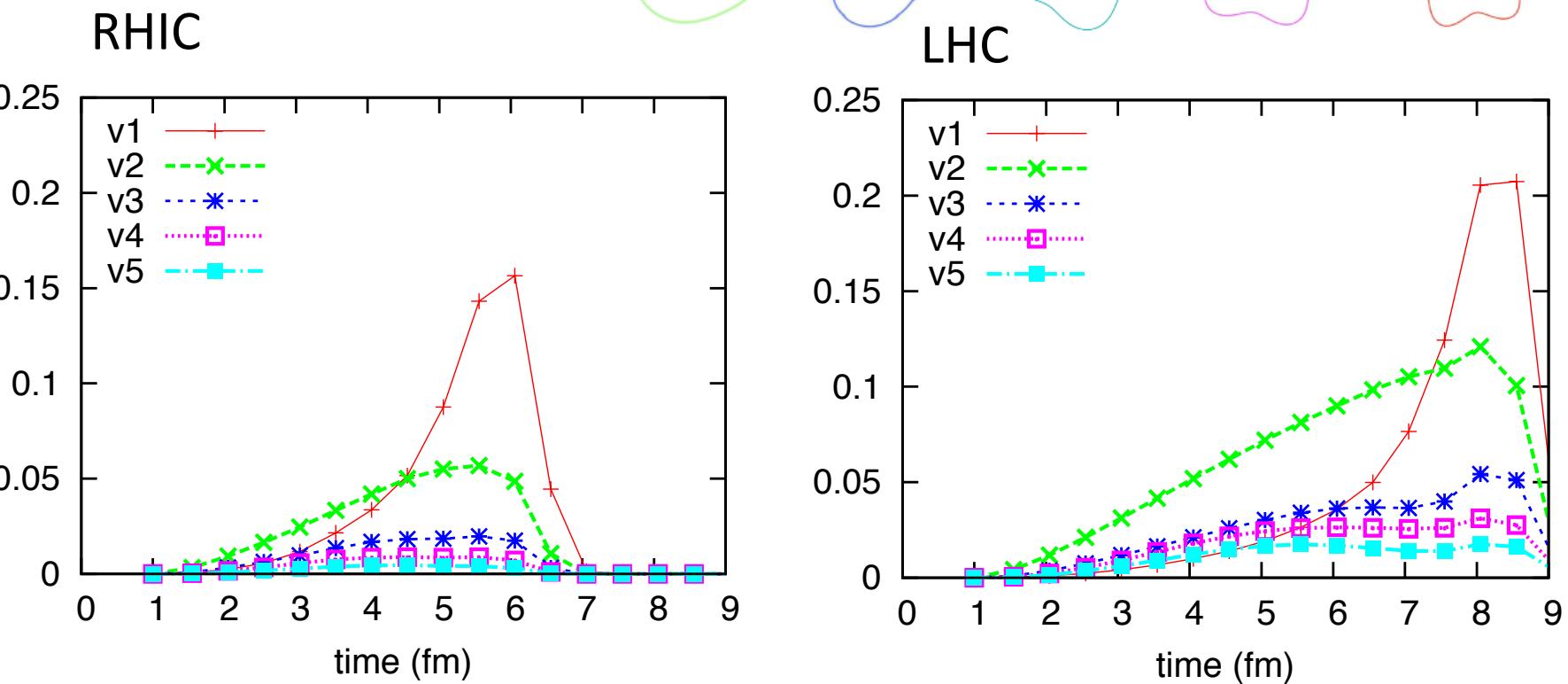
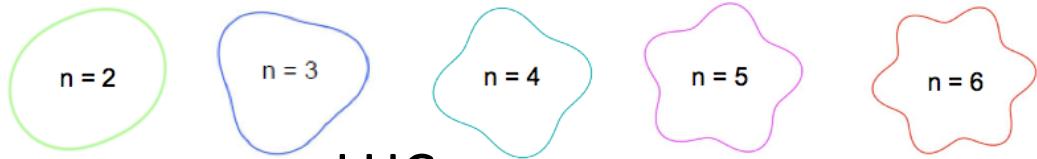
Initial Pressure Distribution

