

Flavor Physics in the Beyond Standard Models

Yuji Omura

Based on PLB744 (2015) 395 with J. Hisano, Y. Muramatsu, M. Yamanaka,
and (arXiv:1505.07636) with T. Kobayashi, F. Takayama, D. Yasuhara

Standard Model $(SU(3)_C \times SU(2)_L \times U(1)_Y)$

is very successful in particle physics

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spin-1/2

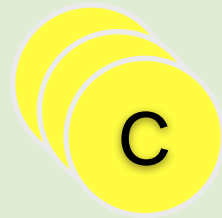
quarks

$SU(3)_c$ -charged

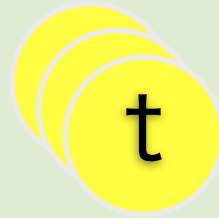
EM-charge



u

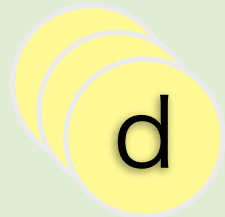


c

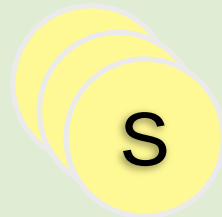


t

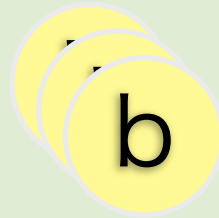
+2/3



d



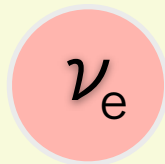
s



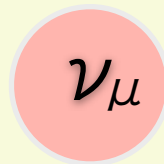
b

-1/3

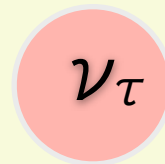
leptons



ν_e



ν_μ

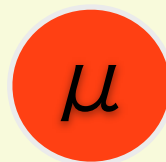


ν_τ

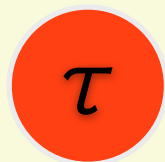
0



e



μ



τ

-1

spin-1

$SU(3)_c$ gauge



g

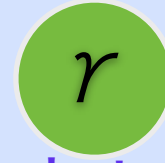
gluon

$SU(2)_L \times U(1)_Y$



Z

massive



γ

photon



$W_{+/-}$

massive

carry forces

spin-0

Higgs

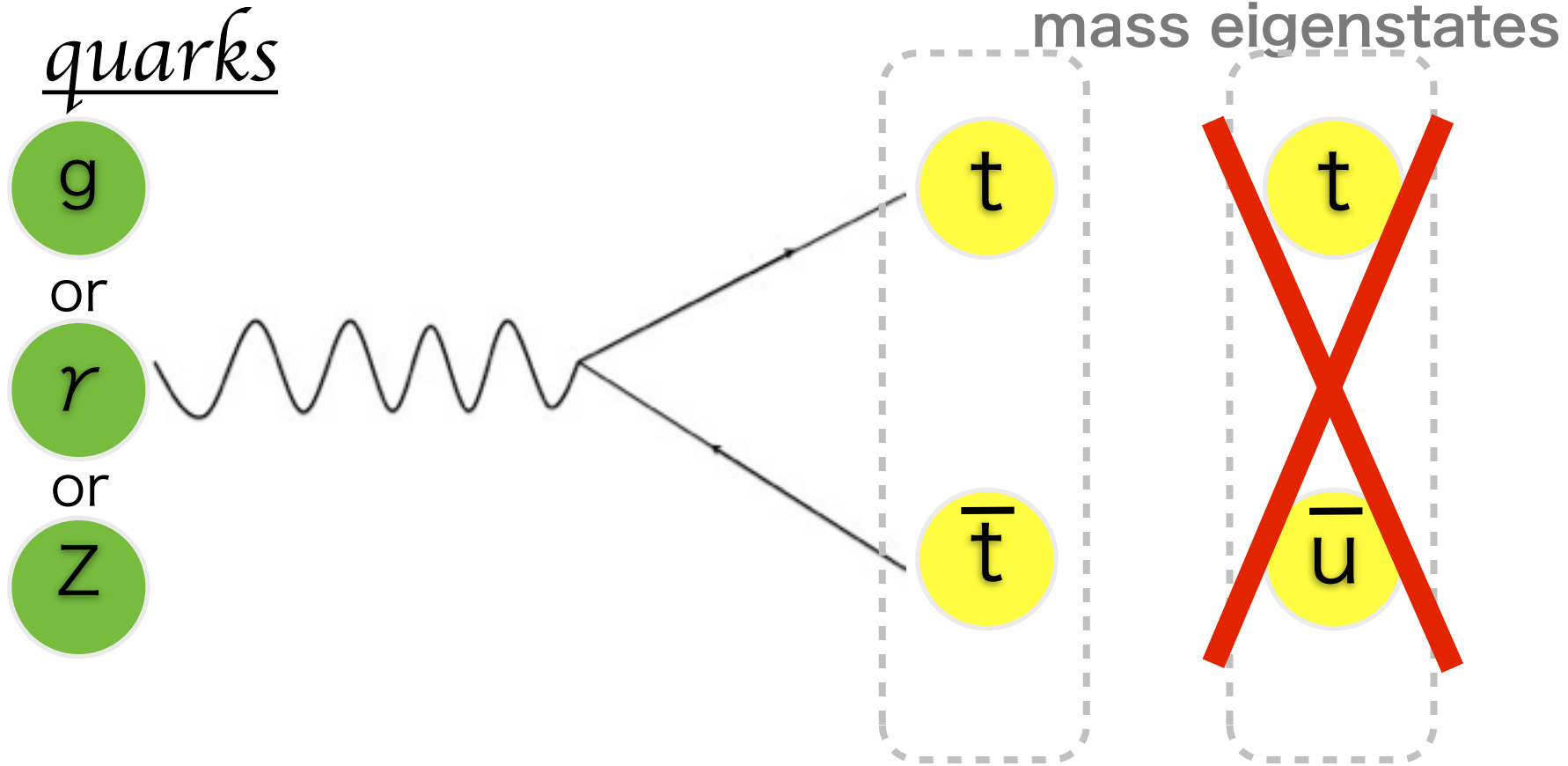


H

breaks

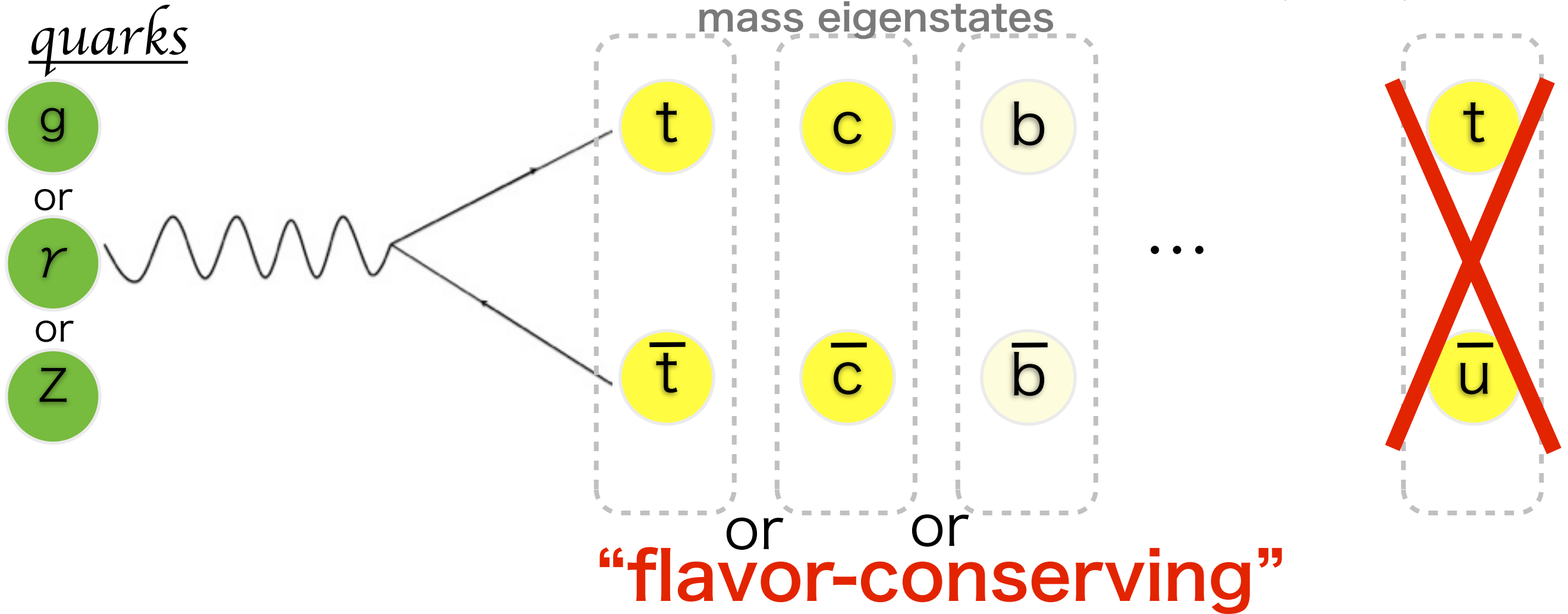
$SU(2)_L \times U(1)_Y$

Interactions in the Standard Model (SM)

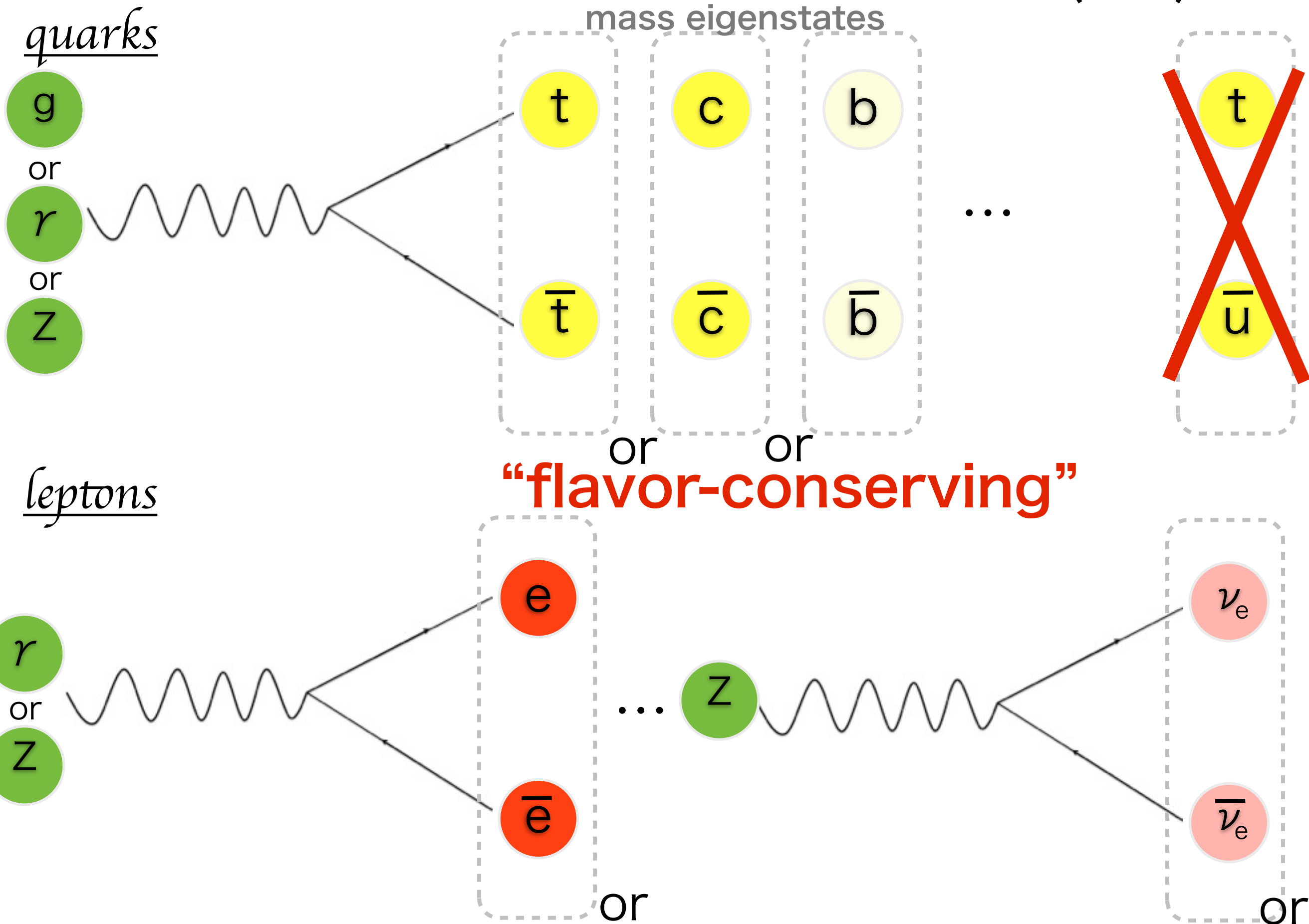


“flavor-conserving”

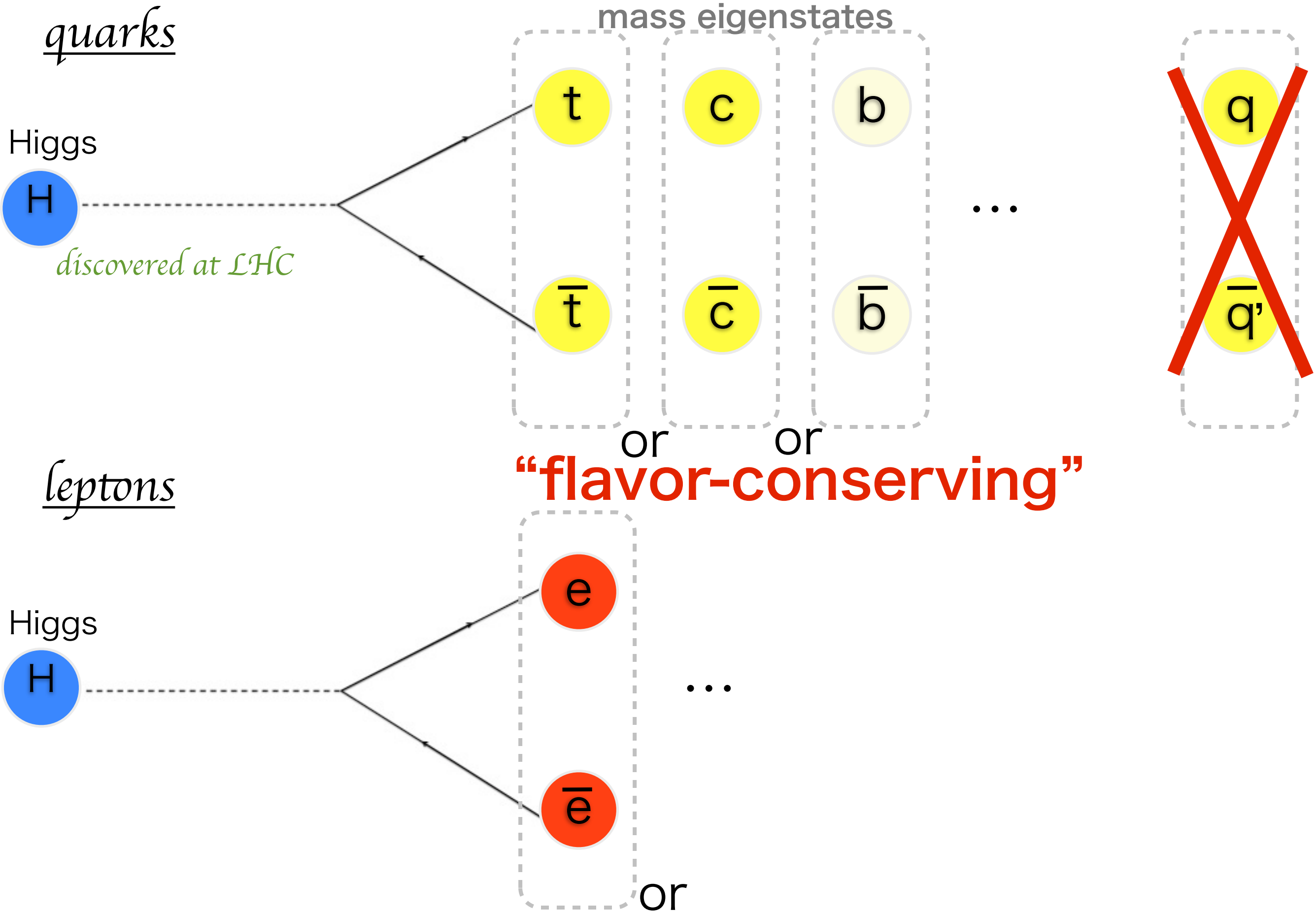
Interactions in the Standard Model (SM)



