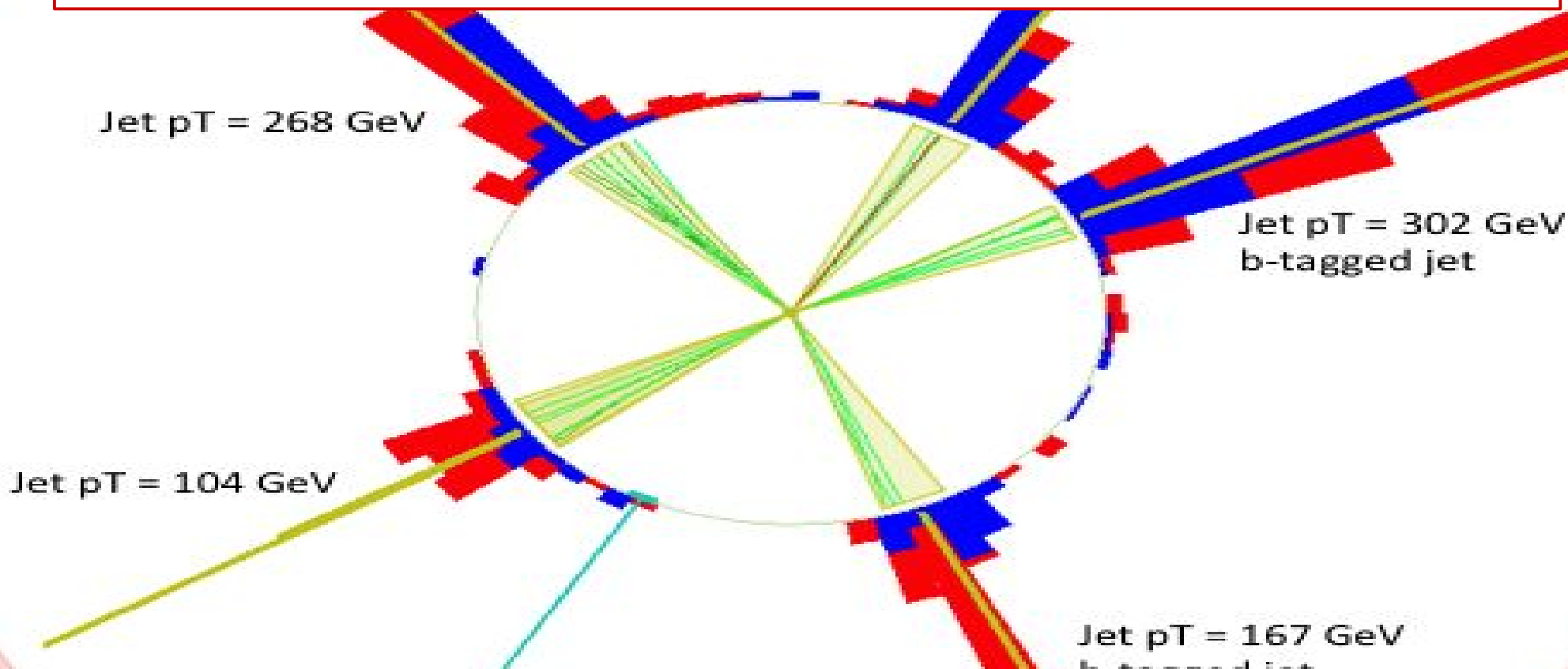


# Searches for R-parity conserving SUSY signatures at CMS



Alessandro Gaz,  
KMI, Nagoya University  
KMI Topics, April 13<sup>th</sup> 2016



# Foreword

- I have been a member of the CMS Collaboration in the period 2009-2014;
- Therefore I have direct experience only of the Run1 data, for Run2, I only participated in the analysis preparation and sensitivity studies.

# The Higgs boson exists!

- After decades of searches, finally the Higgs boson is in the bag (and its mass is  $\sim 125$  GeV);

- It was the last missing piece of the Standard Model, but still many questions to answer:

→ What is the Dark Matter in the Universe?

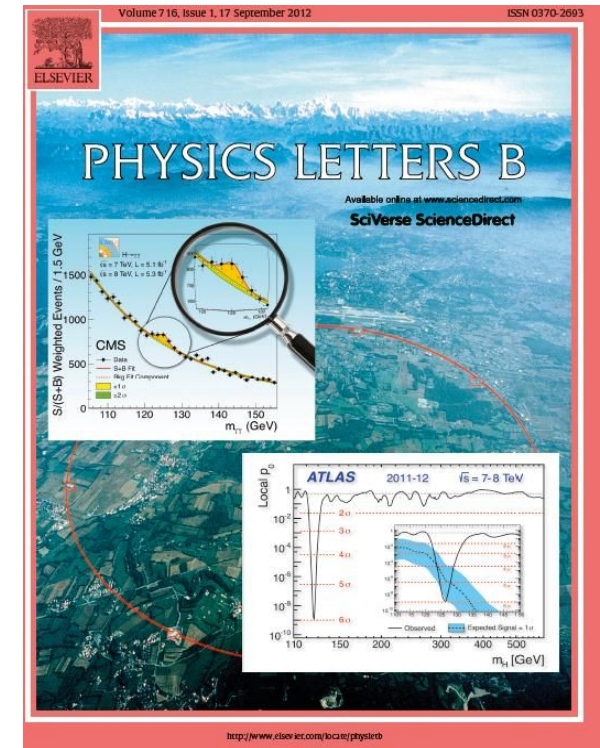
→ What is the origin of the matter/antimatter asymmetry?

→ Neutrinos seem to be special...

→ What's the connection between the Cosmic Inflation and Particle Physics?

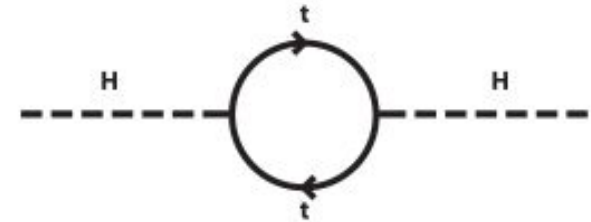
→ ...

- But also explaining why the Higgs is where we found it is not so trivial...



# The Hierarchy problem

- Compared to the other SM particles, the Higgs is “special”;
- We know that New Particles must exist, possibly at energy scales of  $10^{15} - 10^{19}$  GeV;
- The Higgs couples (possibly indirectly, through loops) with those particles, and its mass receives corrections from the interaction with those particles;
- Unlike the other SM particles, the Higgs is not protected by any symmetry...;
- ... thus we would expect its mass to be at the same scale of that of the New Physics particles.



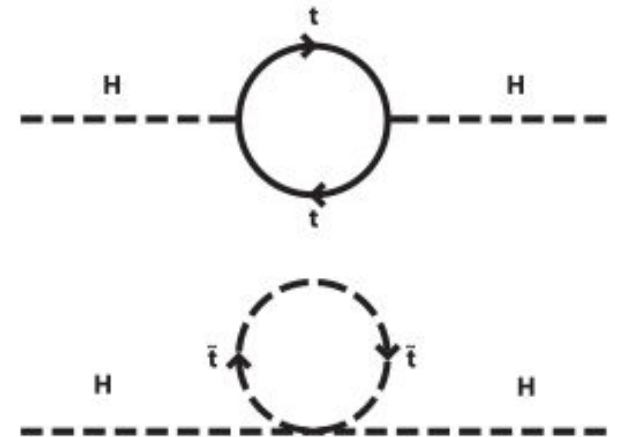
# Possible ways out

How can we explain the fact that the Higgs mass is  $O(100)$  GeV?