- MC-KLN (centrality 15-20%)



Time Evolution of v_n

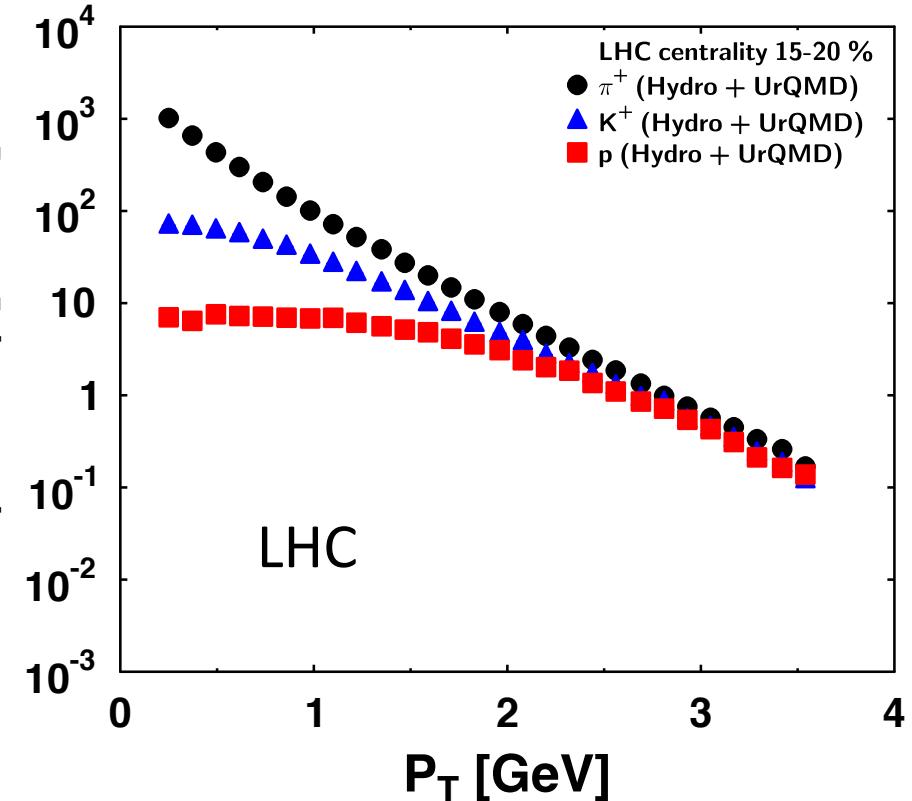
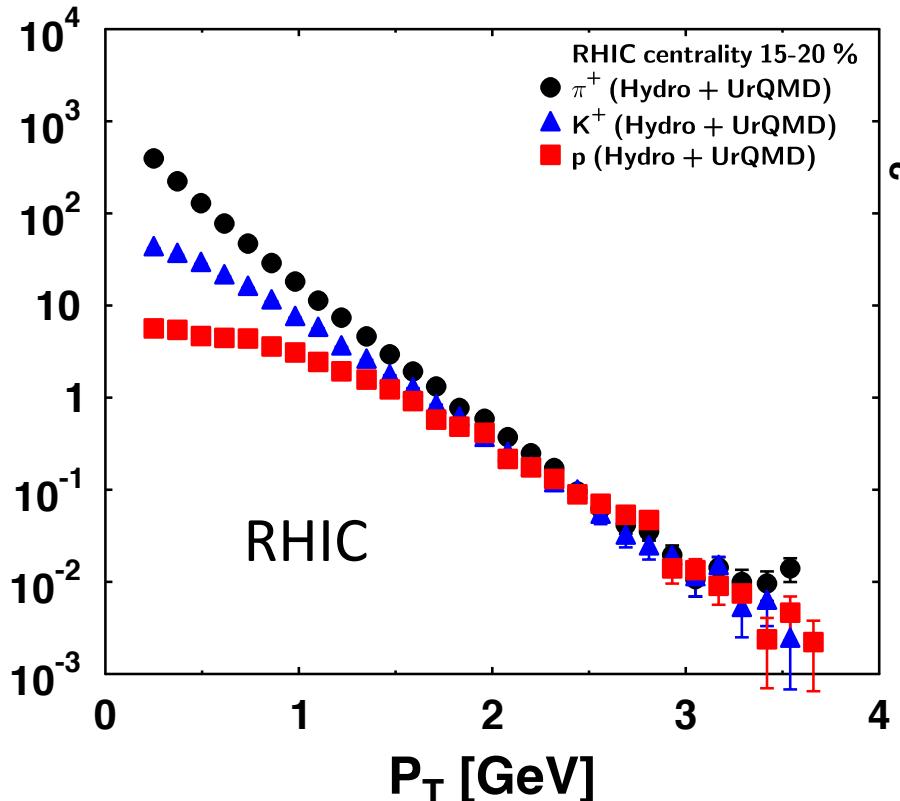
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- Qualitatively RHIC \sim LHC
- v_2 is dominant
- $v_2 > v_3 > v_4 > v_5$