Interactions in the Standard Model (SM)

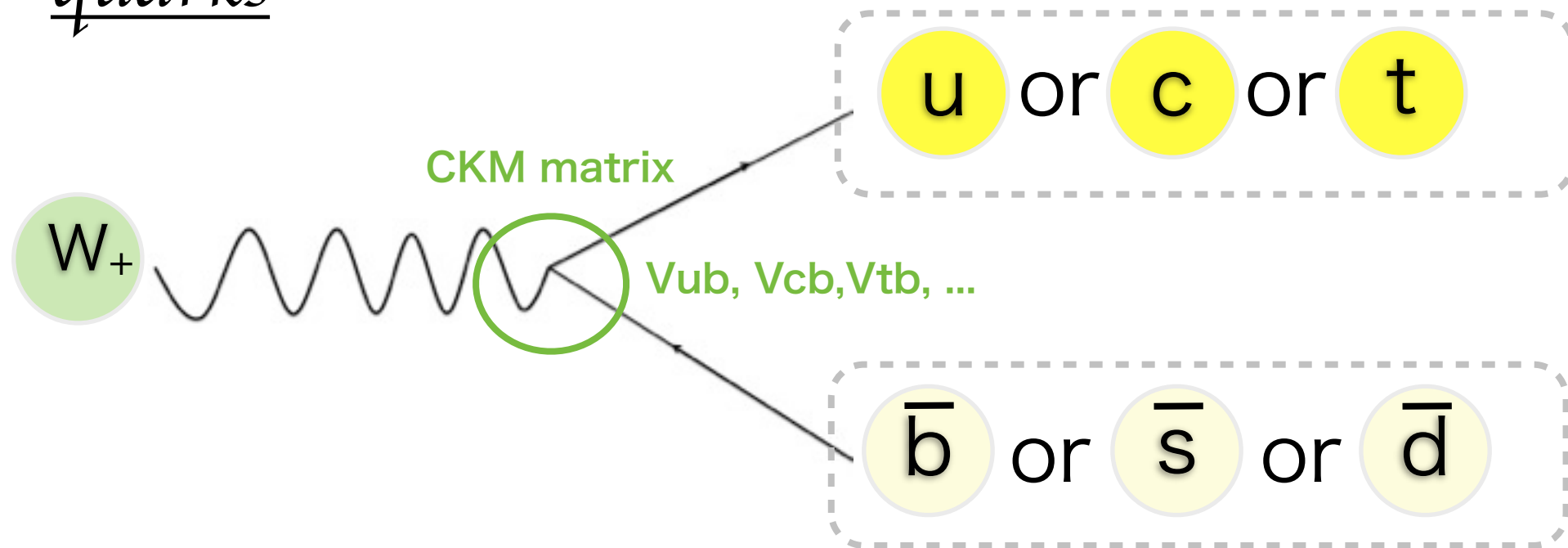


Yukawa Interactions in the SM



W-Interaction in the SM

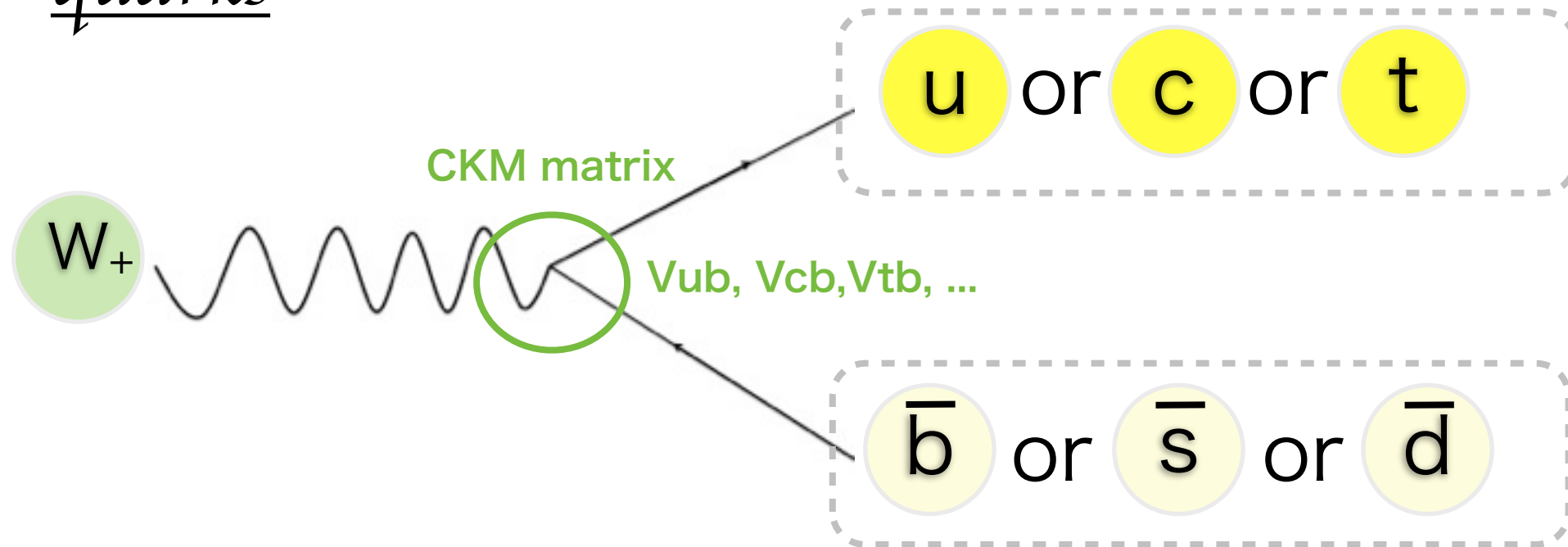
quarks



“flavor-changing”

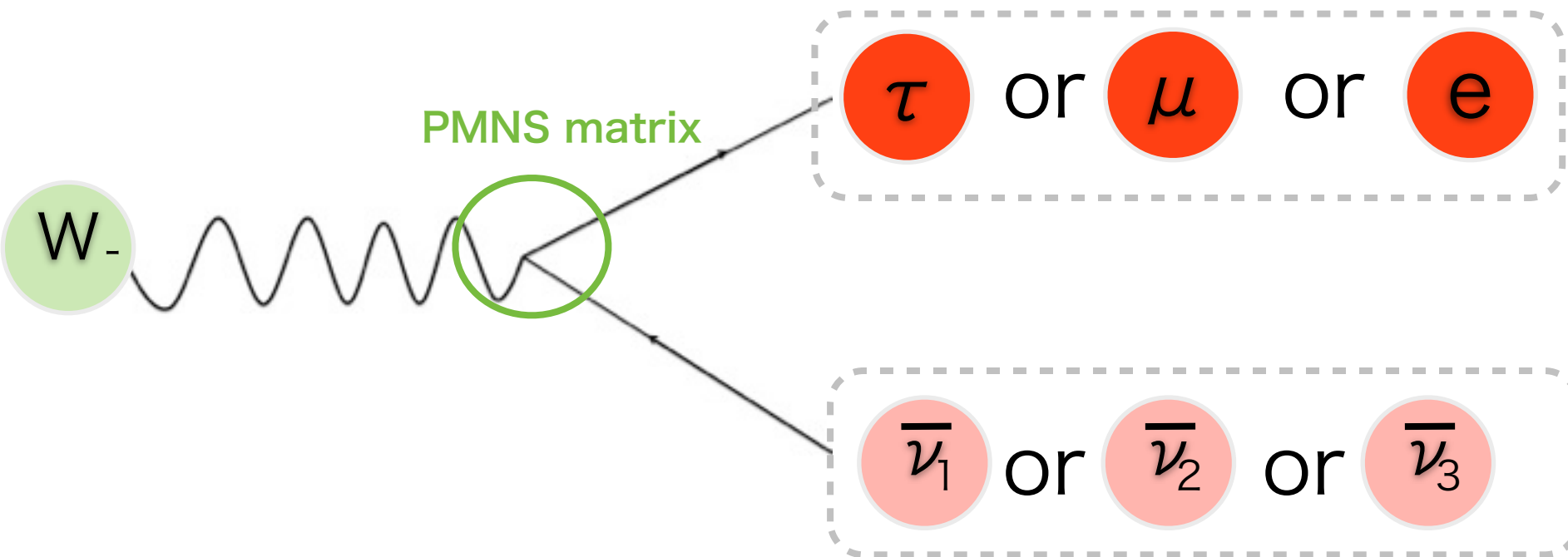
W-Interaction in the SM

quarks

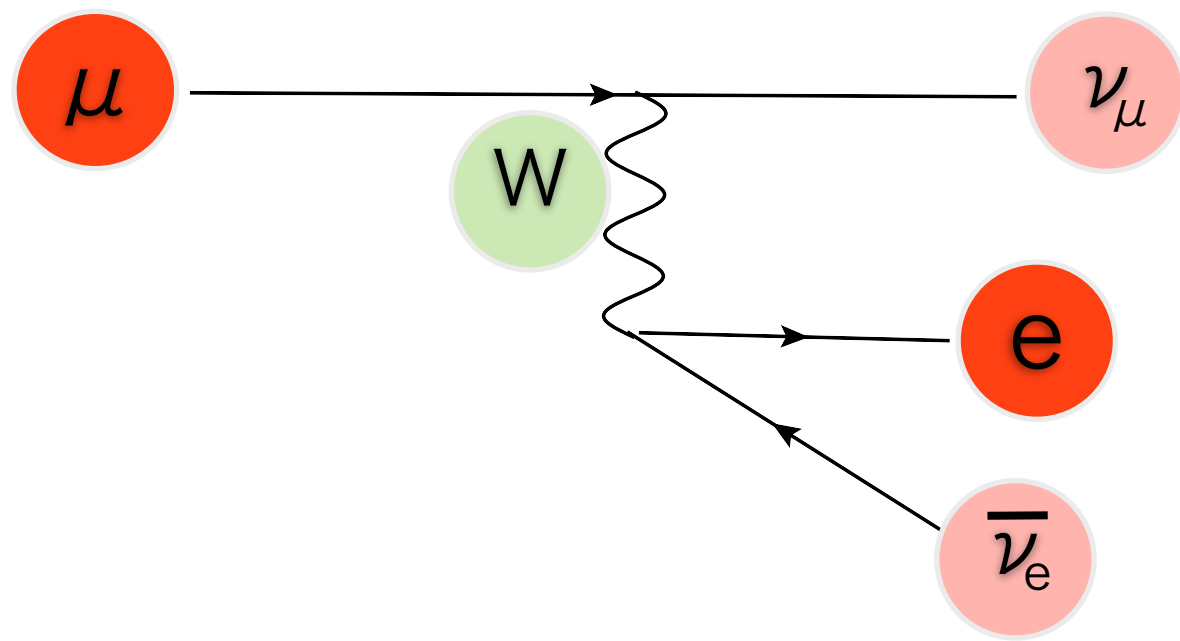


“flavor-changing”

leptons

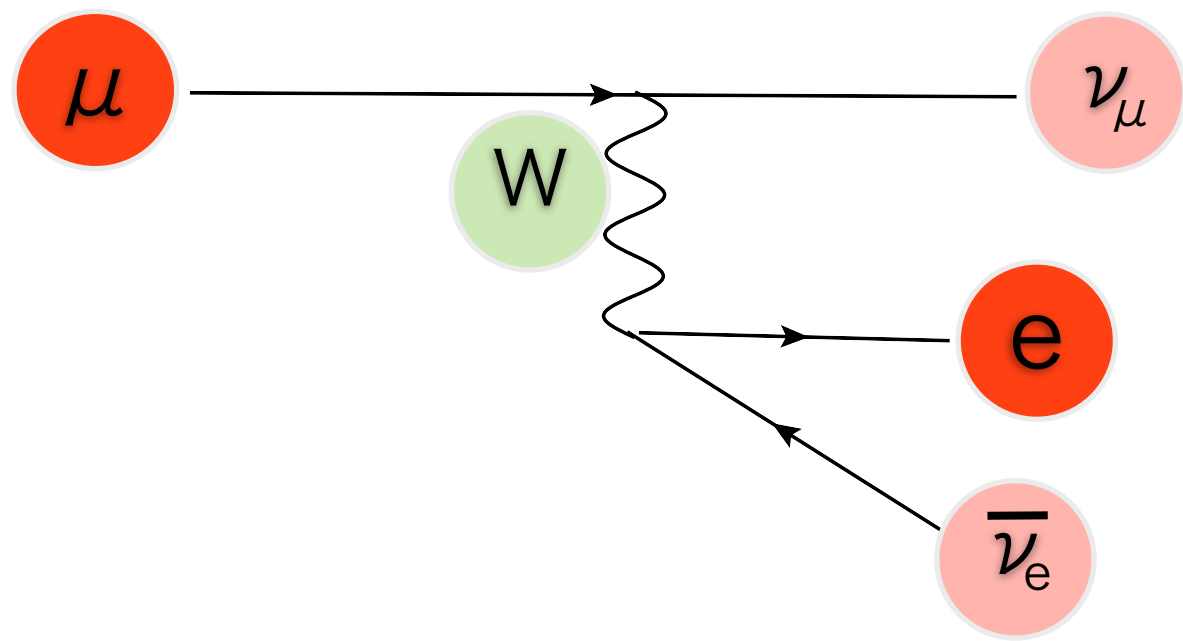


Very tiny flavor changing processes predicted
and this picture successfully describes our nature!