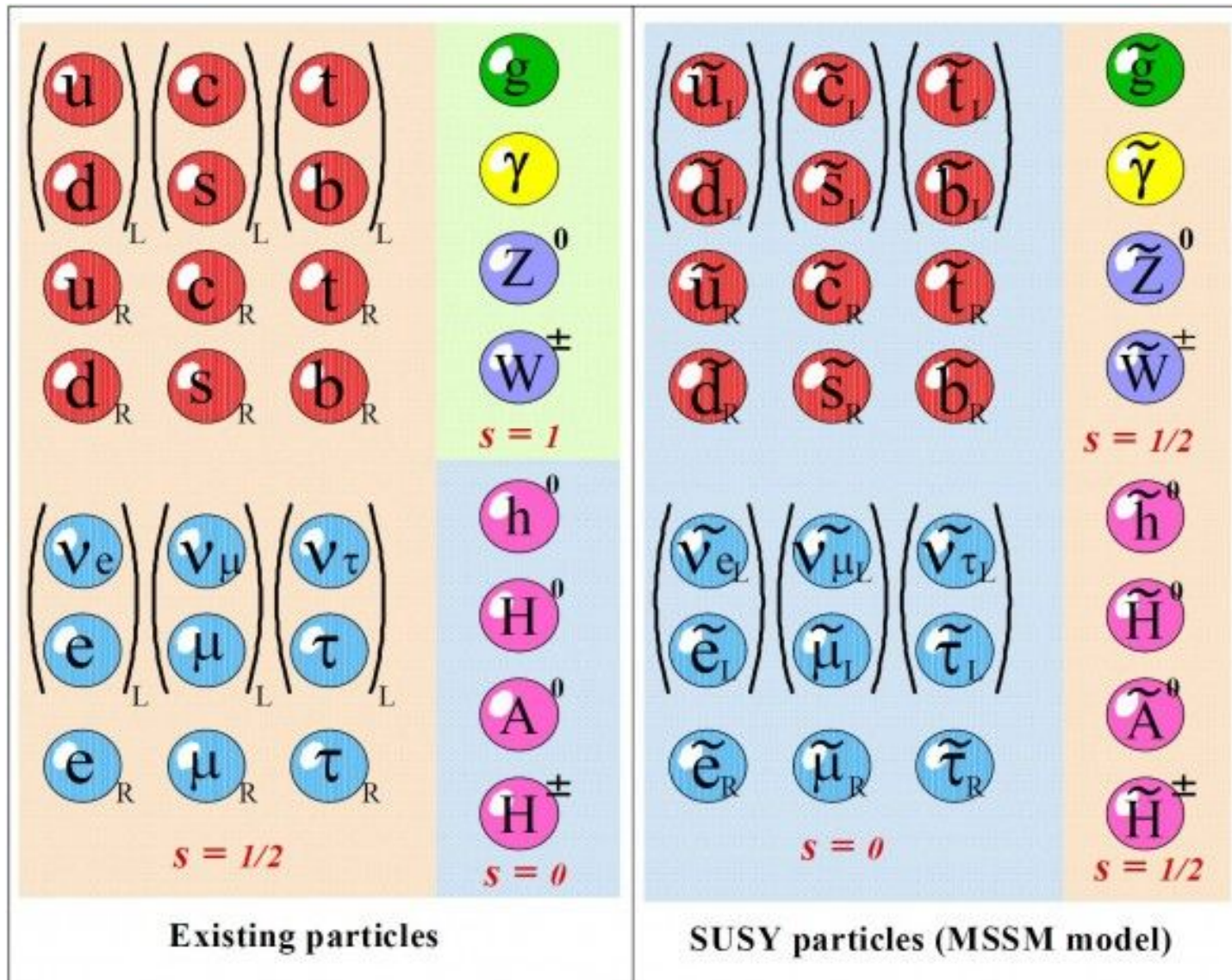
1) There is an extraordinary cancellation among all the corrections to the Higgs mass. “Fine tuning”: required precision better than  $\sim 1/10^{26}$ .

2) There is a mechanism at work so that the sum of all the corrections is reduced to a small value.

Basic idea of Supersymmetry: for each SM particle, associate another particle, with spin differing by  $\frac{1}{2}$  a unit. This would give a contribution with opposite sign, and the closer the masses of the two partners, the better the cancellation mechanism.



# SUSY



# Natural SUSY

- SUSY works well as long as the masses of the s-particles are not much higher than those of the SM particles;
- This is particularly important for the higgsino(s), stop(s), and gluino:

$$\delta m_H^2 \sim -\frac{y_t^2}{\pi^2} \frac{\alpha_s}{\pi} m_{\underline{gluino}}^2 \left( \log \frac{\Lambda}{m_{\underline{gluino}}} \right)^2 \quad \text{two-loops}$$

$$\delta m_H^2 \sim -\frac{3}{8\pi^2} y_t^2 m_{\underline{stop}}^2 \log \frac{\Lambda}{m_{\underline{stop}}} \quad \text{one-loop}$$

$$\delta m_H^2 \sim |\underline{\mu}_{\underline{higgsino}}|^2 \quad \text{tree}$$

- If the masses of these SUSY particles are way beyond the TeV scale, we still need a lot of fine tuning!

# R-parity

- Another ingredient, the **multiplicative** quantum number:

$$P_R = (-1)^{3B+L+2s}$$
$$P_R(\mathbf{x}) = +1$$
$$P_R(\tilde{\mathbf{x}}) = -1$$

- In most popular SUSY realization this quantum number is conserved (it helps making the proton lifetime compatible with the current limits);
- This has two important consequences:
  - 1) by colliding SM particles, SUSY particles can only be produced in pairs;
  - 2) the lightest SUSY particle (LSP) is stable. In most scenarios the LSP is neutral and weakly interacting, so it effectively behaves like a neutrino.  
If its mass is in the 100 GeV – 1 TeV range it is also an interesting candidate for the Dark Matter.



# Outline

- The Large Hadron Collider (LHC) and the Compact Muon Solenoid (CMS) experiment;
- Search strategies;
- Some example Run1 analyses:
  - gluino pair production;
  - direct stop pair production;
  - production of electroweakinos decaying to Higgs bosons;
- A look at Run2 data;
- Future prospects.