Hydro + UrQMD

- Transverse momentum spectrum

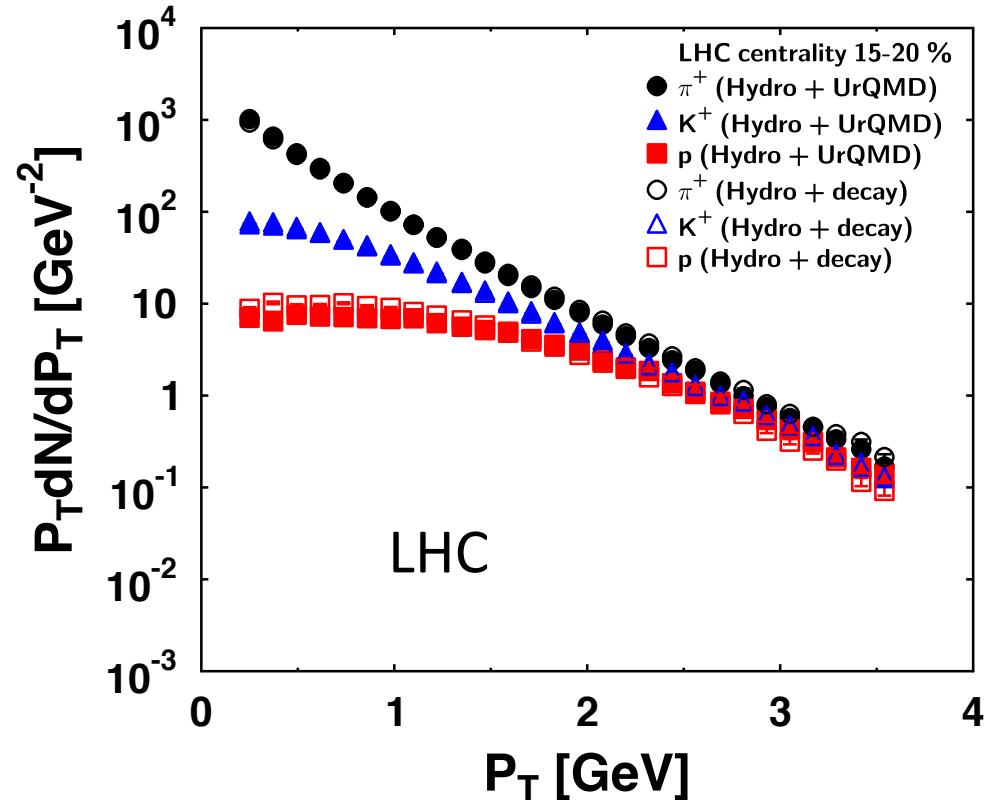
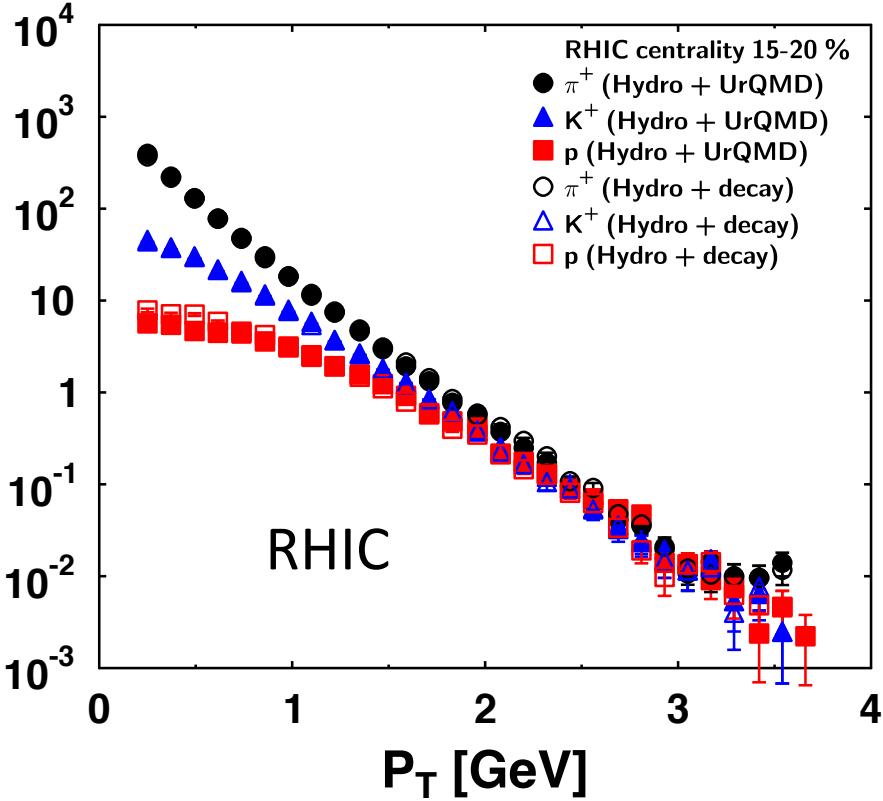