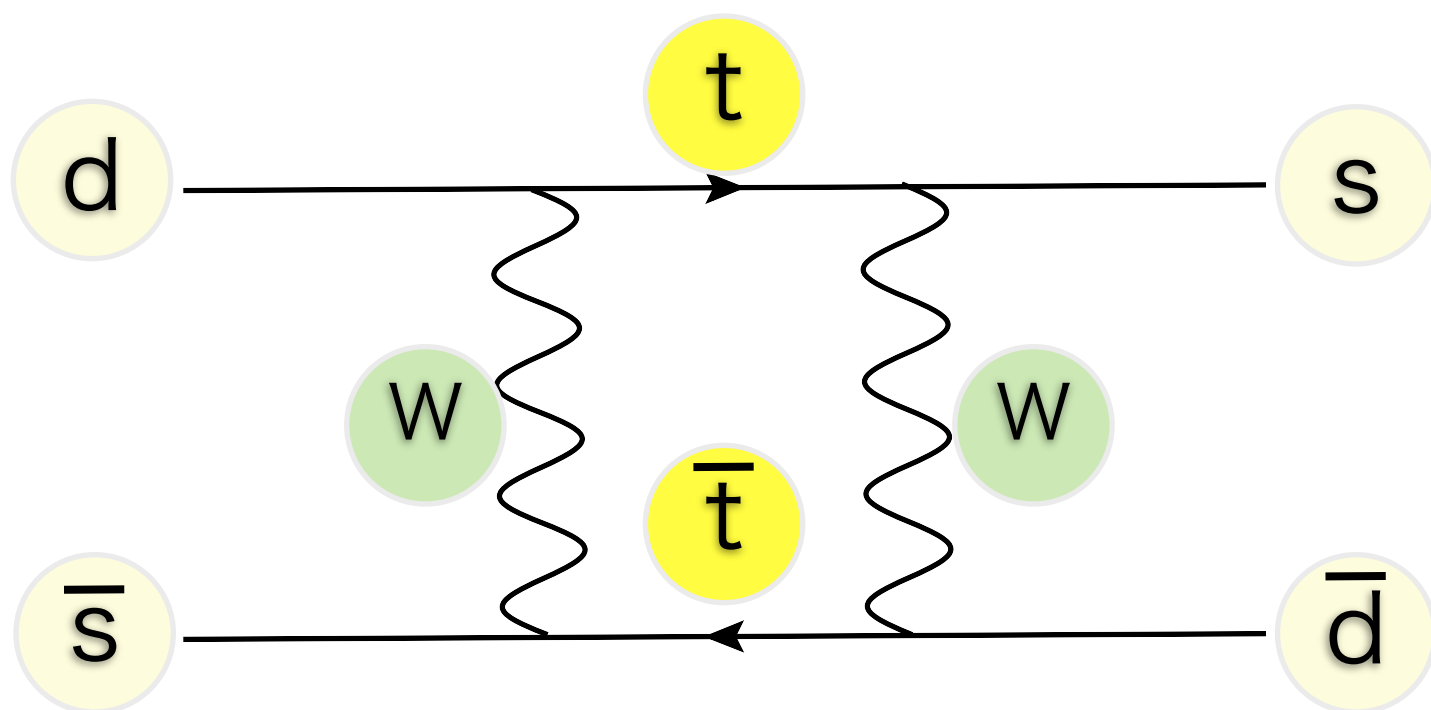


For instance, μ decays
through W exchanging

Very tiny flavor changing processes predicted
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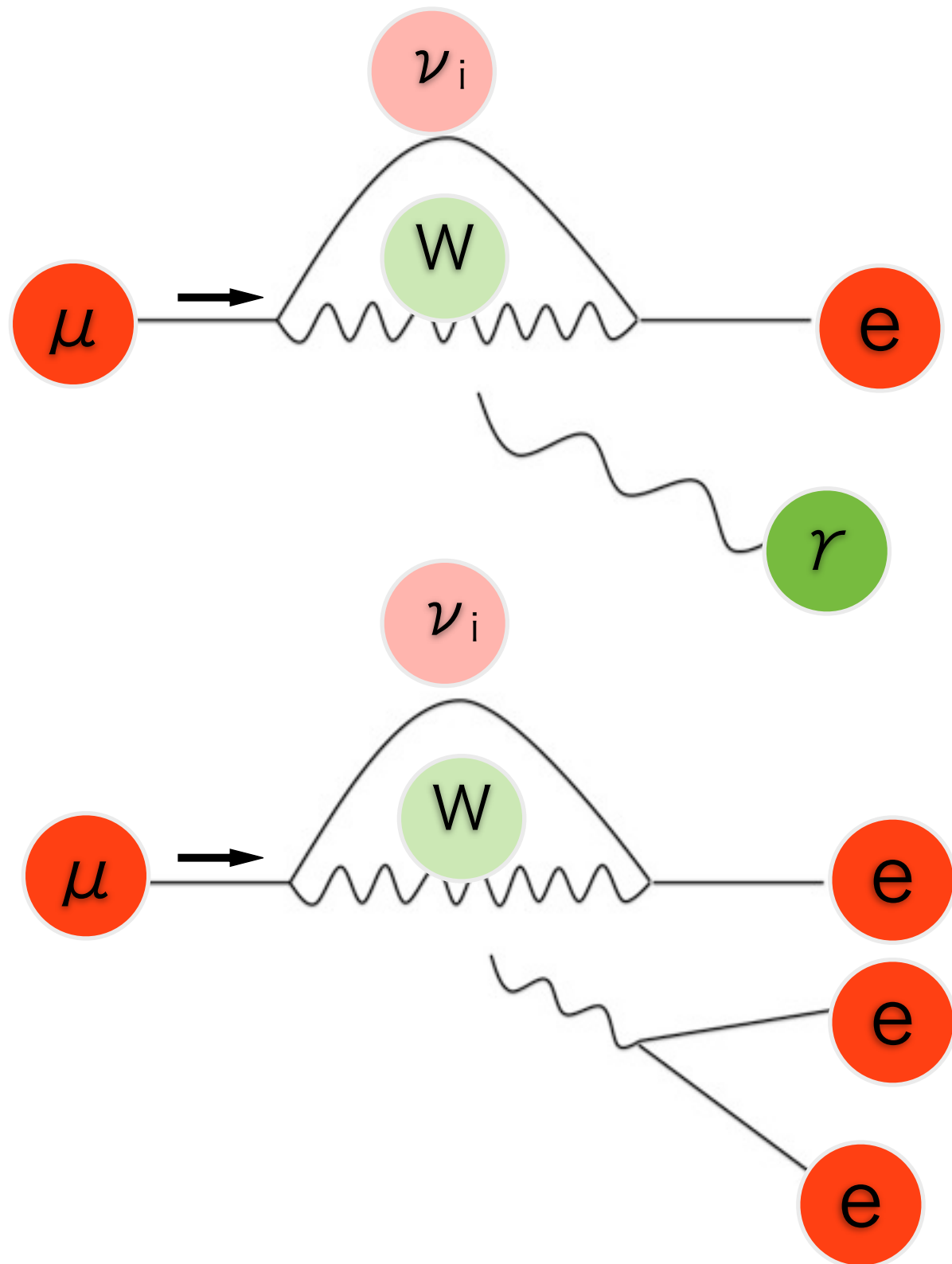


For instance, μ decays
through W exchanging



K_0 - \bar{K}_0 mixing predicted
by the one-loop,
and observed actually.

Very tiny flavor changing processes predicted
and this picture successfully describes our nature!



For instance, $\mu \rightarrow 3e$, $\mu \rightarrow e \gamma$
are very very tiny.

$$Br \propto \left| U_{\mu i}^* U_{ei} \frac{m_{\nu_i}^2}{m_W^2} \right|^2 \ll 10^{-50}$$

They are not still observed,
and consistent with the
experimental bound.

$$Br(\mu \rightarrow e\gamma) < 5.7 \times 10^{-13}$$

$$Br(\mu \rightarrow 3e) < 1.0 \times 10^{-12}$$

Is this the end of the story???

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No, No, No!

There are many “why” in the SM.

Standard Model $(SU(3)_c \times SU(2)_L \times U(1)_Y)$

is very successful in particle physics

Why?

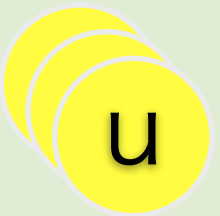
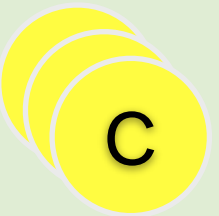
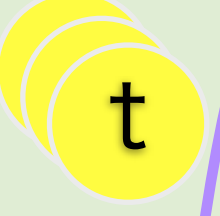
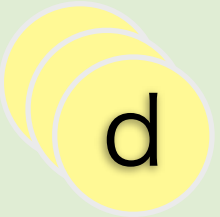
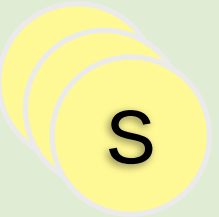
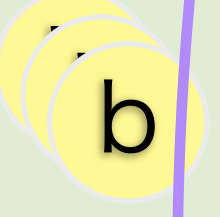
spin-1/2

Why?

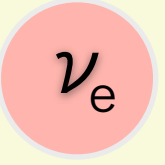
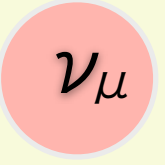
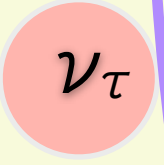

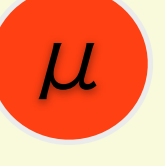
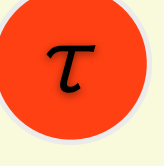
spin-1

spin-0


quarks $SU(3)_c$ -charged

			$+2/3$
			$-1/3$


leptons


			0
			-1

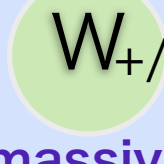
$SU(3)_c$ gauge

 gluon


$SU(2)_L \times U(1)_Y$

 Z
massive

 γ
photon

 $W_{+/-}$
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Higgs

 H

breaks $SU(2)_L \times U(1)_Y$

carry forces

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spin-1/2

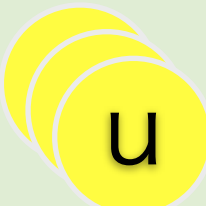
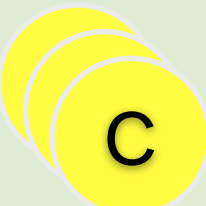
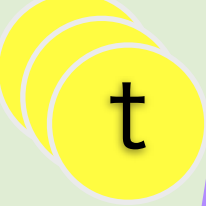
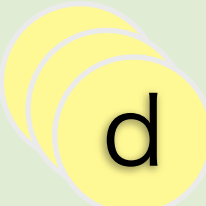
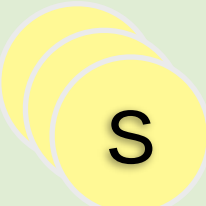
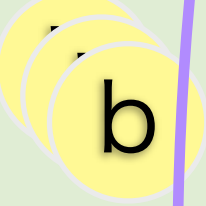
Why?

spin-1

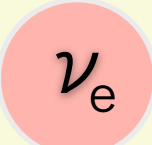

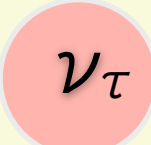



spin-0

quarks $SU(3)_c$ -charged


EM-charge

 u	 c	 t	+2/3
 d	 s	 b	-1/3


leptons


 ν_e	 ν_μ	 ν_τ	0
 e	 μ	 τ	-1

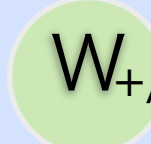
$SU(3)_c$ gauge

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Higgs

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breaks $SU(2)_L \times U(1)_Y$

Why?

carry forces

Supersymmetry (SUSY)

Grand Unified Theory (GUT)

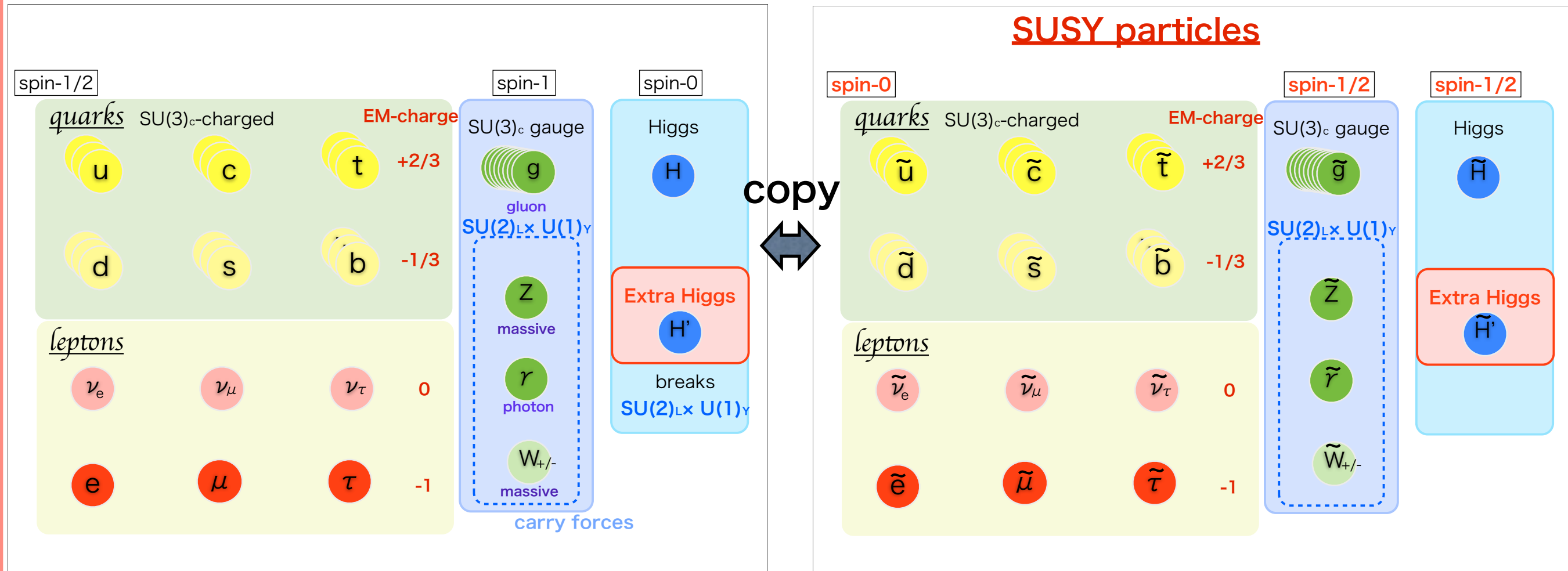
are very natural explanations!

Answer to the origin of EW scale:

Supersymmetry (SUSY)

SUSY can explain why $SU(2)_L \times U(1)_Y$ breaking happen around 200 GeV

Supersymmetric SM



because of no quadratic divergence
 nonrenormalization theorem
 radiative $SU(2)_L \times U(1)_Y$ breaking

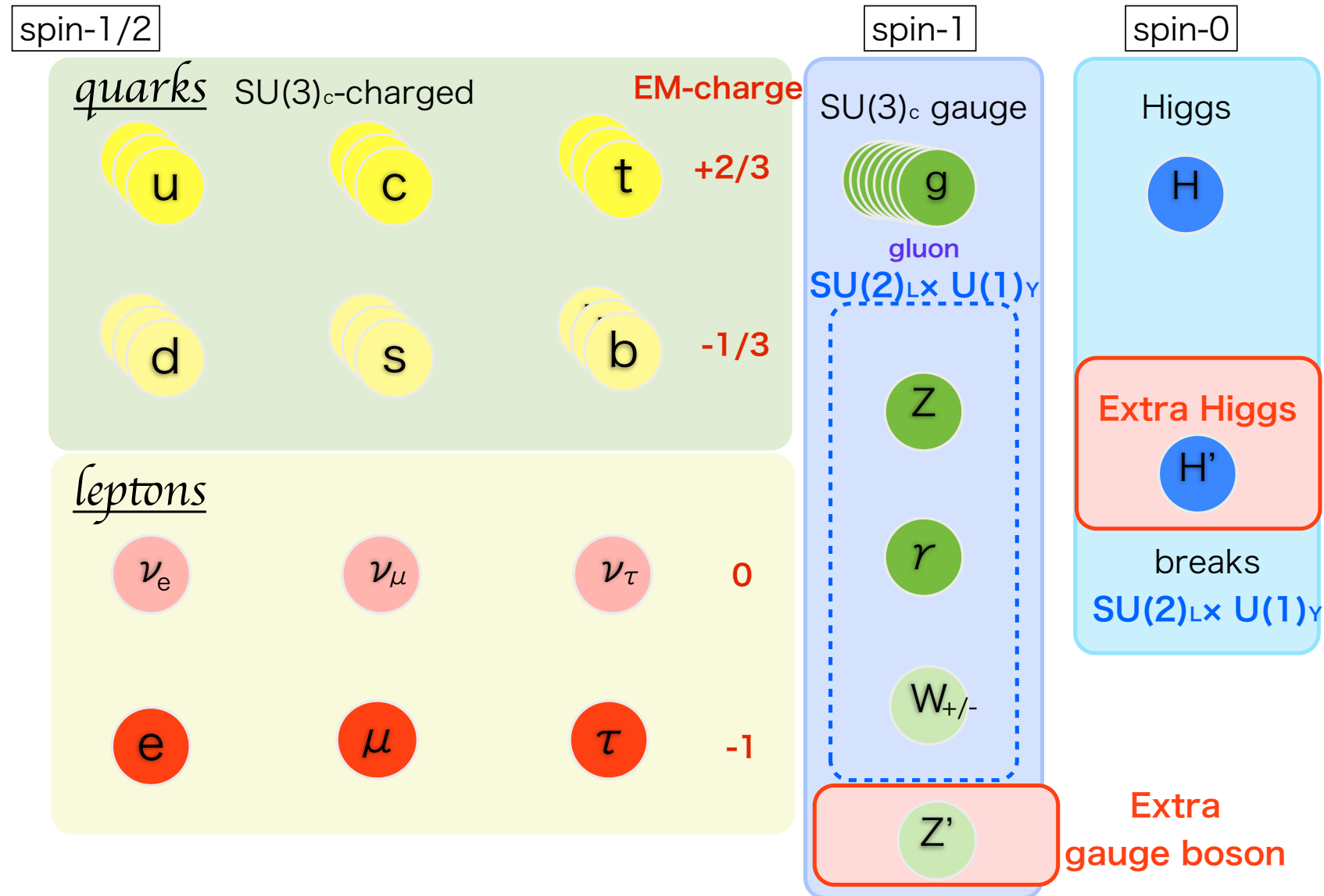
SUSY scale relates to EW scale

Answer to the origin of gauge symmetry:

Grand Unified Theory (GUT)