# The LHC

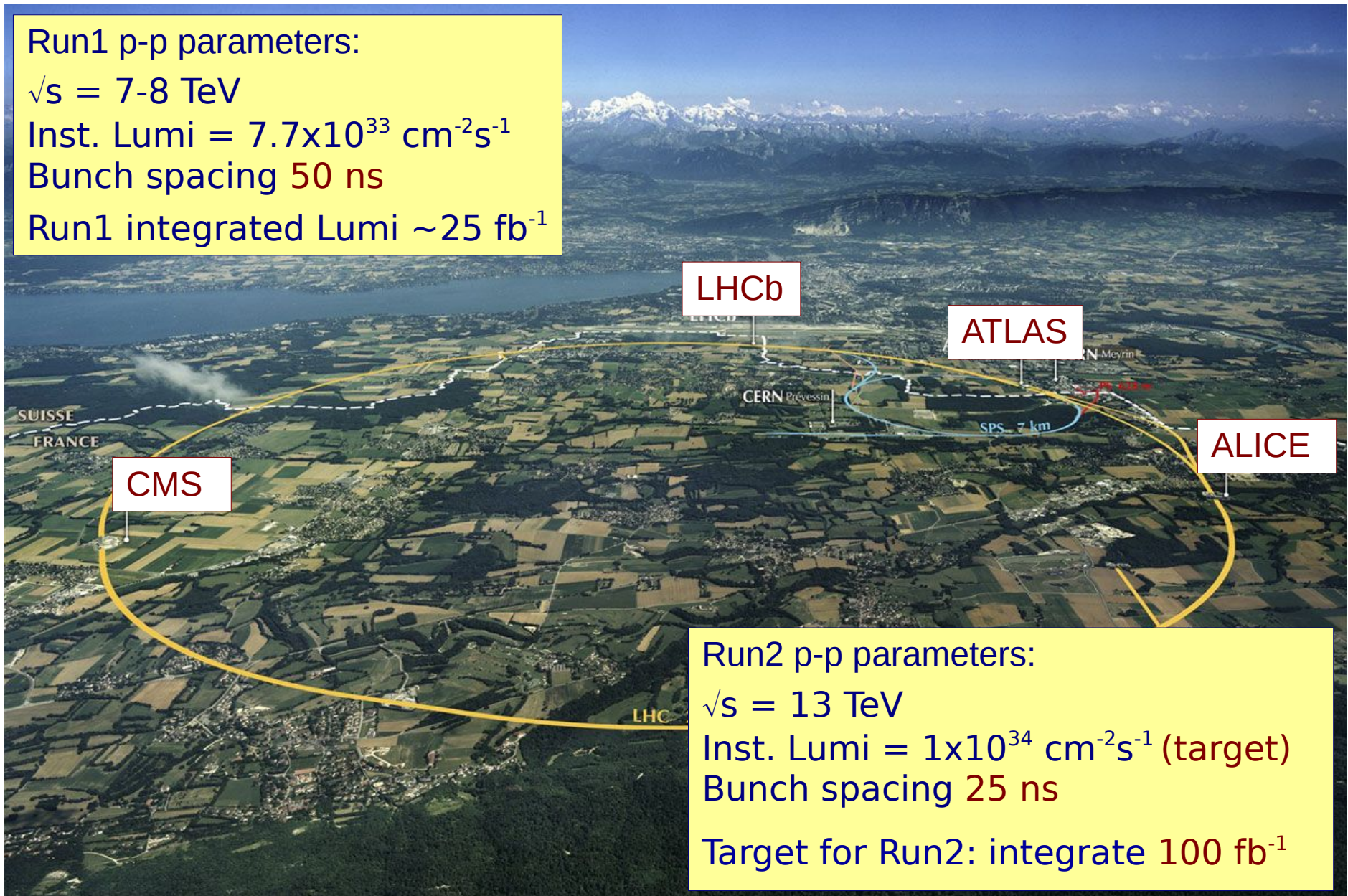
Run1 p-p parameters:

$$\sqrt{s} = 7\text{-}8 \text{ TeV}$$

$$\text{Inst. Lumi} = 7.7 \times 10^{33} \text{ cm}^{-2}\text{s}^{-1}$$

Bunch spacing **50 ns**

Run1 integrated Lumi  $\sim 25 \text{ fb}^{-1}$



Run2 p-p parameters:

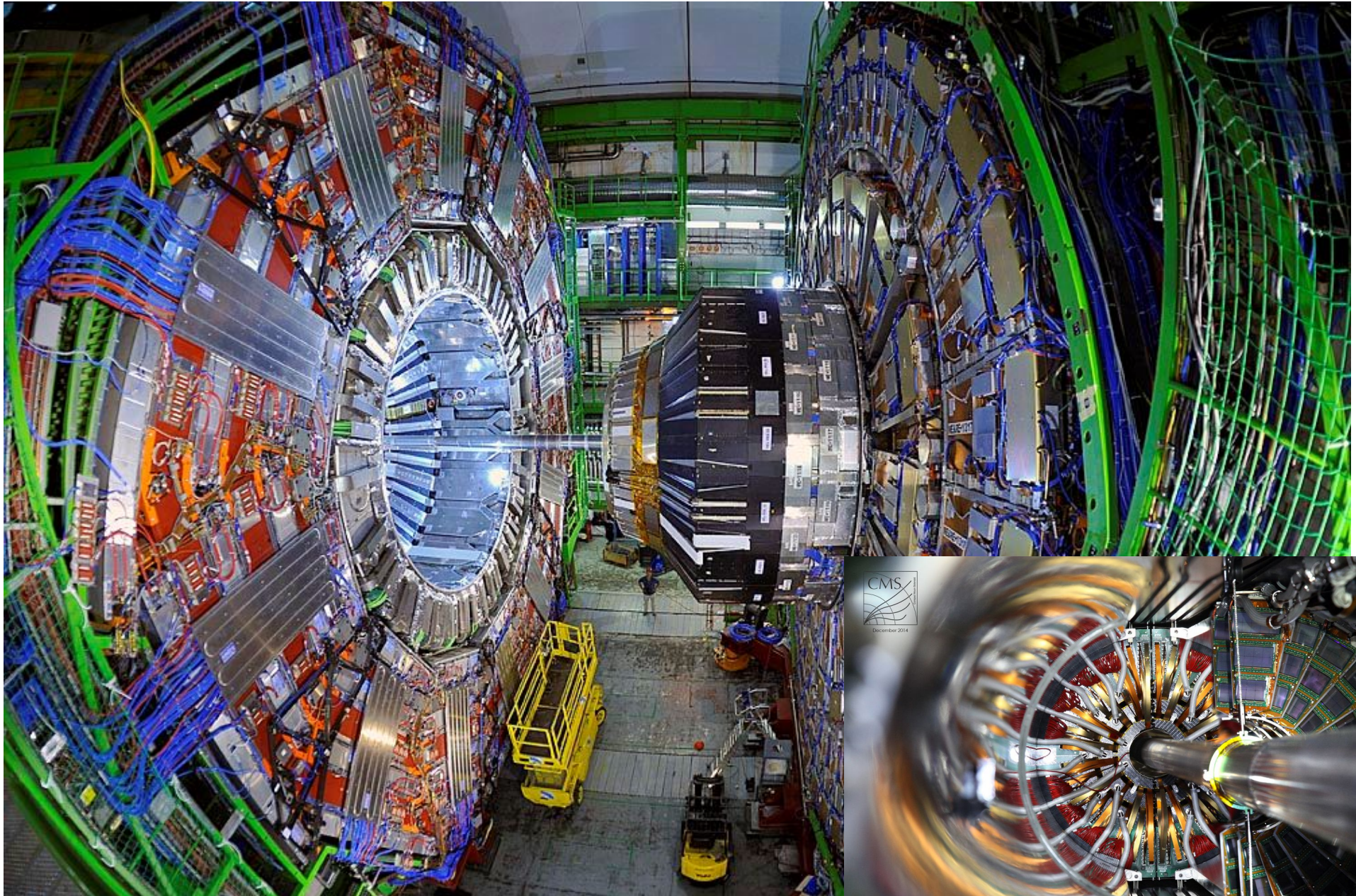
$$\sqrt{s} = 13 \text{ TeV}$$

$$\text{Inst. Lumi} = 1 \times 10^{34} \text{ cm}^{-2}\text{s}^{-1} \text{ (target)}$$

Bunch spacing **25 ns**

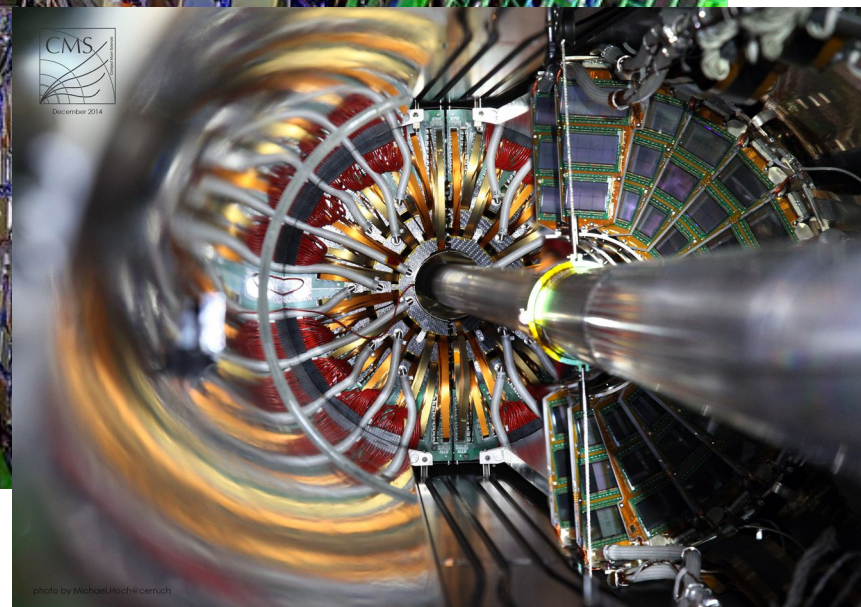
Target for Run2: integrate **100 fb<sup>-1</sup>**