- P_T distribution at LHC has flatter slope

→ Larger radial flow at LHC

Effect of Hadronic Interaction

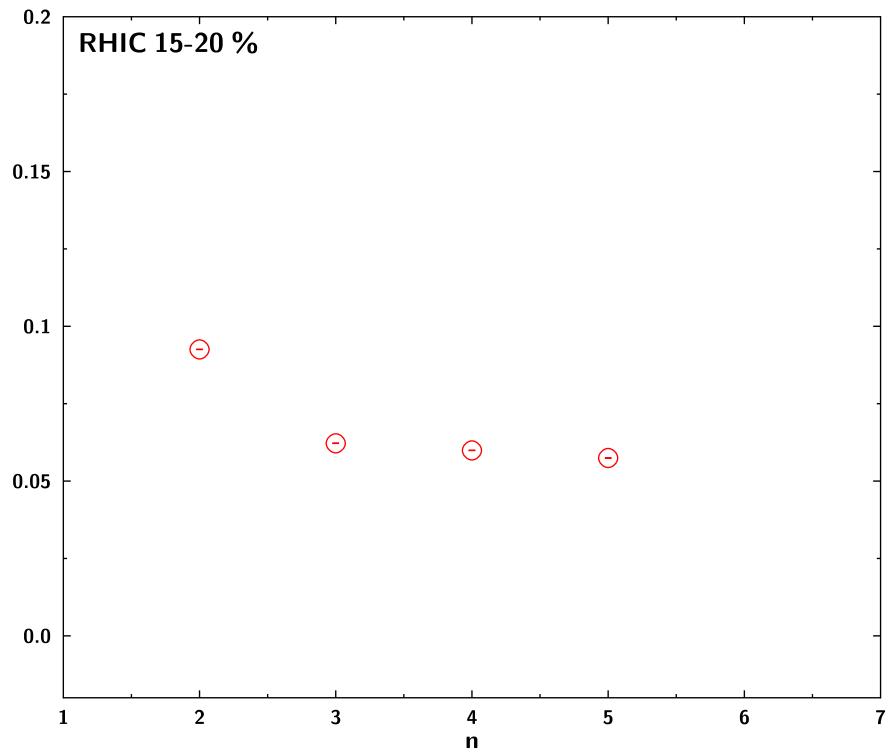
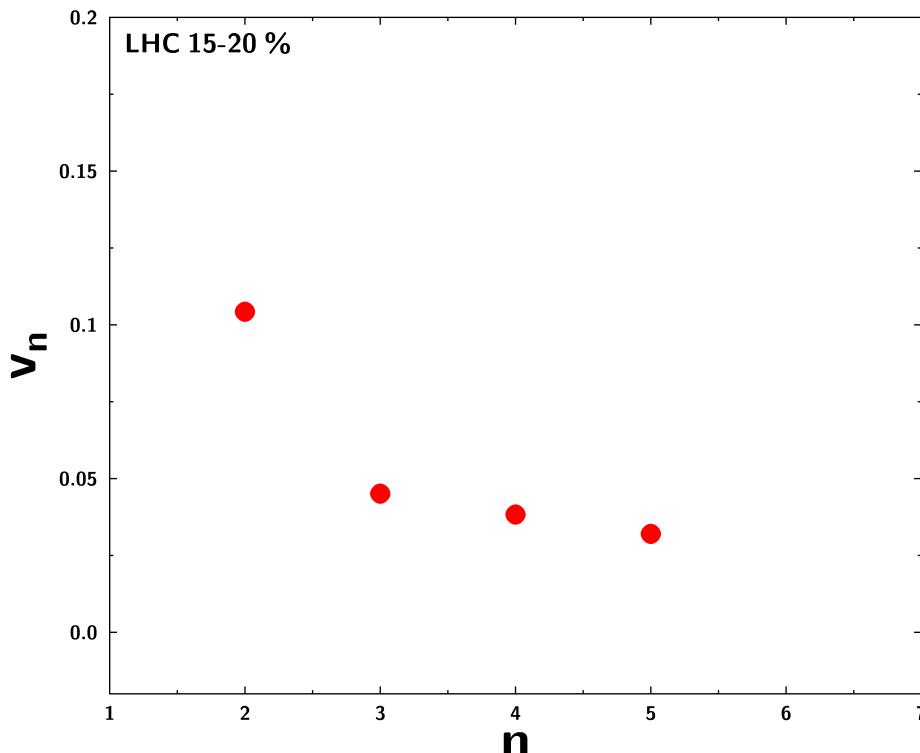
- Transverse momentum distribution



- Effect of final state interactions is small
- Slope of proton Pt spectra become flatter

Higher harmonics from Hydro + UrQMD

- Effect of hadronic interaction



Summary

Importance of numerical scheme in Hydrodynamic Models

- We develop a state-of-the-art numerical scheme
 - Shock wave capturing scheme: Godunov method

Our algorithm

- Less artificial diffusion: crucial for viscosity analyses
- Stable for strong shock wave

- Construction of a hybrid model
 - Fluctuating initial conditions + Hydrodynamic evolution +

UrQMD

- Higher Harmonics
 - Time evolution, hadron interaction