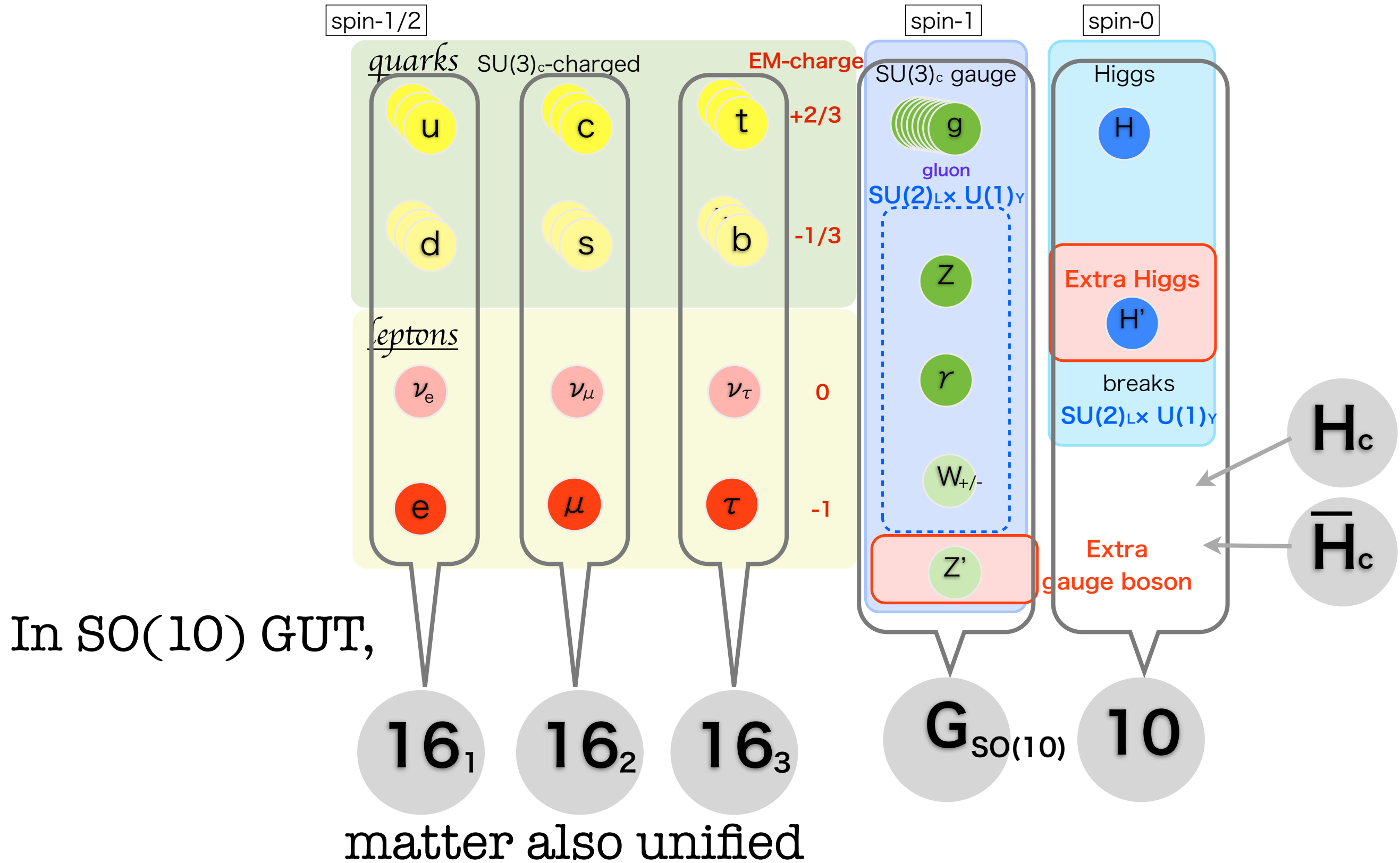
SM gauge groups naturally embedded into **GUT**

SO(10) Embedding: $SU(3)_c \times SU(2)_L \times U(1)_Y \times U(1)_X \rightarrow SO(10)$
 slightly extended SM extra



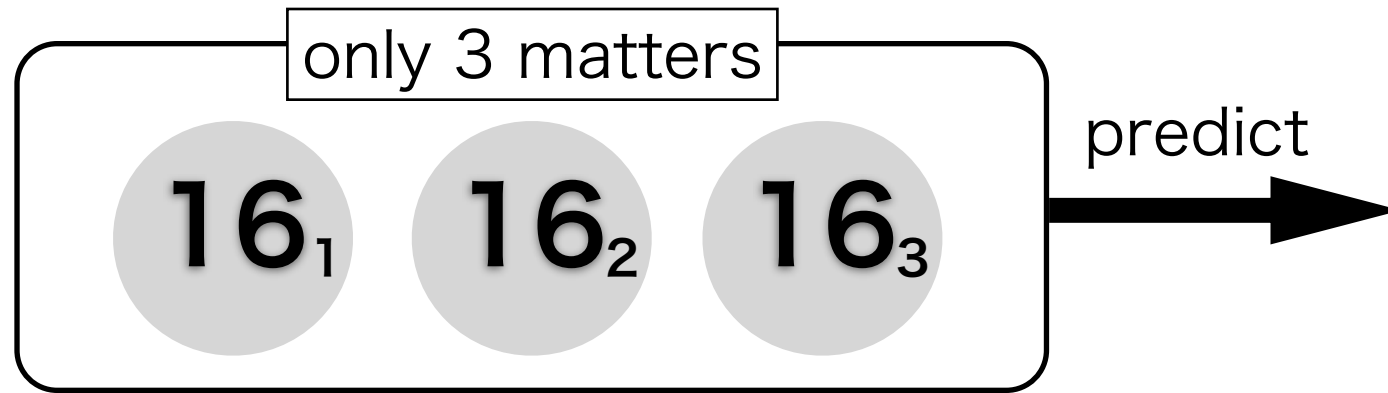
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supersymmetric $SO(10)$ GUT looks very elegant and natural
but story is not so simple...

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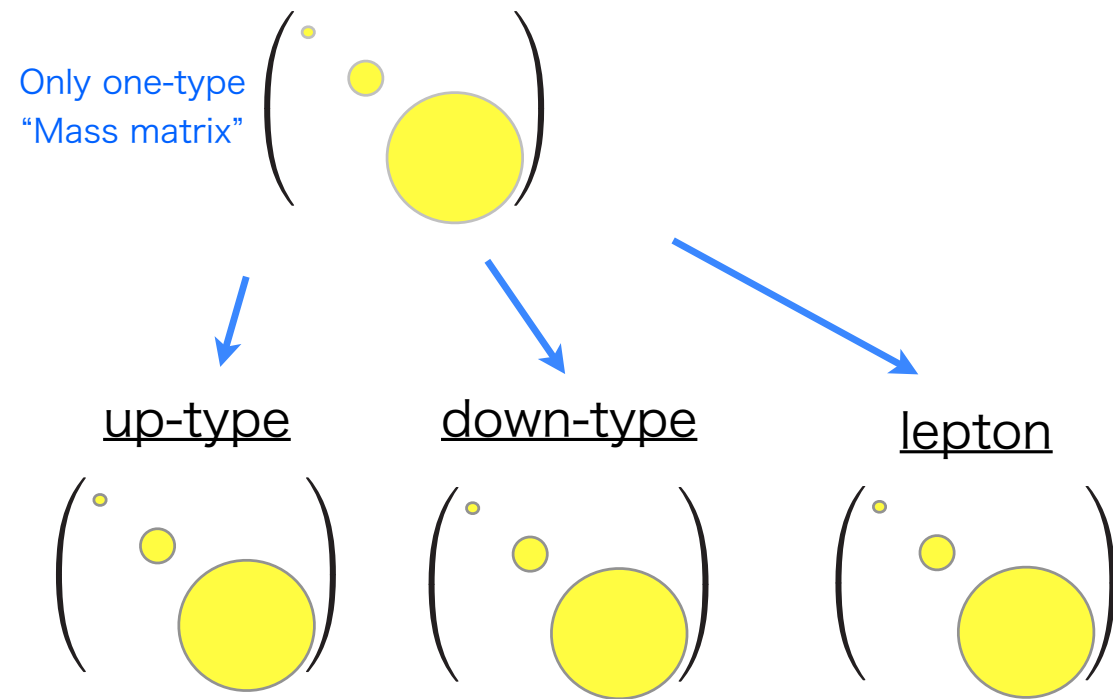


same mass matrices
 of (u,c,t), (d,s,b), (e, μ , τ)
 no CKM mixing

SO(10) GUT unifies quarks and leptons into **16** rep.

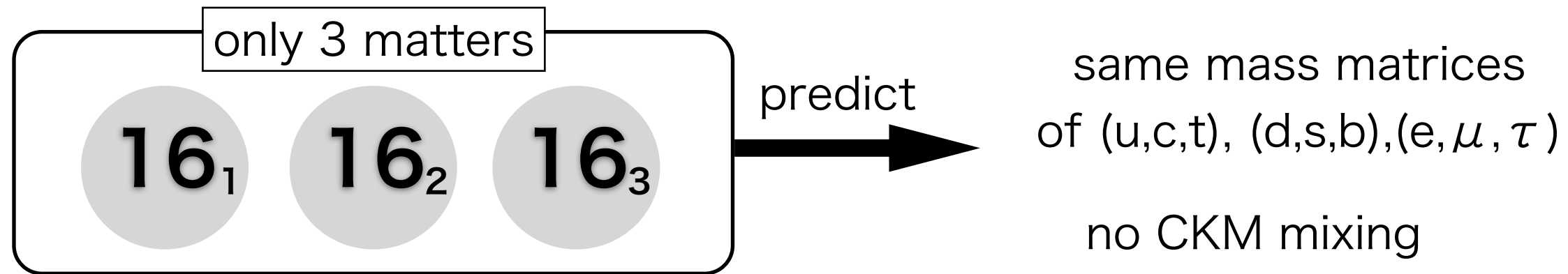
$$W_{\min} = h_{ij} 16_i 16_j 10_H$$

$$@ \tan \beta = 1$$



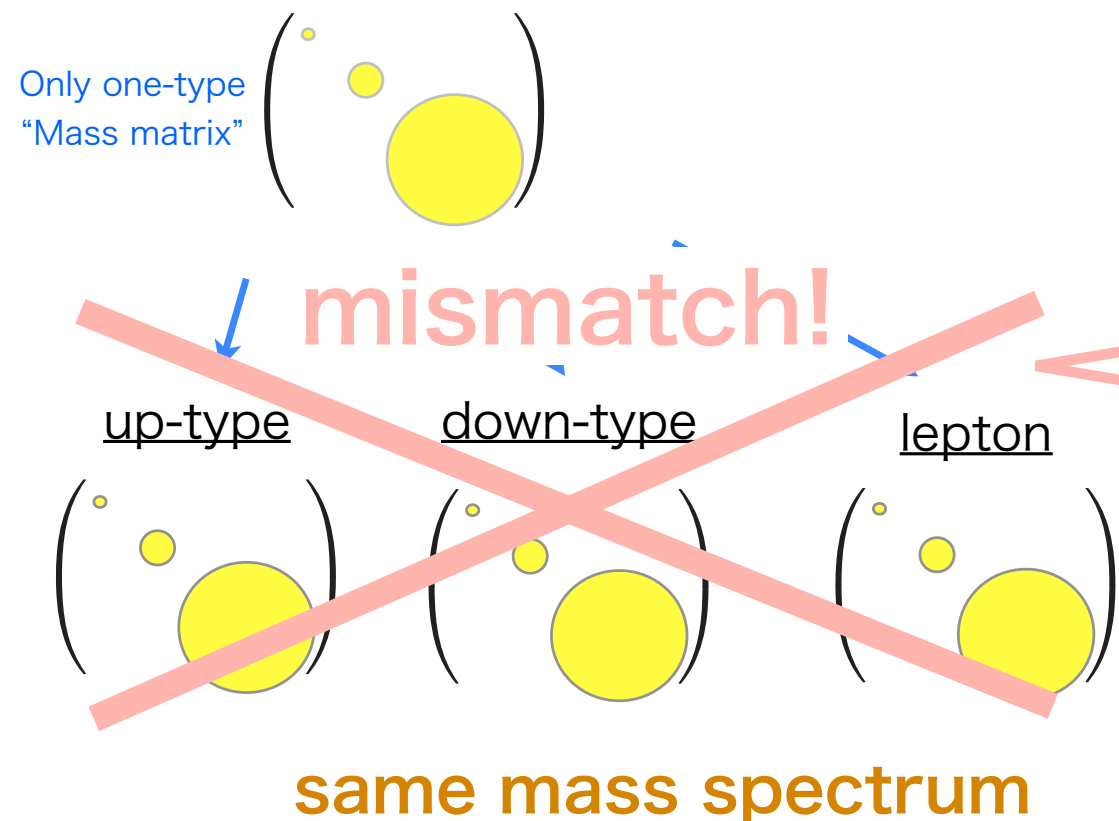
same mass spectrum

supersymmetric SO(10) GUT looks very elegant and natural
but story is not so simple...



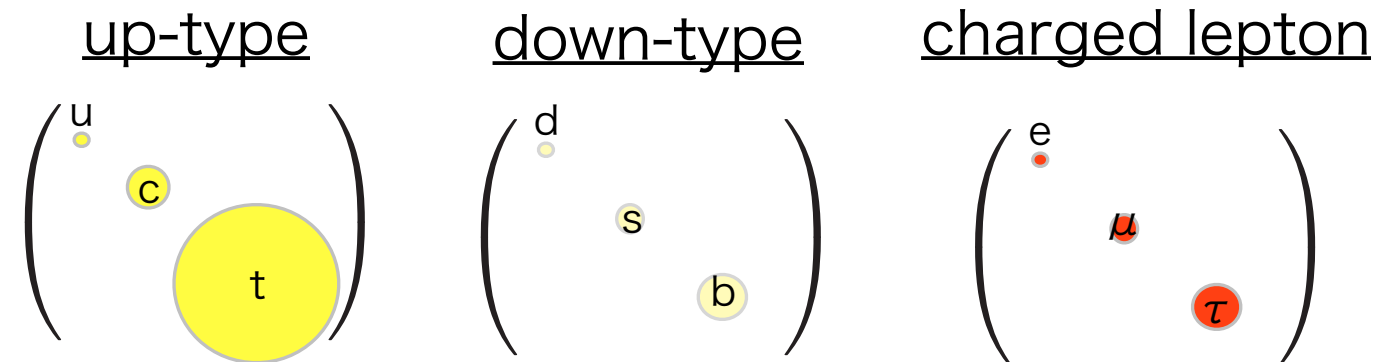
SO(10) GUT unifies quarks and leptons into

$$W_{\min} = h_{ij} 16_i 16_j 10_H \quad @ \tan \beta = 1$$



Realistic Mass Forms

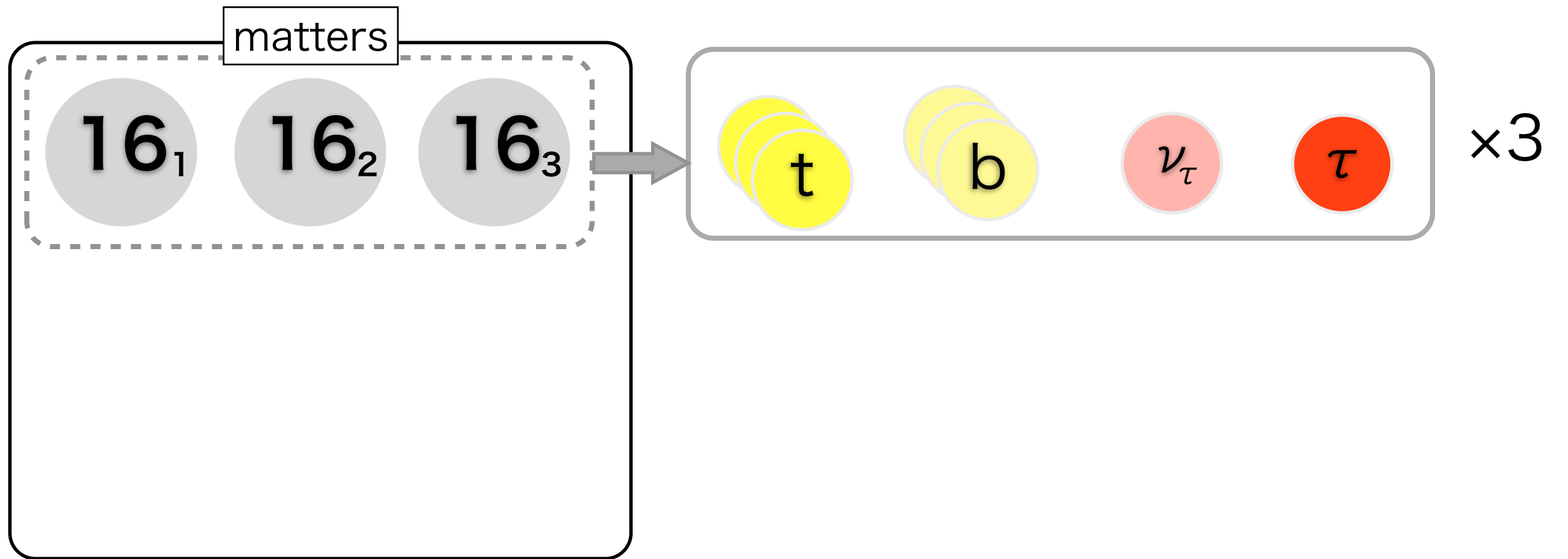
hierarchical masses



CKM mixing from the different mass matrices

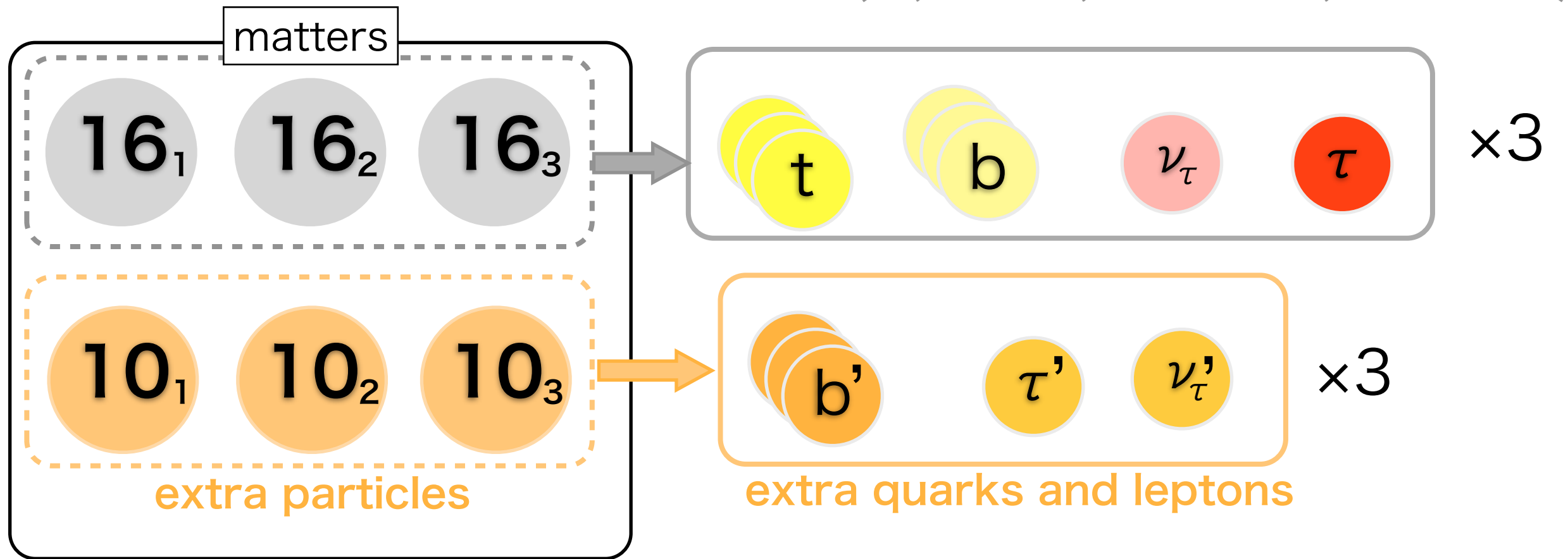
We need extra something.