# The CMS Experiment

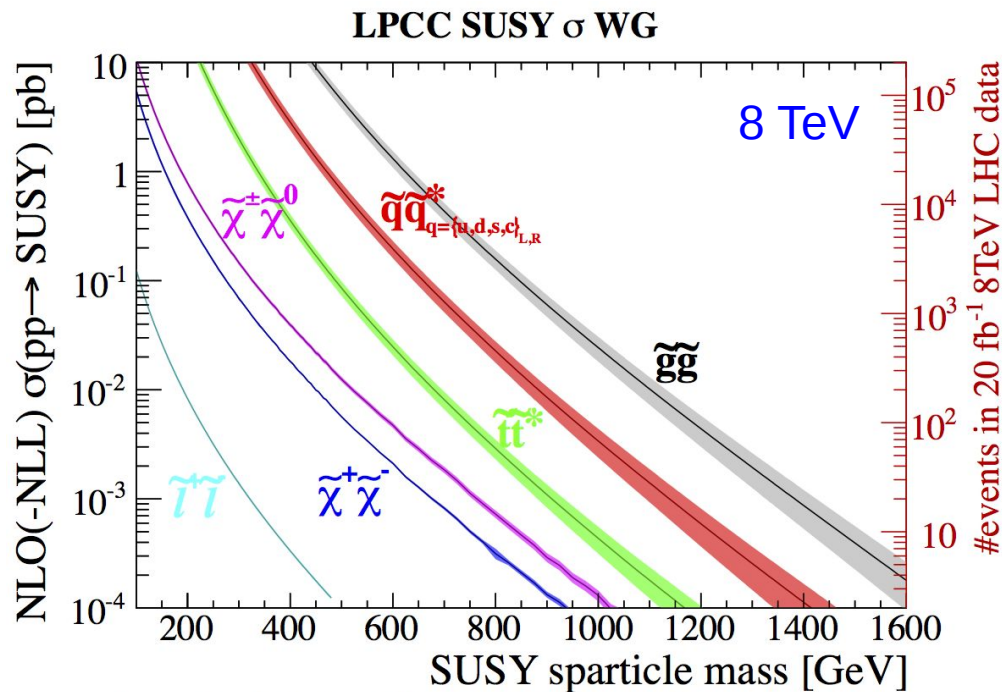


April 13th 2016

A. Gaz

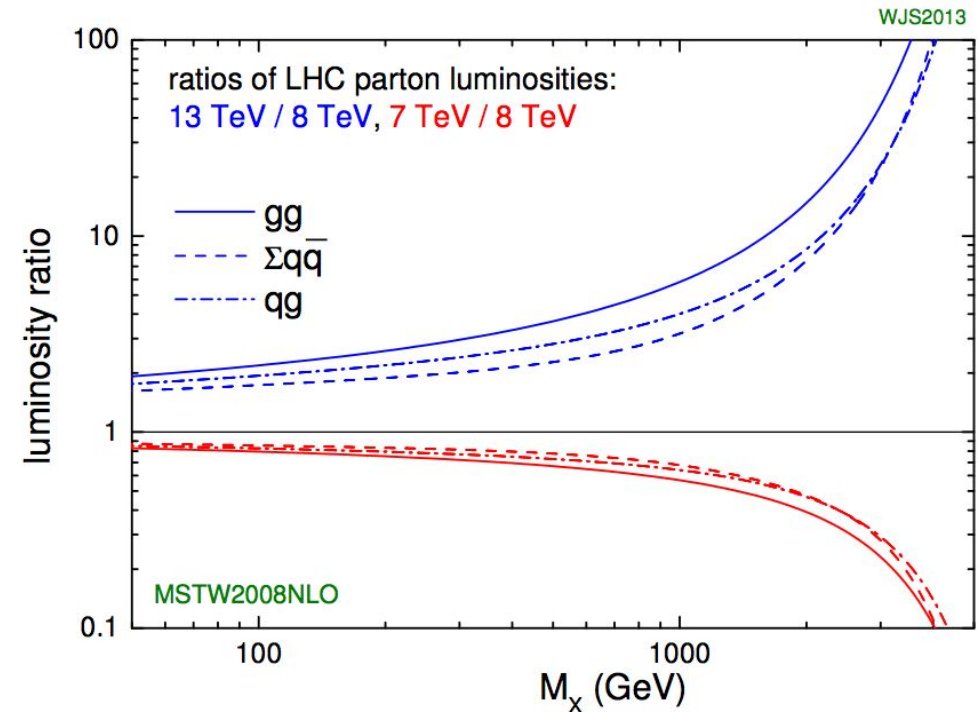


# SUSY production at LHC



<https://twiki.cern.ch/twiki/bin/view/LHCPhysics/SUSYCrossSections>

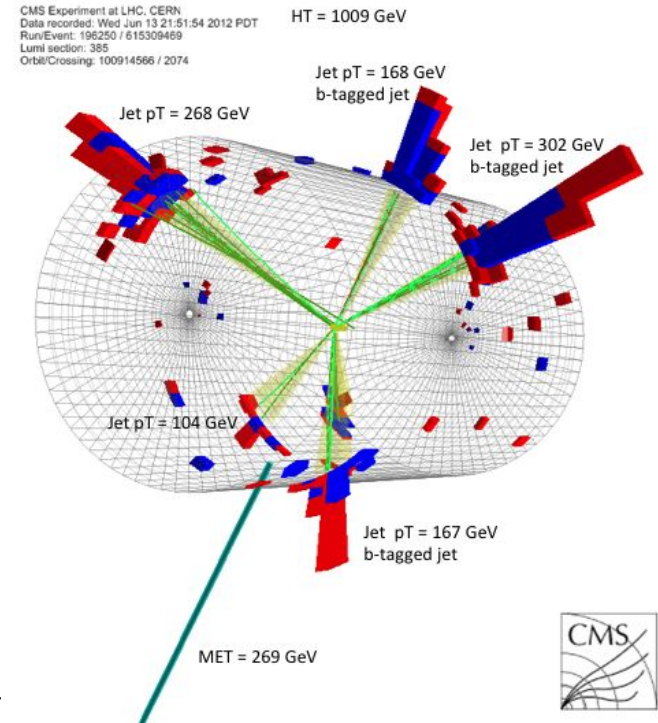
arXiv:1206.2892



- Production of colored (i.e. strongly interacting) particles clearly favored at a hadron collider...
- ... still neutralinos are expected to be among the lightest SUSY particles, so the production could be significant!

# Search strategies

- I will focus on RP-conserving scenarios: the LSP is neutral, weakly interacting, and stable. So it behaves like a (possibly heavy) neutrino;
- The common feature for these events is the presence of large **missing transverse energy** (MET);
- Together with MET, other selection criteria are applied, to increase sensitivity to specific topologies and reduce the impact of SM backgrounds:
  - presence or absence of ( $\tau$ ) leptons;
  - presence of vector bosons, Higgs, ...;
  - total energy in the event;
  - ...

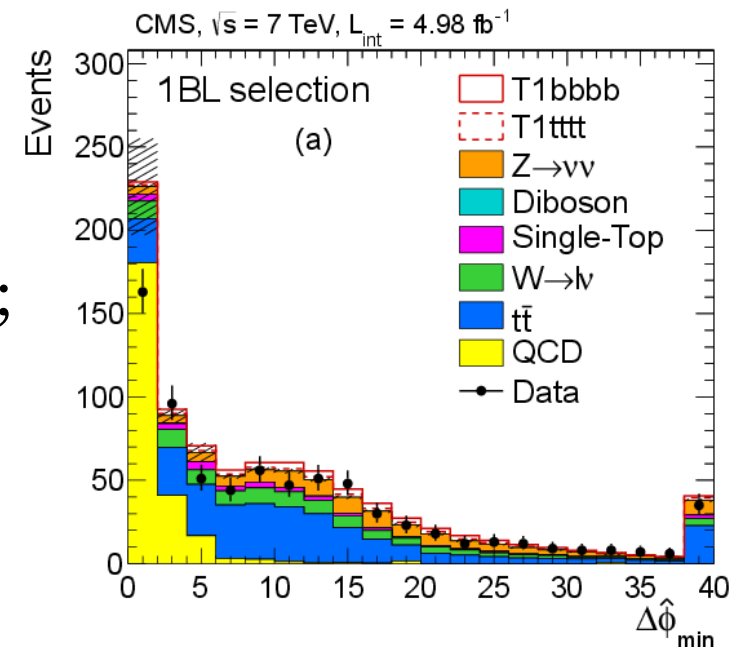


# Standard Model backgrounds

- Missing transverse energy is not exclusively a signature for SUSY (or other New Physics), many SM processes can mimic SUSY events;

- Examples:

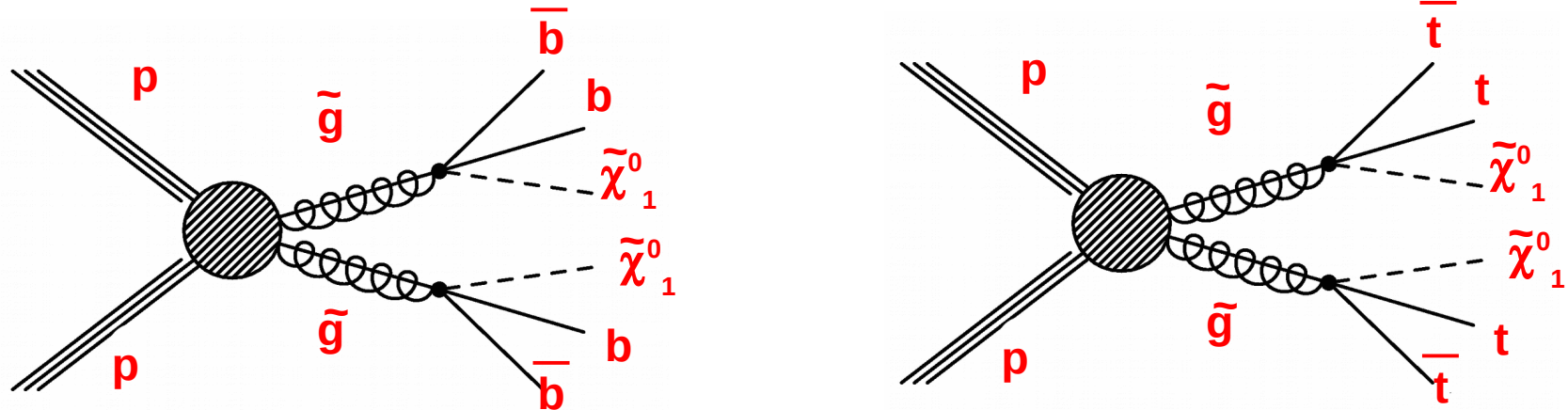
- multijet events, in which MET arises from incorrectly measured jet energy;
- $t\bar{t}$  events, with one top decaying semileptonically ( $t \rightarrow b l \nu$ )
- $Z \rightarrow \nu\nu$  events;



- A large fraction of the analysis effort is devoted to carefully estimating the SM backgrounds (from data control samples, as much as possible);
- The hope is to observe some topology in which the event yields are incompatible with the SM background predictions.

# Search for gluino pair production

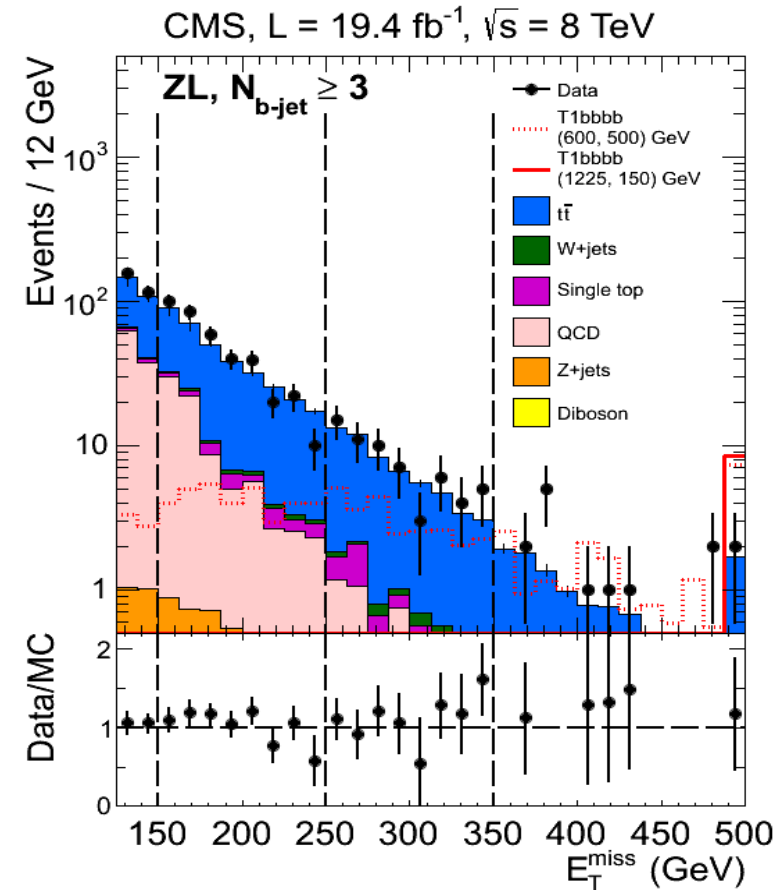
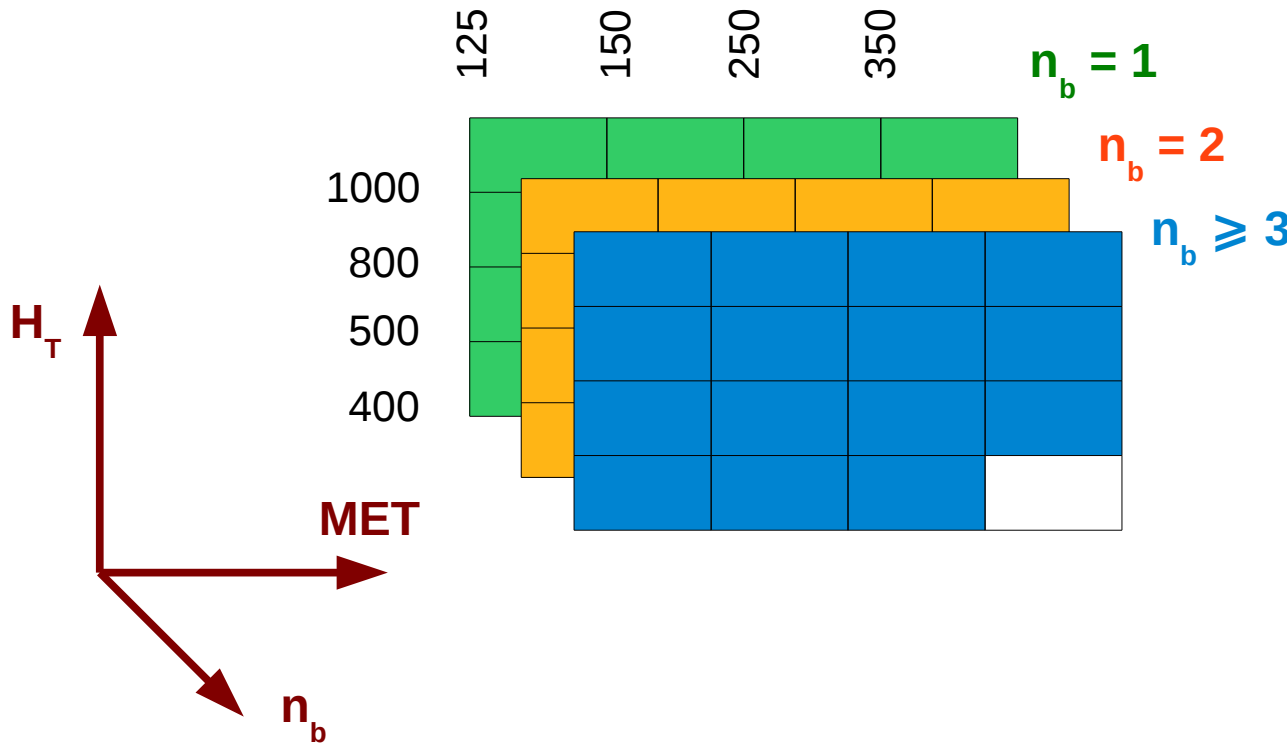
- Target topologies:



- In Natural SUSY scenarios, s-tops and s-bottoms are expected to be relatively light, so searches with “3<sup>rd</sup> generation” particles are particularly important;
- Main signatures:
  - large MET and hadronic energy;
  - high b-jet multiplicity.

arXiv: 1305.2390 [hep-ex],  
Phys. Lett. **B725** 243 (2013)

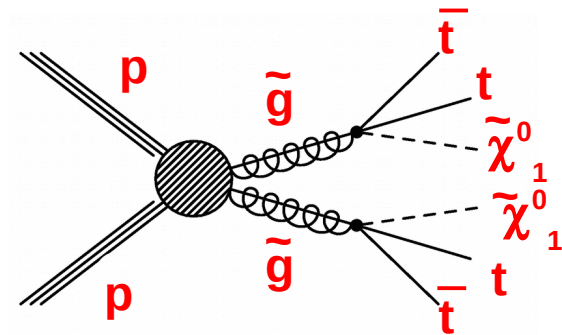
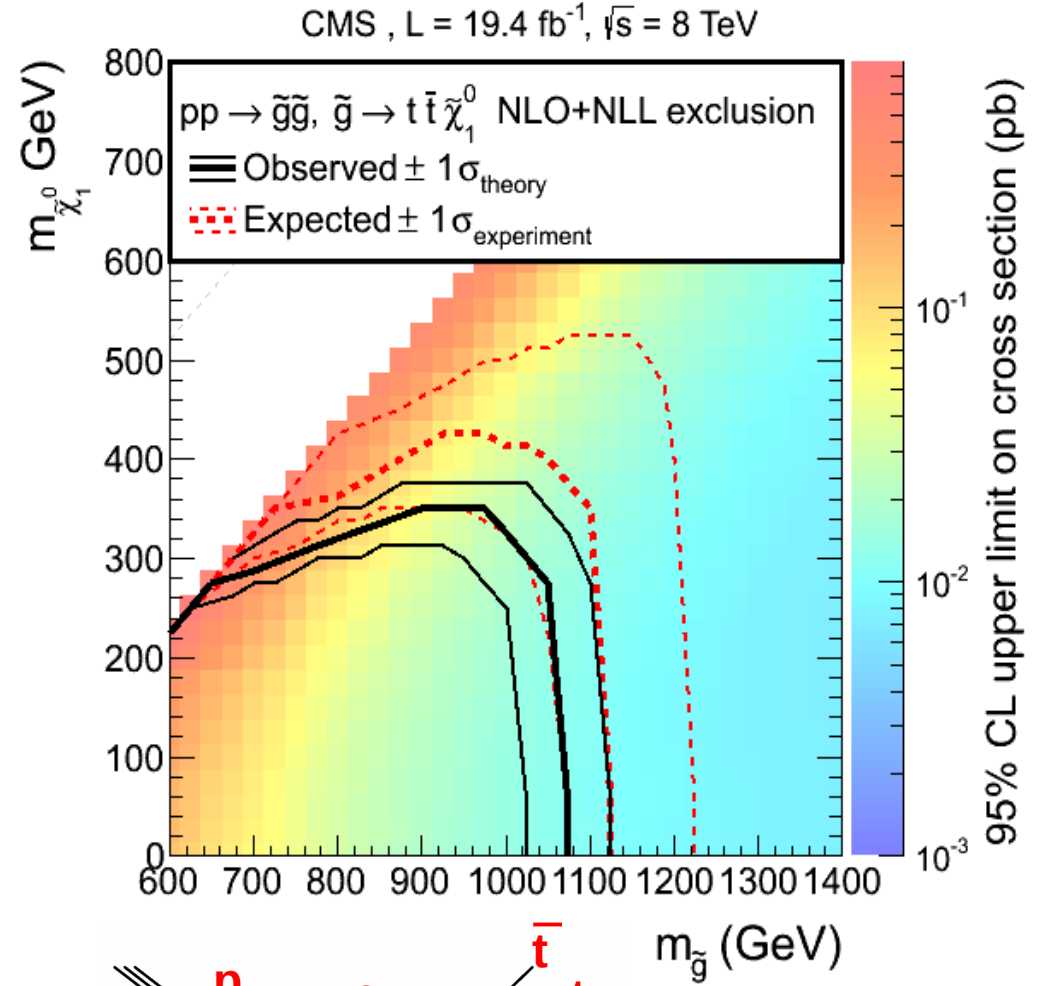
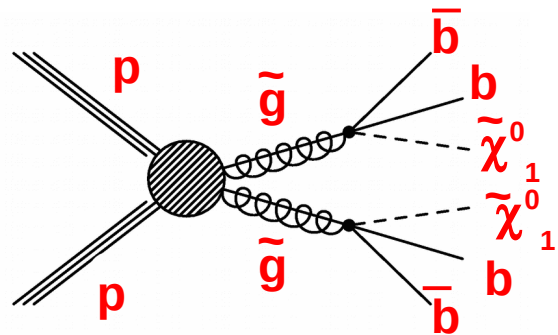
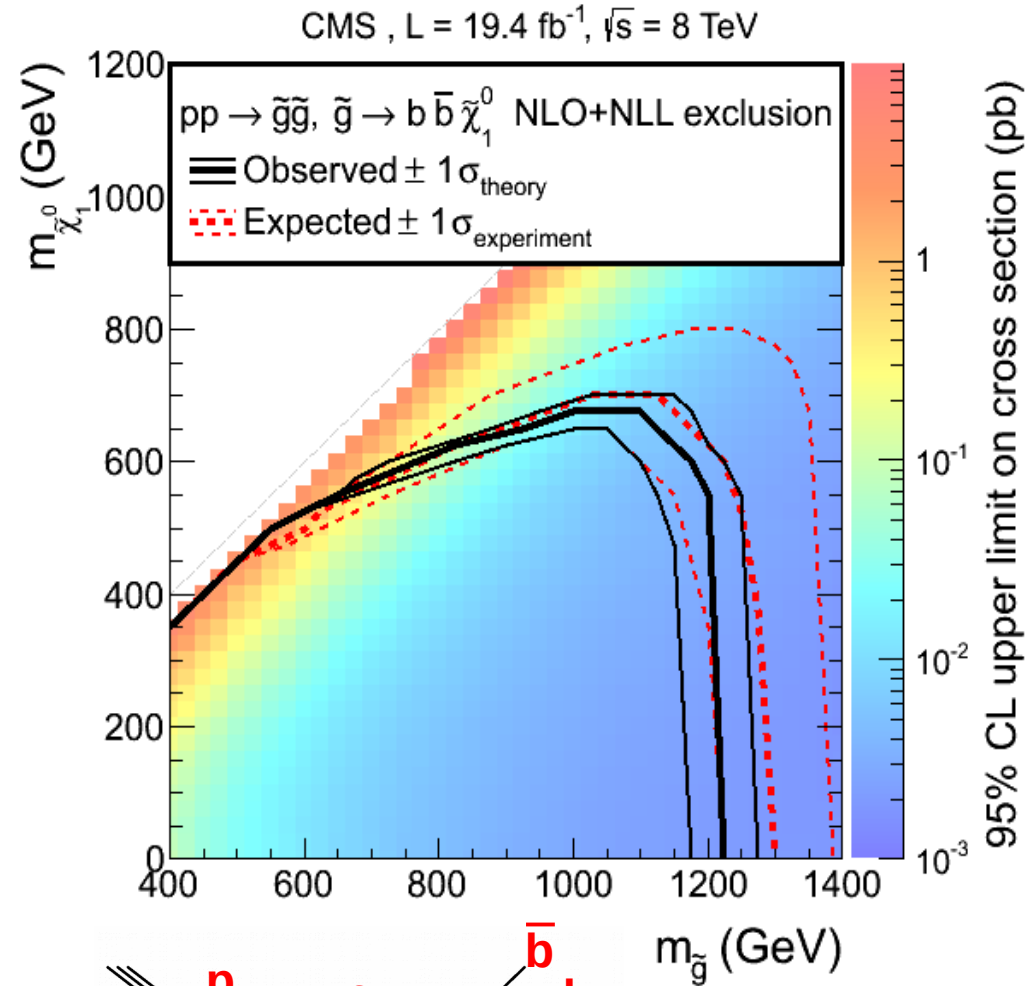
# Search for gluino pair production