minimal setup



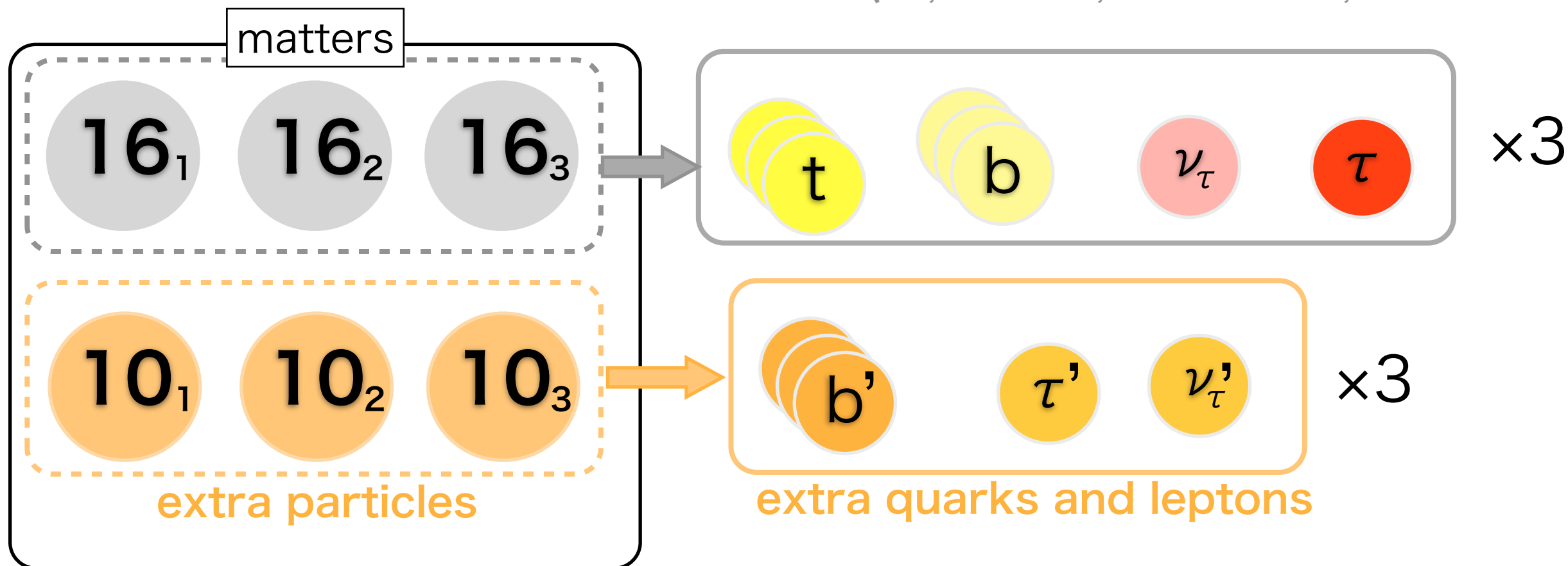
Let me add extra matters ($\mathbf{10}_1, \mathbf{10}_2, \mathbf{10}_3$)

(YO, J. Hisano, Y. Muramatsu, M. Yamanaka)



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Then SM particles are given by the linear combinations:

ex)

minimal

extended

$$|(b_R^c)_{SM}\rangle = |b_R^c\rangle \quad \Rightarrow \quad |(b_R^c)_{SM}\rangle = (\cos \theta_b) |b_R^c\rangle + (\sin \theta_b) |b'^c_R\rangle$$

realization of b mass by the mixing

minimal

extended

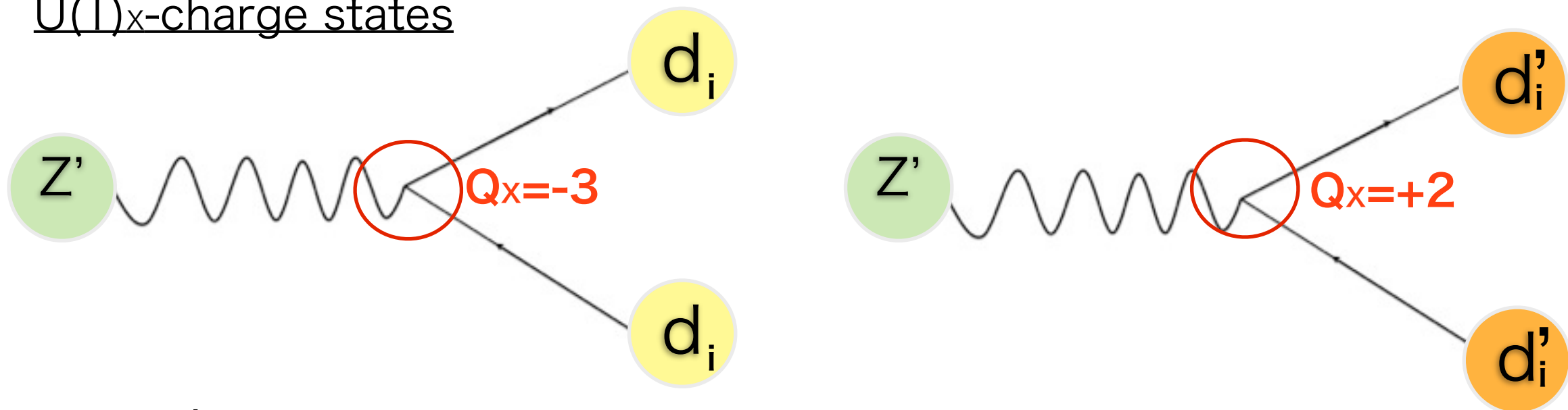
$$m_b = m_t \quad \Rightarrow \quad m_b = \begin{pmatrix} m_t \end{pmatrix} \times \begin{pmatrix} \cos \theta_b \end{pmatrix}$$

@ $\tan \beta = 1$

Crucial fact is

d_i and d'_i carry different $U(1)_X$ charges
extra

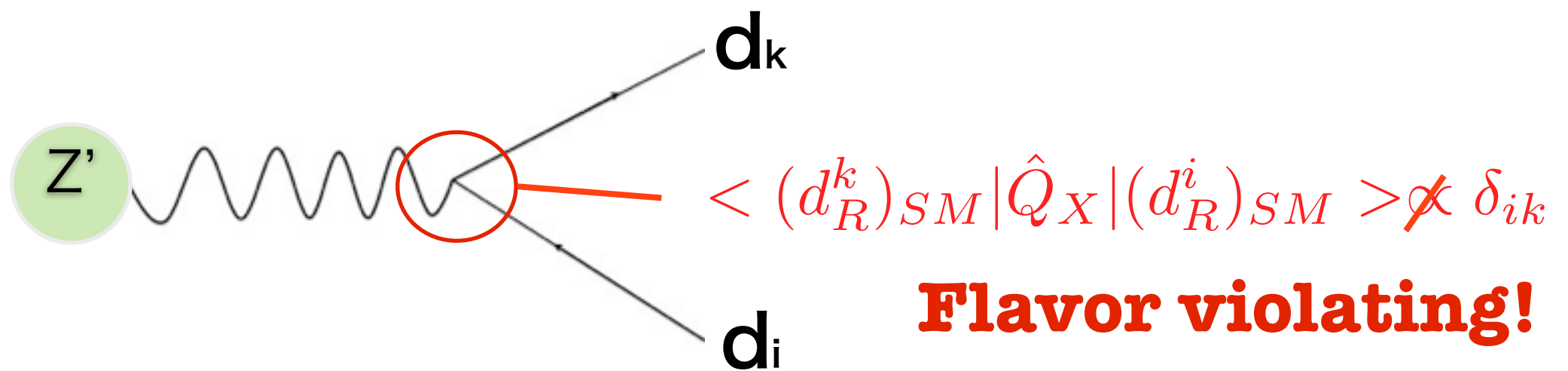
$U(1)_X$ -charge states



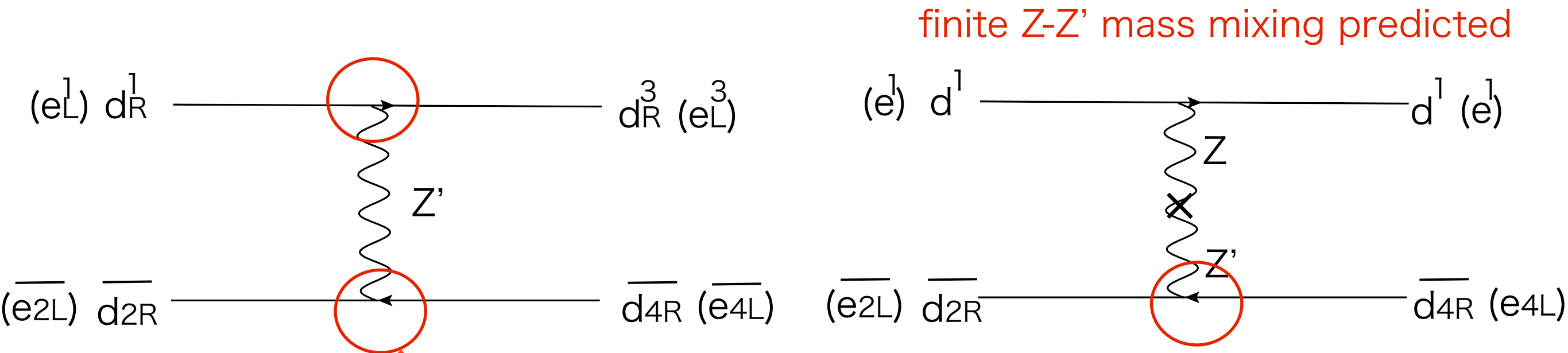
mass eigenstates

down-type quarks

$$|(d_R^i)_{SM}\rangle = U_{ij} |d_R^j(-3)\rangle + U'_{ij} |d_R^j(+2)\rangle$$



Contributions to Flavor Physics



finite Z - Z' mass mixing predicted

flavor changing
almost flavor-universally

$(d,s), (d,b), (s,b), (e,\mu), (e,\tau), (\mu,\tau)$ elements are $O(1)$ and complex.

Very nice predictions

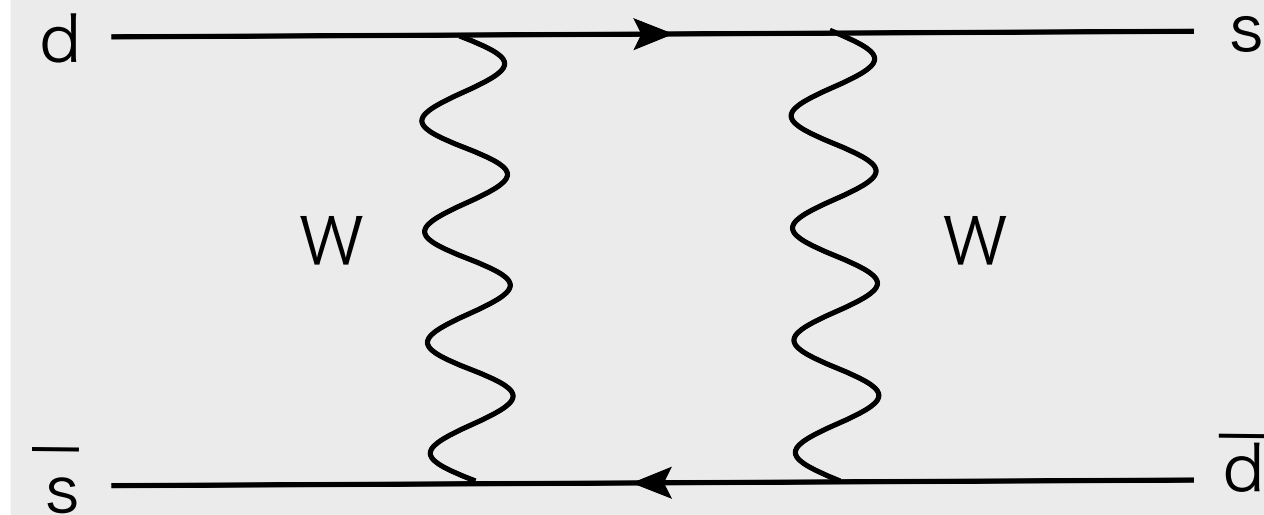
because it is usually difficult to prove GUTs.

Z' scale may be also determined by the 125 GeV Higgs and Z' coupling is predicted by the GUT relation.

Predictability is very high!

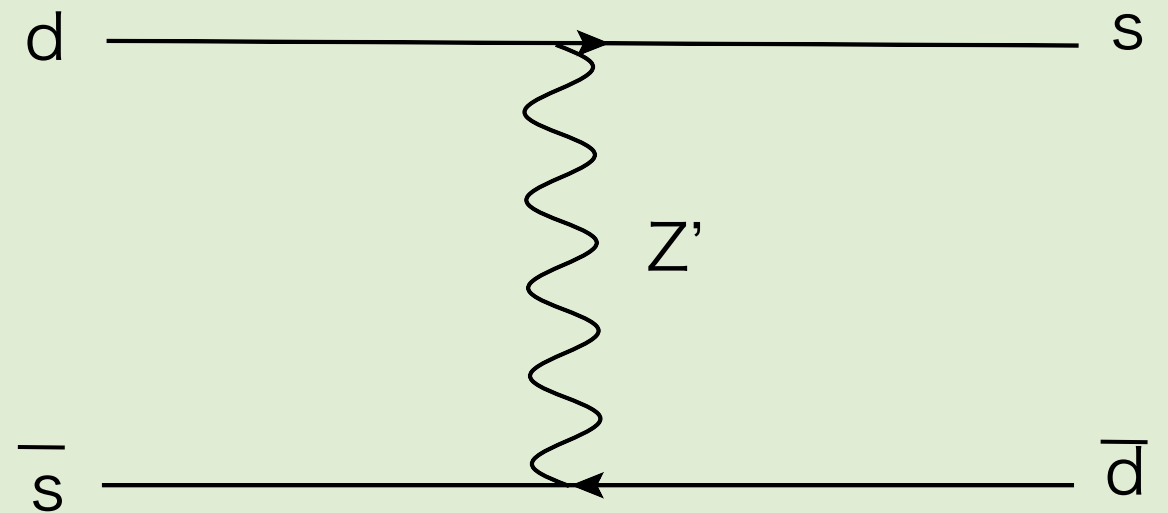
Important processes are K system and μ physics

SM prediction



vs

Z' contribution



$$\underbrace{|V_{ts}^* V_{td}|}_{K \text{ system}} \sim 5 \cdot 10^{-4} \ll \underbrace{|V_{tb}^* V_{td}|}_{B_d \text{ system}} \sim 10^{-2} < \underbrace{|V_{tb}^* V_{ts}|}_{B_s \text{ system}} \sim 4 \cdot 10^{-2}$$

K system is more sensitive to new physics

Experimental constraints on μ are stronger

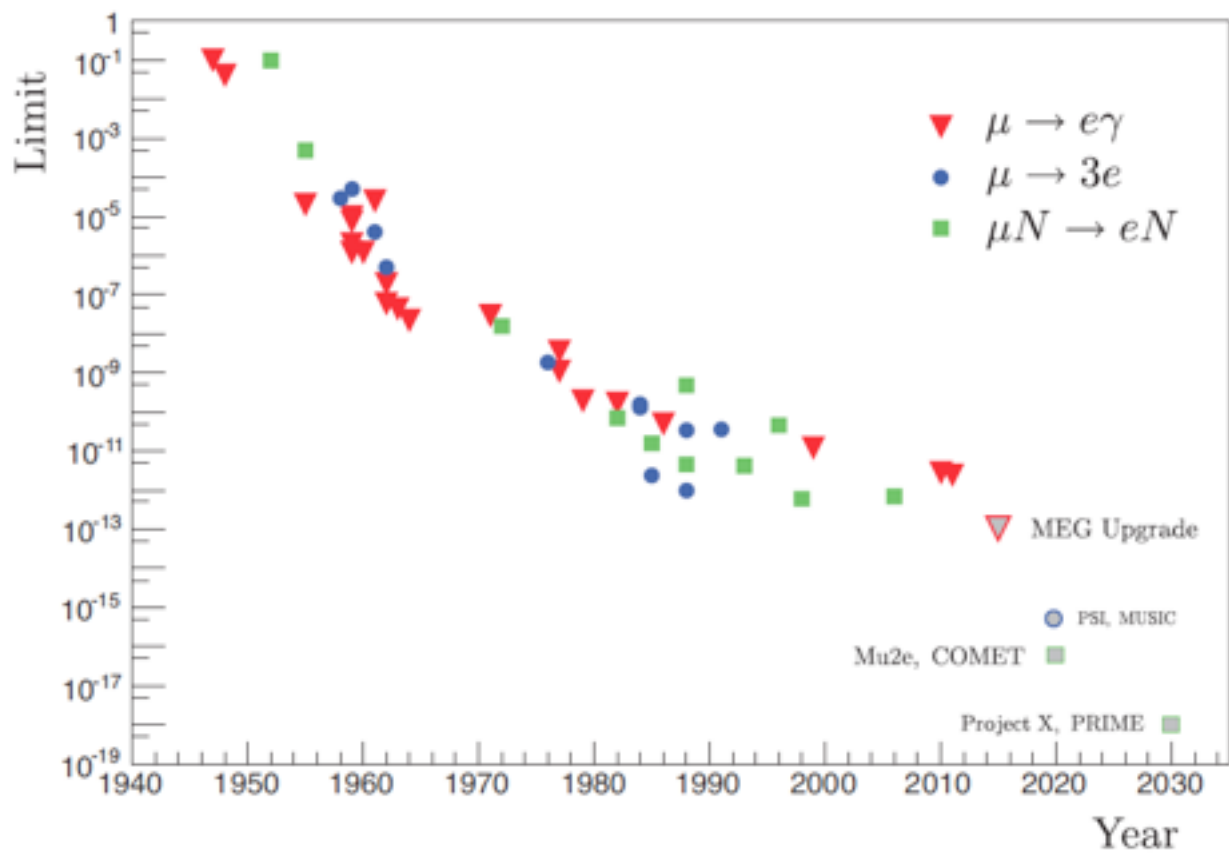
τ decay

$$\text{BR}(\tau \rightarrow l_i l_j l_k) \lesssim 10^{-8}$$

$$\text{BR}(\tau \rightarrow l_i \gamma) \lesssim 10^{-8}$$

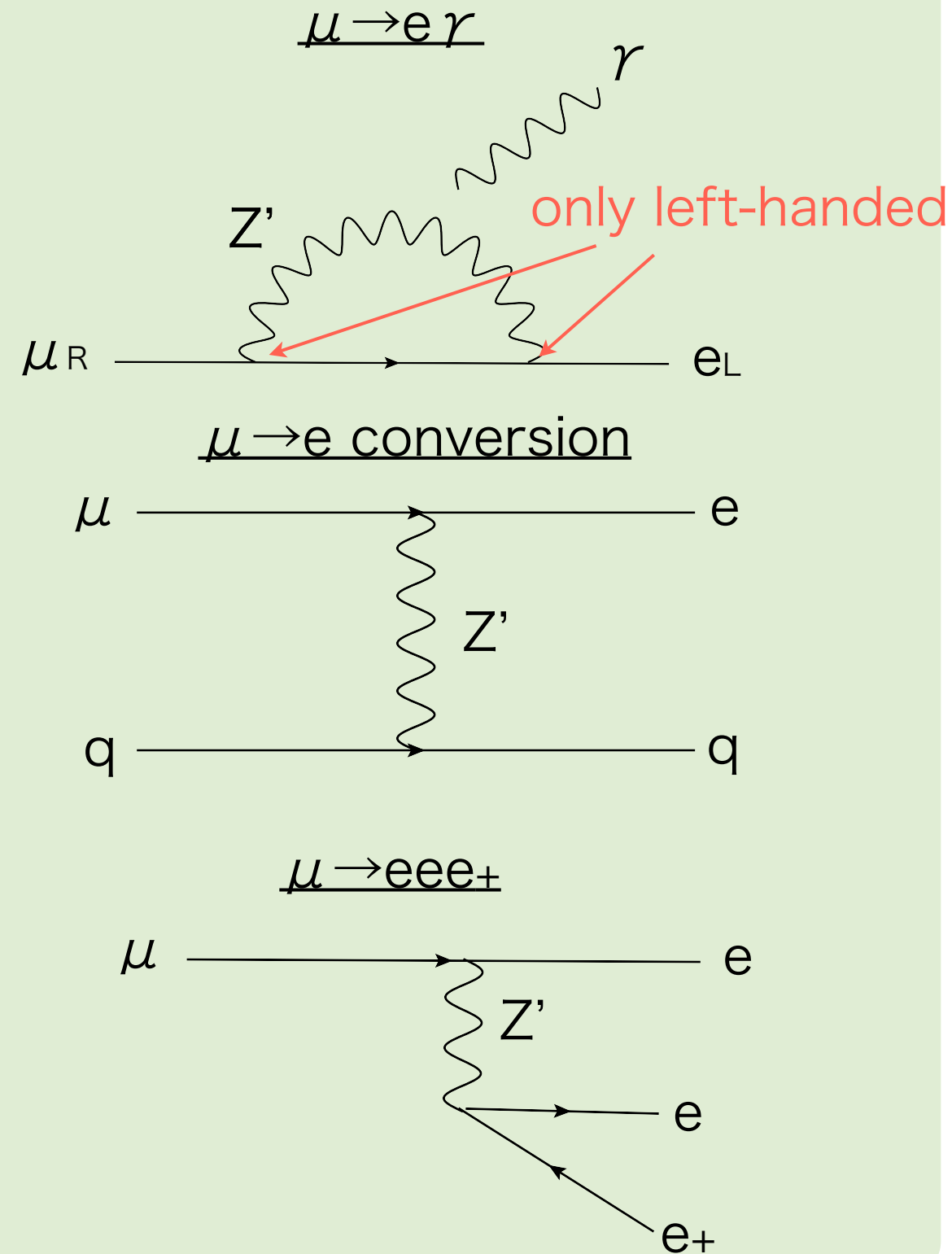
(arXiv:1001.3221;0908.2381)

μ decay



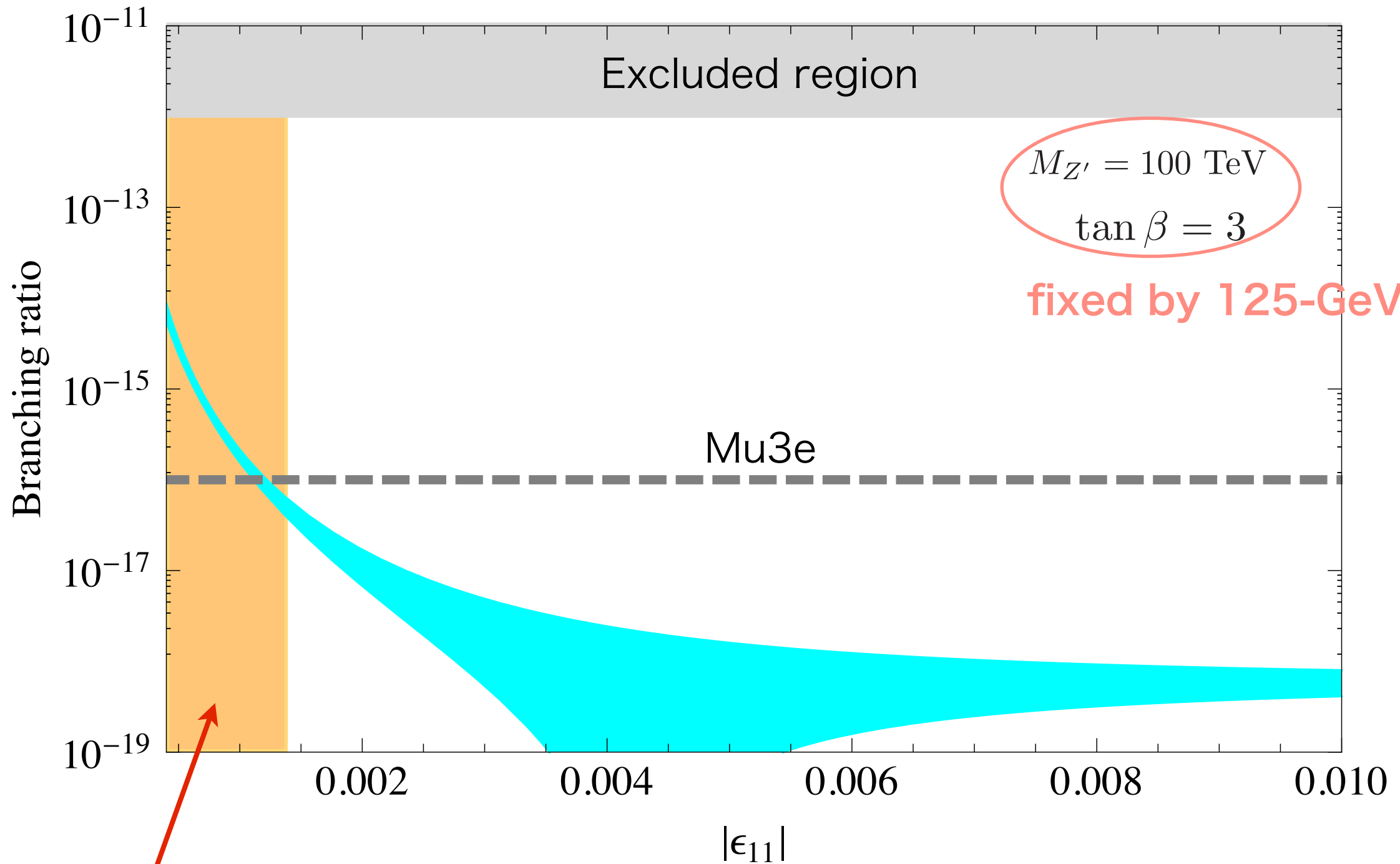
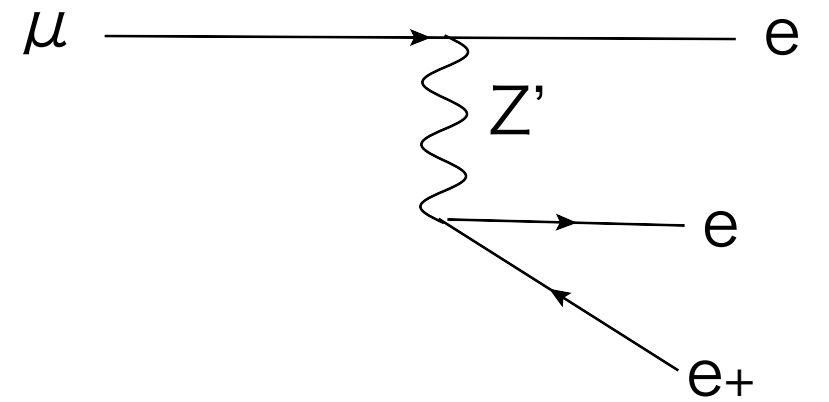
(arXiv:1307.5787)

VS



$\mu \rightarrow 3e$, μ -e conversion are most important

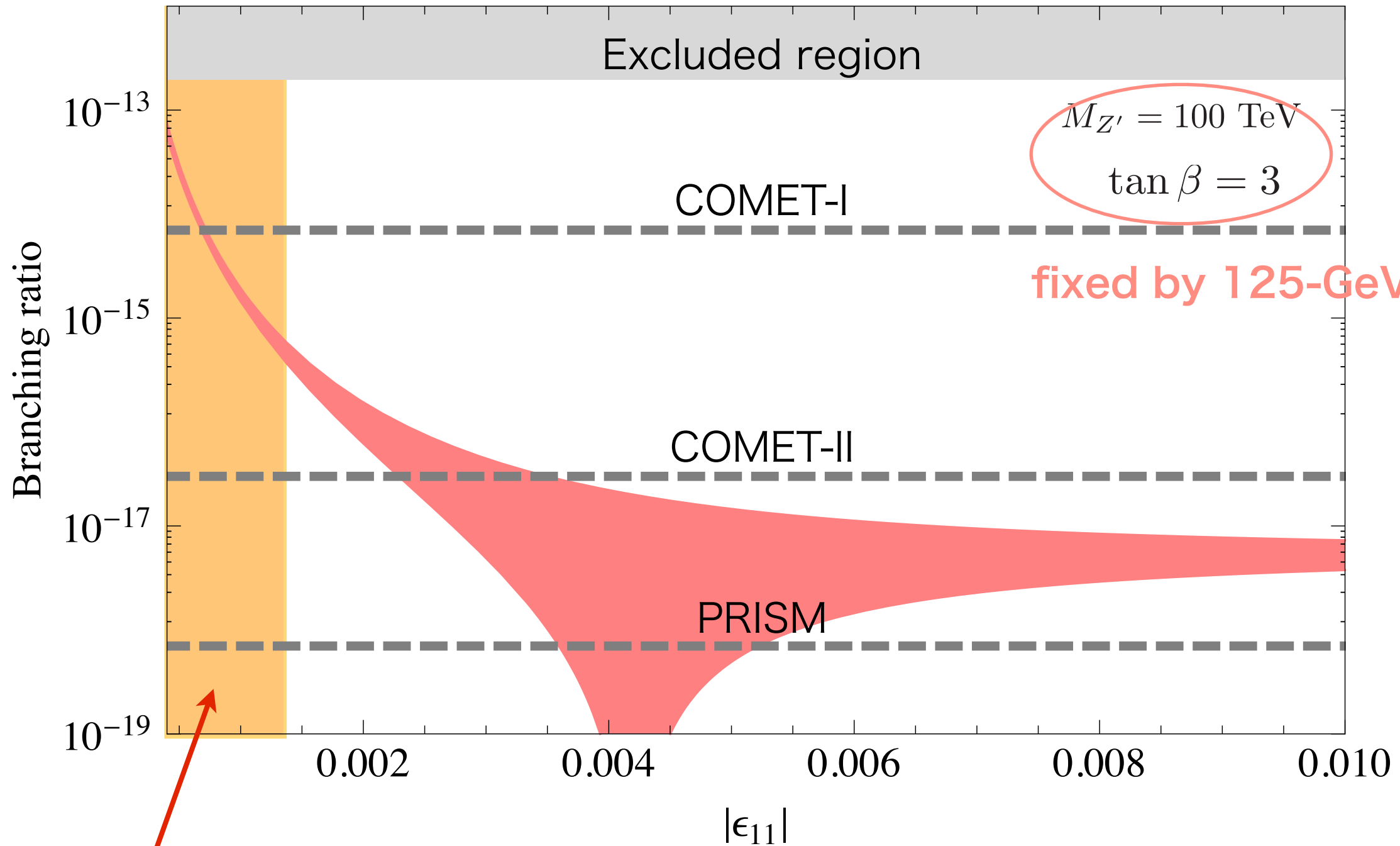
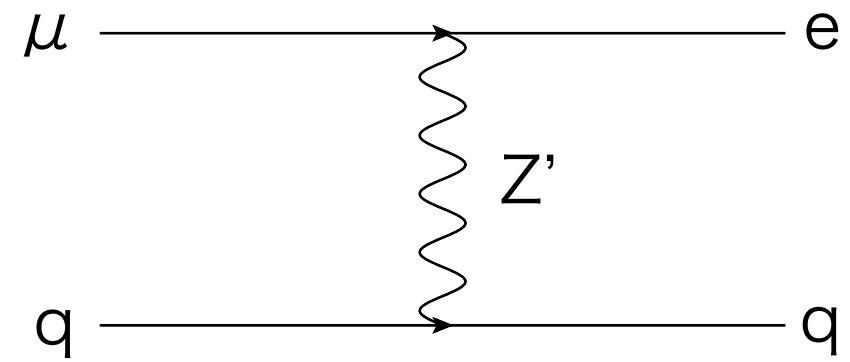
Allowed region for $\mu \rightarrow 3e$



Excluded by K-Kbar mixing

deviations of $K_L \rightarrow \mu \mu, \mu e, K_+ \rightarrow \pi_+ \nu \nu$ are less than O(1)%.