- 3-dimensional analysis on the variables MET, scalar sum of jet momenta ( $H_T$ ), and b-jet multiplicity;
- 45 independent signal regions, correlations from nuisance parameters taken into account;
- No excess wrt SM background predictions...

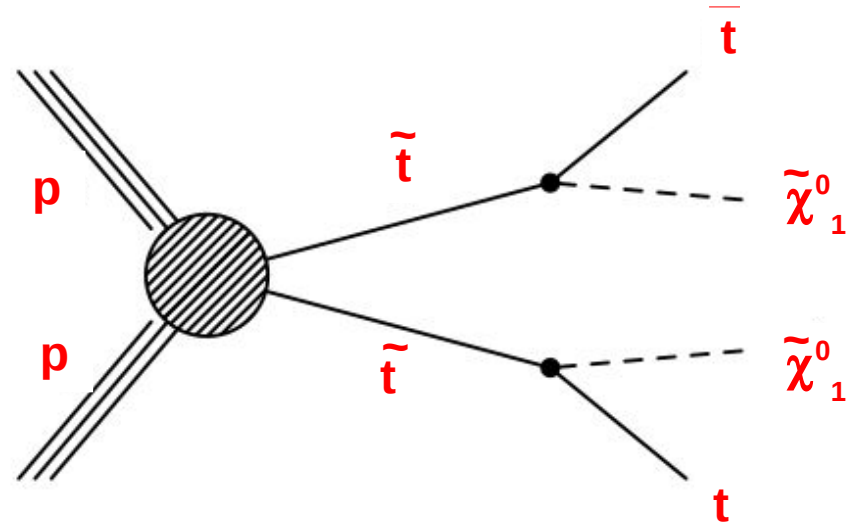


# Search for gluino pair production



# Search for direct stop pair production

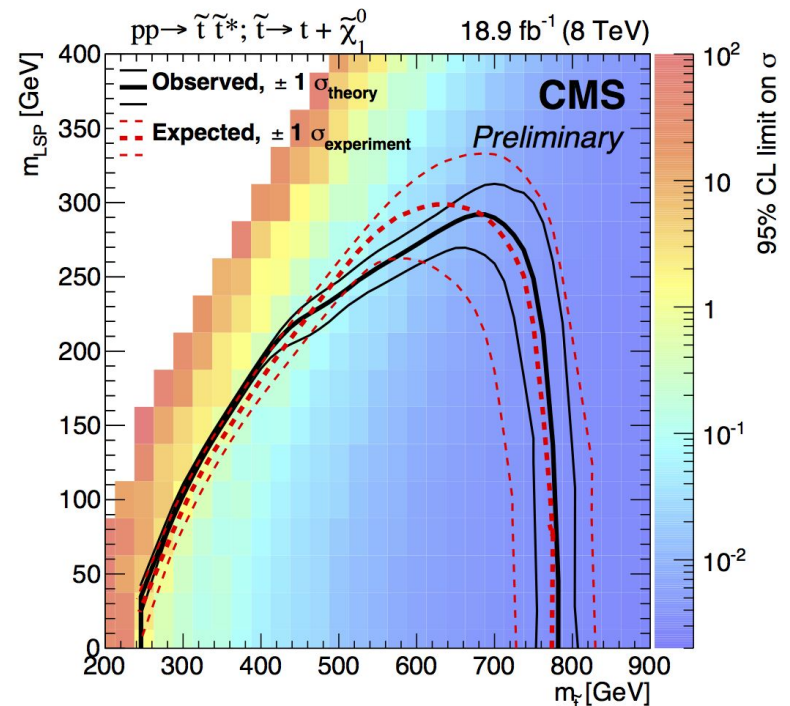
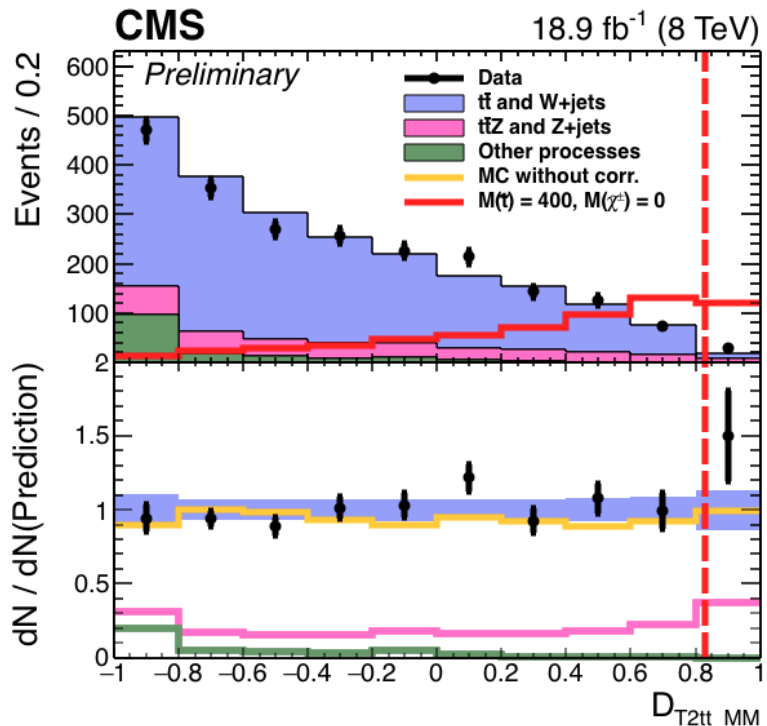
- The stop is expected to be light, in order to keep the fine tuning reasonably low;
- Expect cross section for direct production to be reasonably high, but...;
- The signature is similar to that of a  $t\bar{t}$  event, the only difference is the presence of MET;
- If  $m(\tilde{t}) \sim m(t) + m(\tilde{\chi}_1^0)$ , the neutralinos carry little momentum (and the events will look like SM  $t\bar{t}$ );
- Need a strategy to cover also the “blind spots” of the phase space: maybe the stop is there but we haven't been able to see it yet.