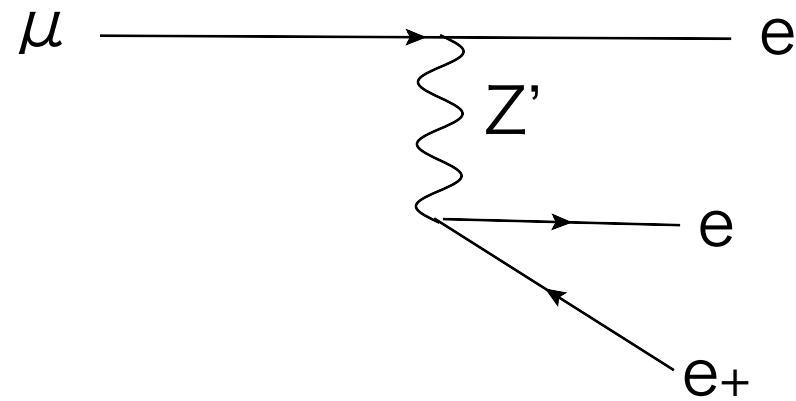
Allowed region for $\mu \rightarrow e$ in Al



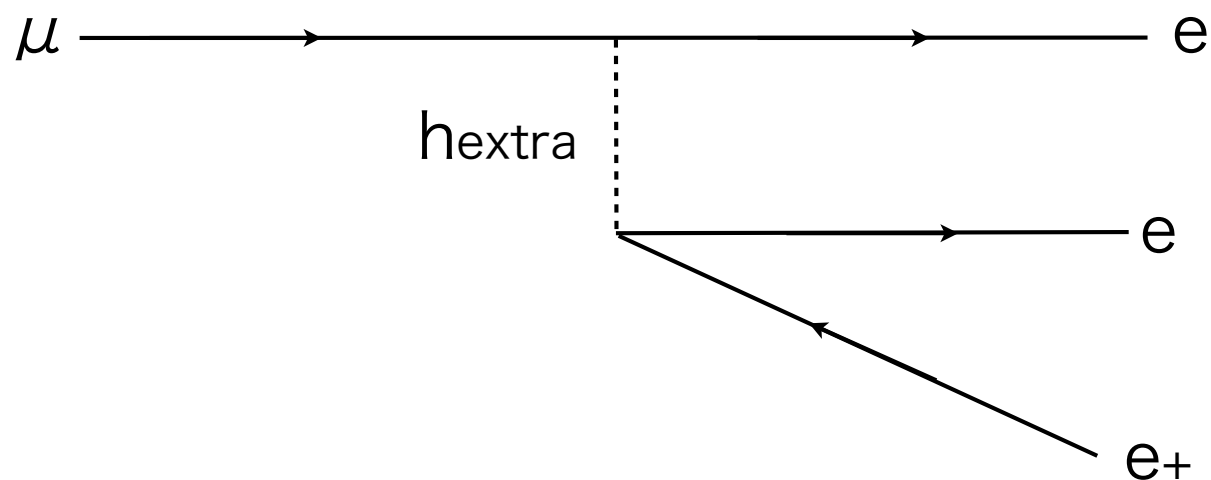
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I introduced Z' exchanging flavor violation



extra scalar (spin-0) exchanging flavor violation is also widely discussed in the BSMs.



Higgs mediated Flavor violation

Higgs mediation given by Yukawa couplings

Extra Higgs



$$Y_{ij} h_{\text{extra}} \bar{l}_L^i l_R^j$$

$$(Y_{ij}) = \begin{pmatrix} Y_{ee} & Y_{e\mu} & Y_{e\tau} \\ Y_{\mu e} & Y_{\mu\mu} & Y_{\mu\tau} \\ Y_{\tau e} & Y_{\tau\mu} & Y_{\tau\tau} \end{pmatrix}.$$

tend to be large in many BSMs,
because of the mass hierarchy

The processes involving third generations (τ , b , t) are important

LFVs involving electron suppressed because of small Y_{ee} , $Y_{\mu e}$ and $Y_{e\mu}$.

1 example of the BSMs

which predict the Higgs-mediated Flavor violation.

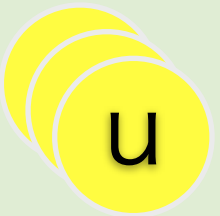
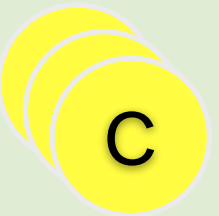
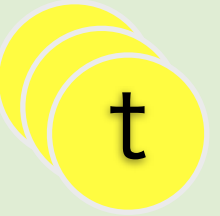
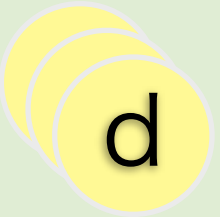
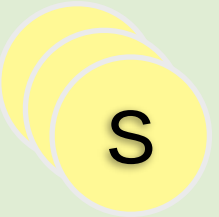
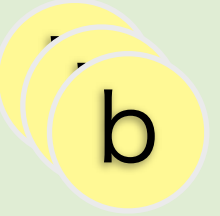
Standard Model $(SU(3)_c \times SU(2)_L \times U(1)_Y)$

is very successful in particle physics

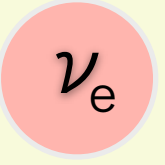
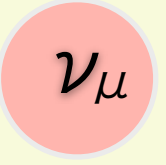
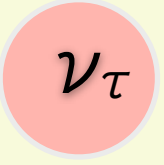

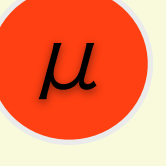
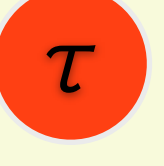
spin-1/2

quarks $SU(3)_c$ -charged

EM-charge


 u	 c	 t	+2/3
 d	 s	 b	-1/3

leptons


 ν_e	 ν_μ	 ν_τ	0
 e	 μ	 τ	-1


spin-1

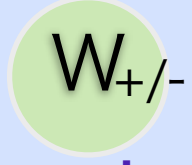
$SU(3)_c$ gauge

 g
gluon

$SU(2)_L \times U(1)_Y$

 Z
massive


 γ
photon

 $W_{+/-}$
massive

carry forces

spin-0

Higgs

 H

breaks $SU(2)_L \times U(1)_Y$

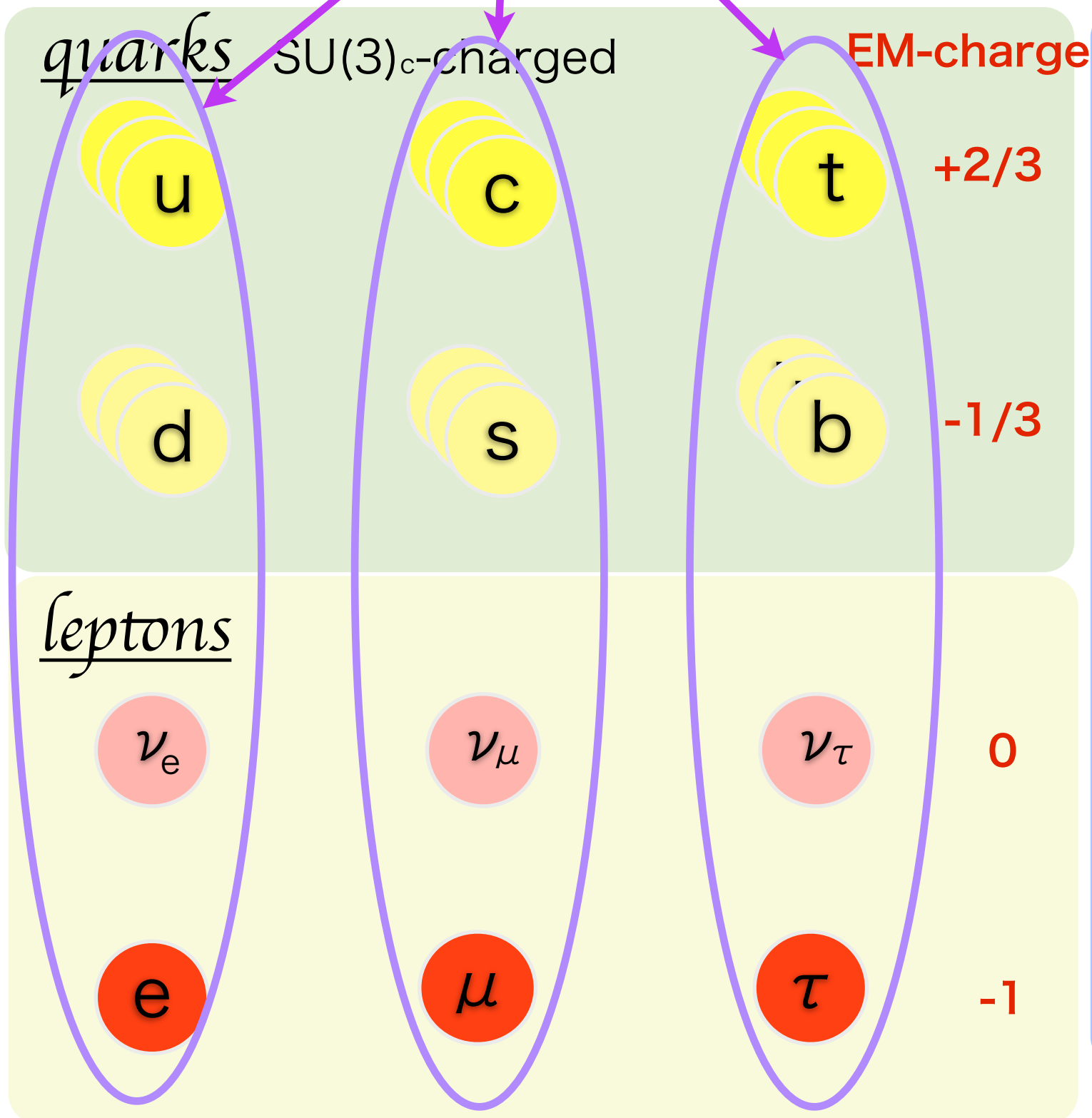
Why are there 3?

Are there flavor symmetry?

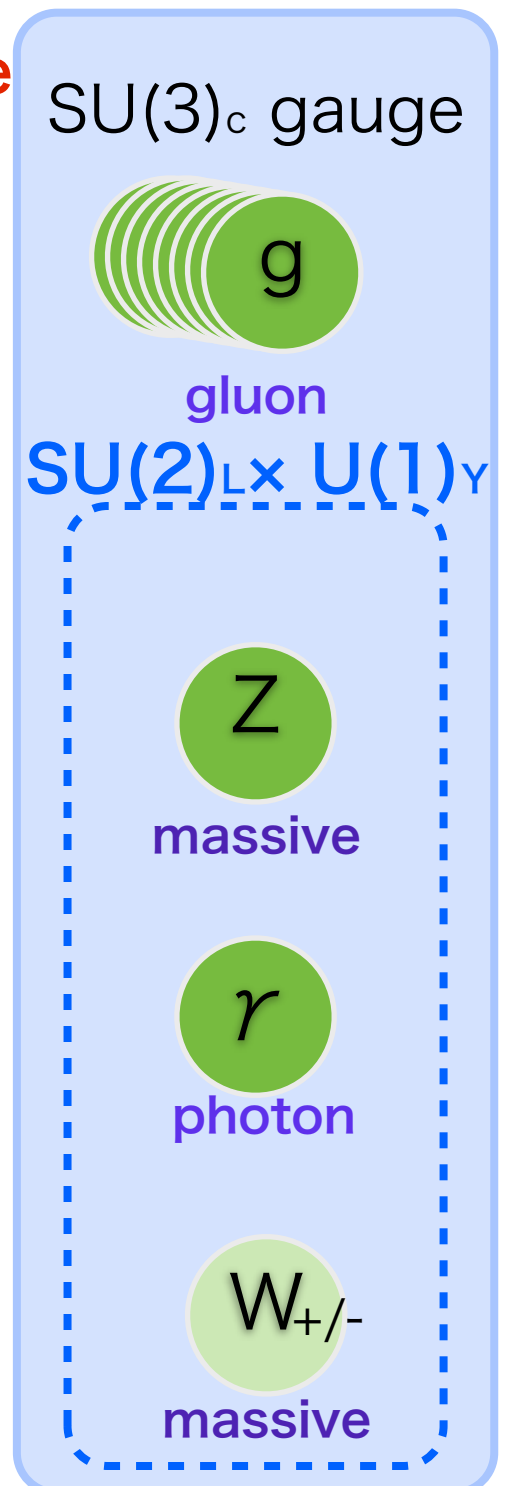
$$SU(2)_L \times U(1)_Y$$

particle physics

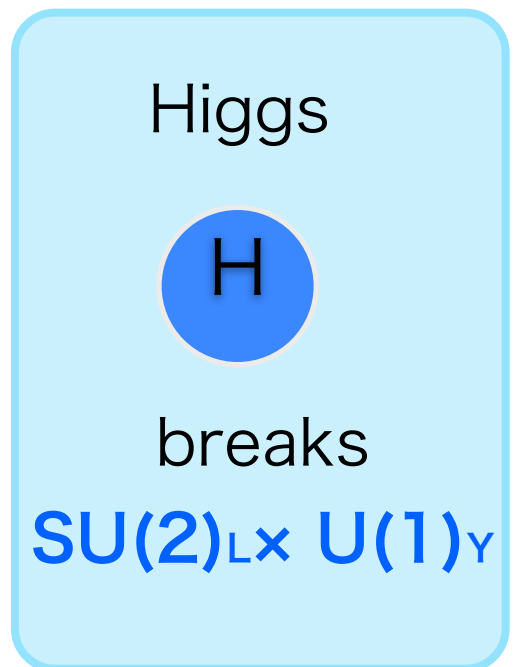
spin-1/2



spin-1



spin-0



Why are there 3?

Are there flavor symmetry?

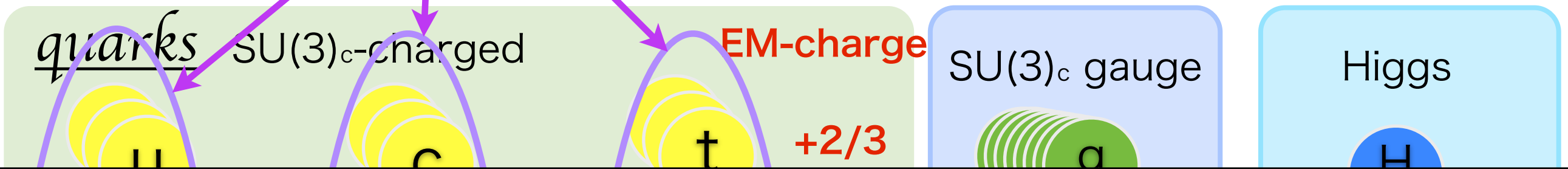
$$SU(2)_L \times U(1)_Y$$

particle physics

spin-1/2

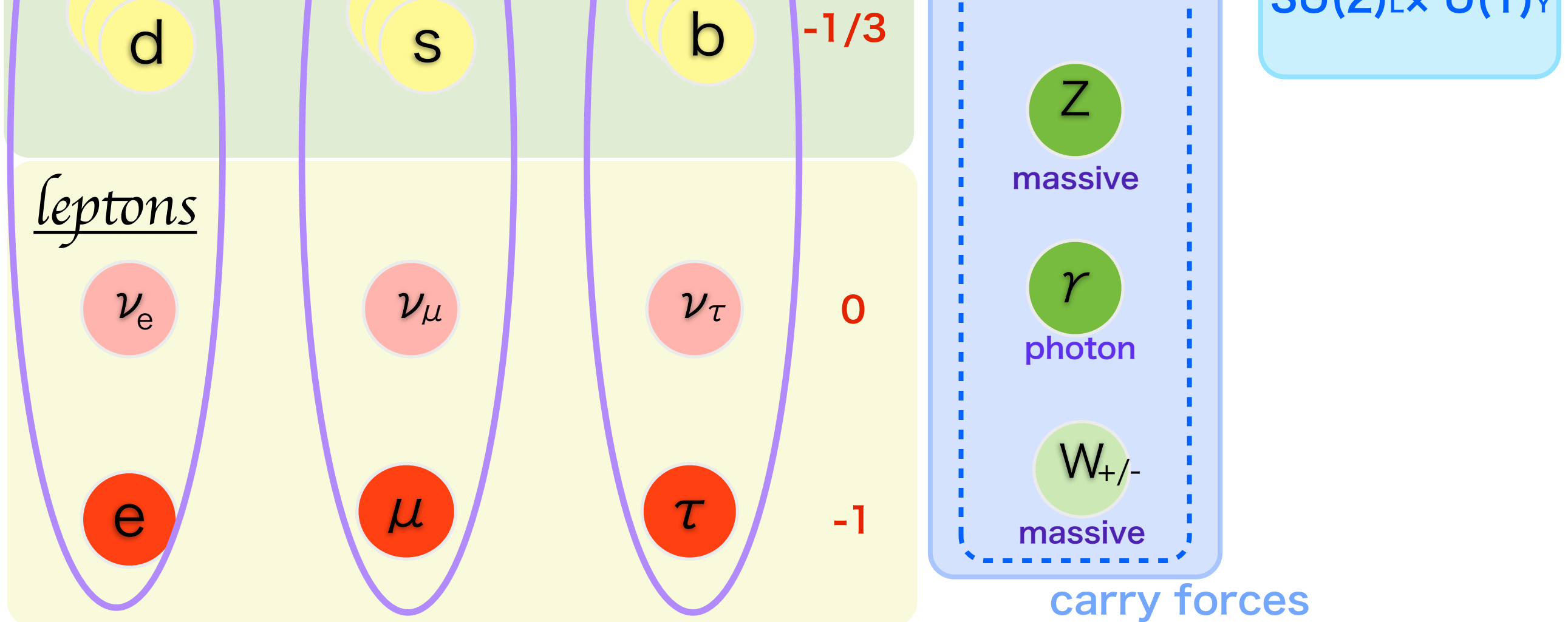
spin-1

spin-0

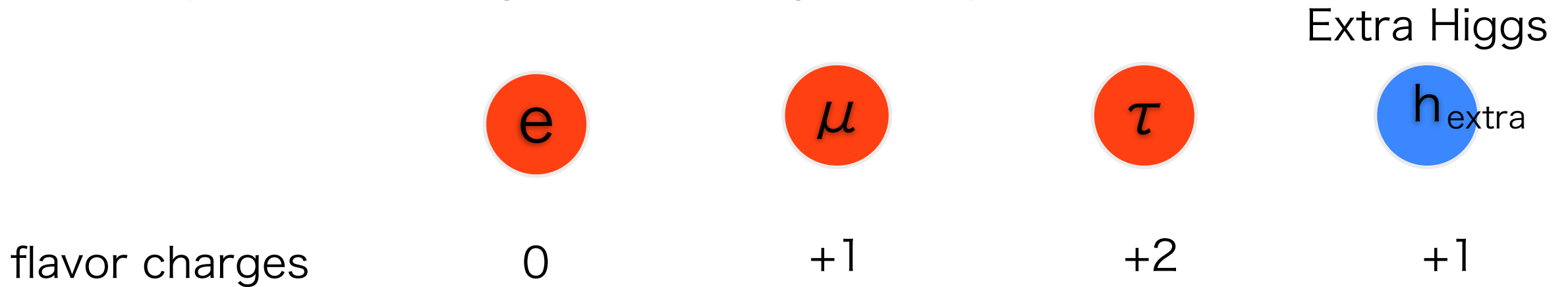


Let's assume

3 Generations carry different quantum numbers

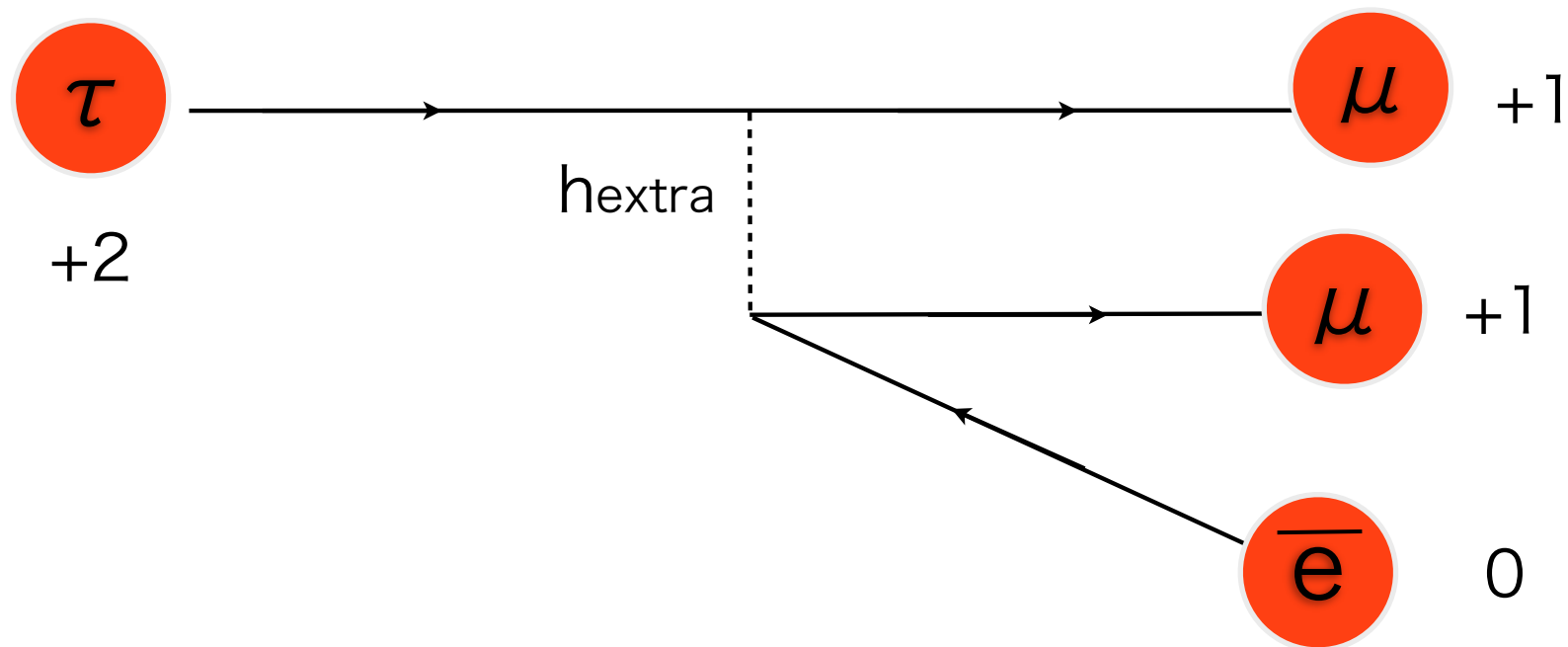


For example, let me assign flavor charges to leptons



(In fact, many flavor models suggest this feature.)

prediction

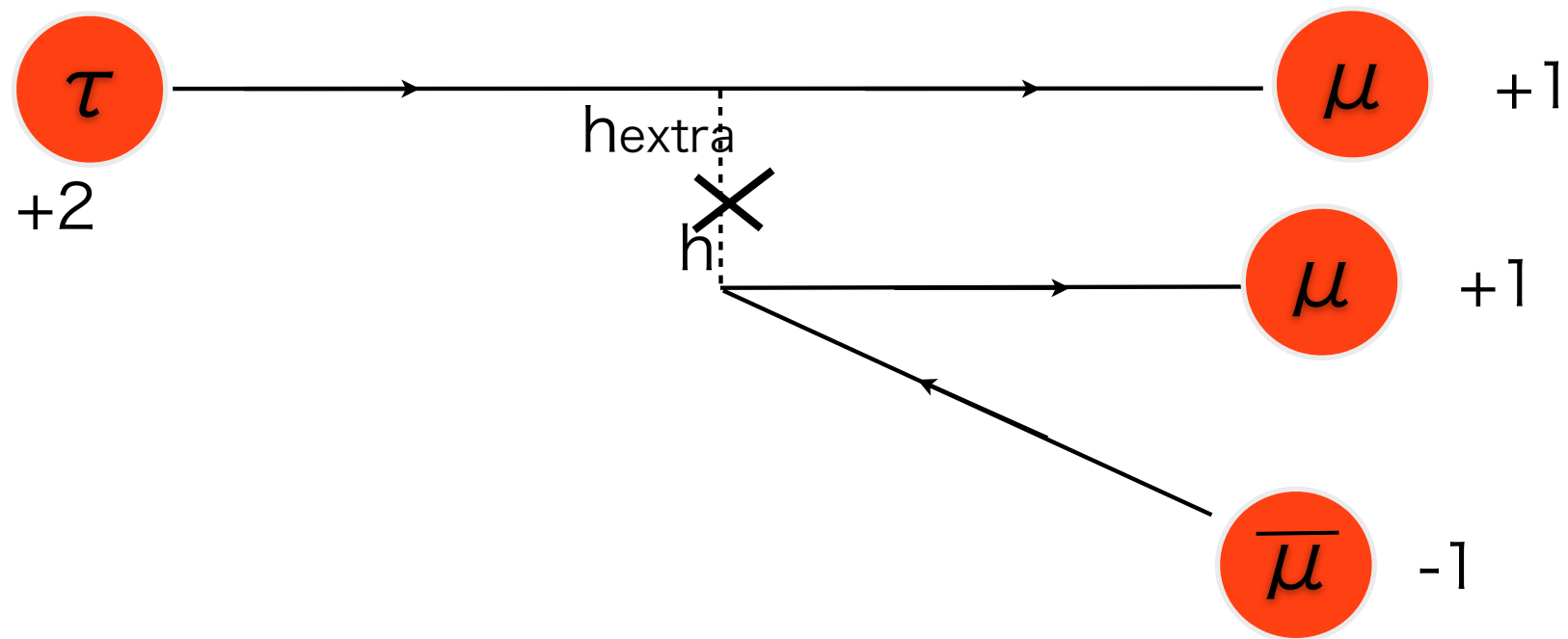


This symmetry cannot be hold in the full lagrangian, but works very well.

(neutrino mass matrix breaks this, but the masses are quite small.)

The other processes are suppressed by the breaking terms in scalar potential

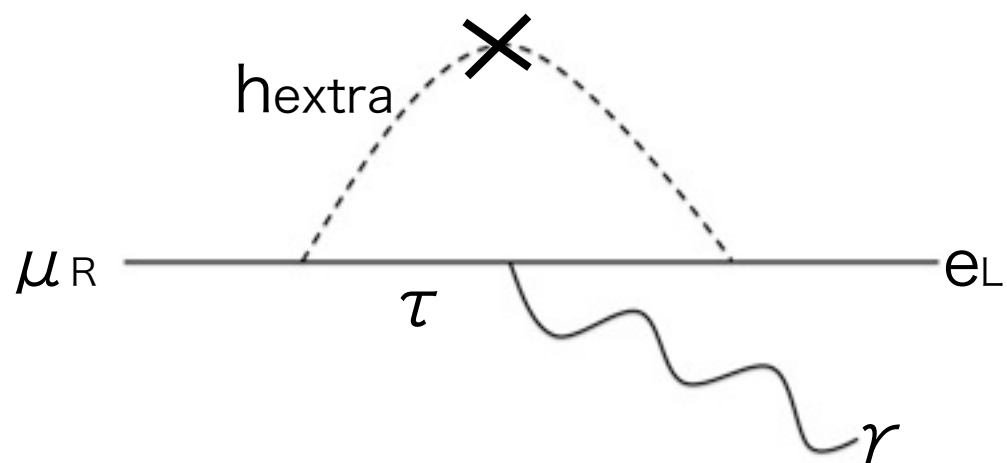
prediction



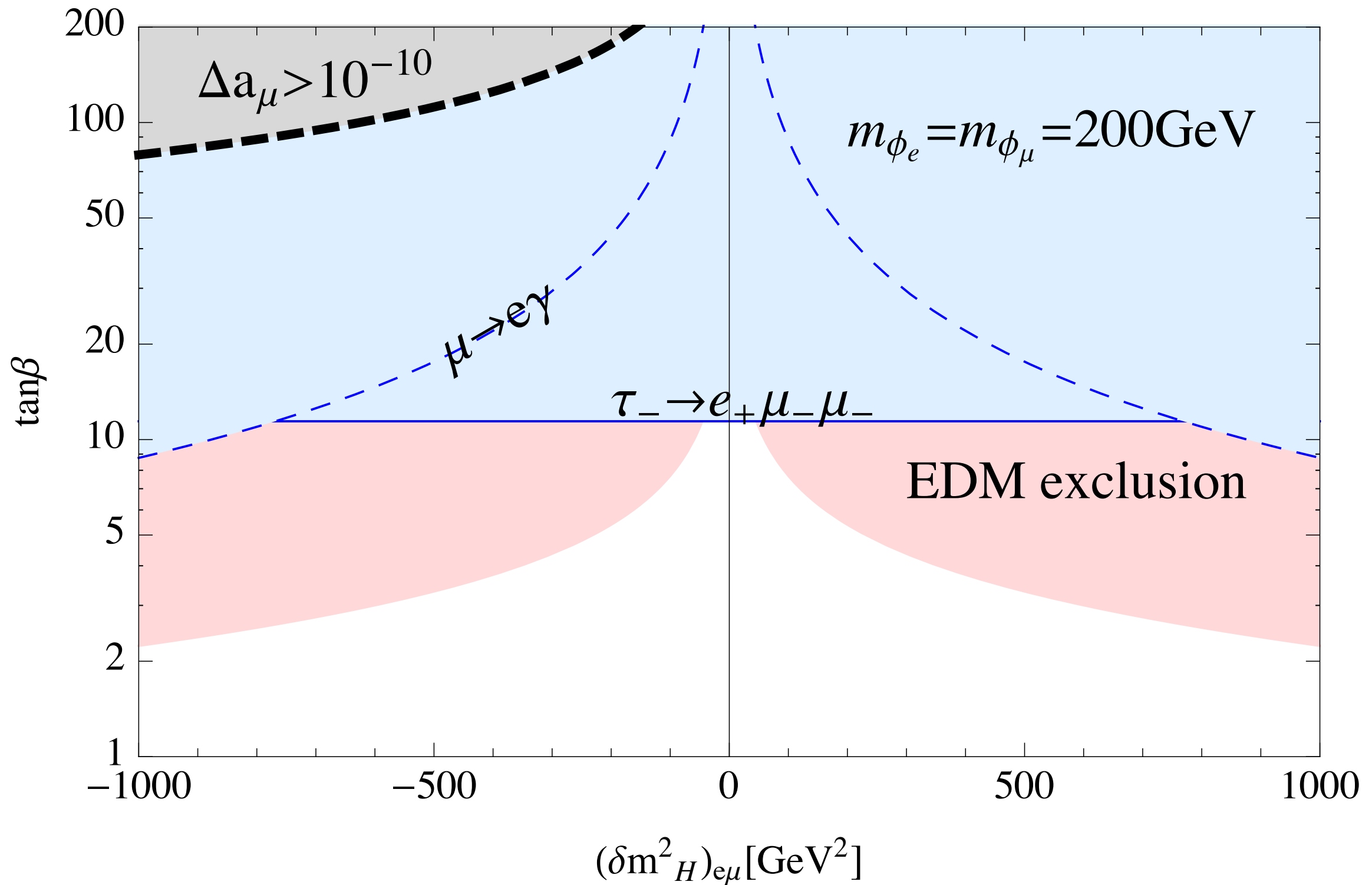
$$\text{Br}(\tau^- \rightarrow 3\mu) \lesssim O(0.01) \times \text{Br}(\tau^- \rightarrow e^+ \mu^- \mu^-)$$

$$\text{Br}(\tau^- \rightarrow 3e) \ll \text{Br}(\tau^- \rightarrow e^+ \mu^- \mu^-)$$

Not only τ decay, but $\mu \rightarrow e \gamma$ process is also important, because of large $Y_{\tau\mu}$, $Y_{e\tau}$ and τ mass enhancement.



$$\text{Br}(\tau^- \rightarrow e\gamma) \lesssim \text{Br}(\mu^- \rightarrow e\gamma)$$



EDM is relevant, depending on phases of Yukawa.

difficult to explain the anomaly of muon $g-2$.

Summary

- Higgs discovered! What is next?
- I hope it is not the end of the story.
- We are waiting for discovery of new physics.
- Flavor violations will be very good processes to find the evidence of new physics. SM predicts tiny flavor violation.

- I introduce flavor violating process in the BSMs, motivated by GUT and flavor symmetry.
- SO(10) GUT predicts flavor violating Z'
 - flavor-universally Z' coupling
 - K physics, μ physics are important.
- Flavor symmetry
 - scalar exchanging causes flavor violation
 - τ , b , t physics important.

Thank you