# Search for direct stop pair production

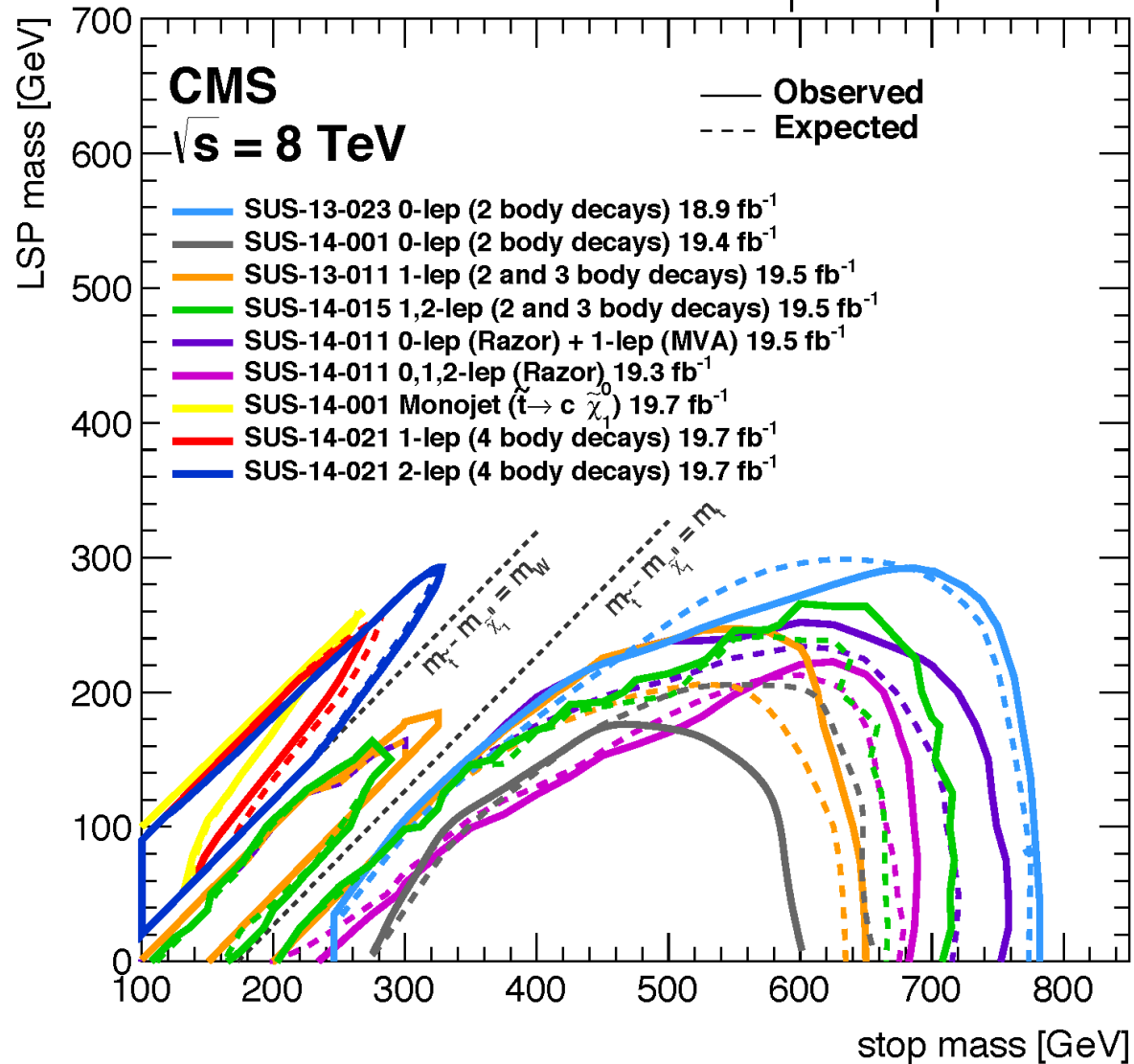
- Search performed on fully hadronic final states;
- Variable-cone jet reconstruction to optimize top selection;
- Multivariate BDT discriminant to enhance the potential signal from the dominant  $t\bar{t}$  background.

**CMS-SUS-13-023**



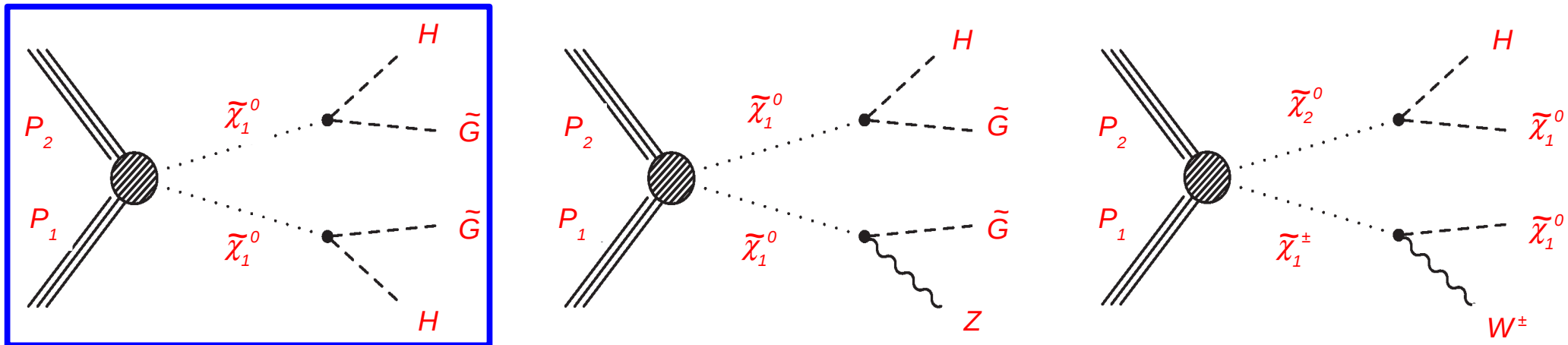
# Search for direct stop pair production

$\tilde{t}\tilde{t}^*$  production,  $\tilde{t} \rightarrow t \tilde{\chi}_1^0 / c \tilde{\chi}_1^0$



# Ewkinos production with Higgs in the final state

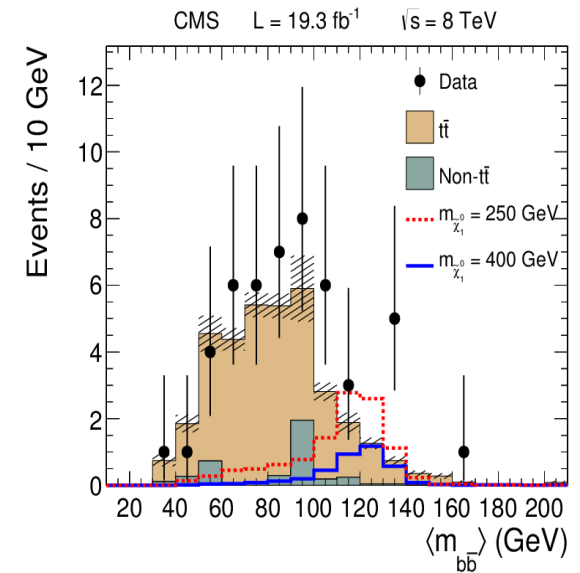
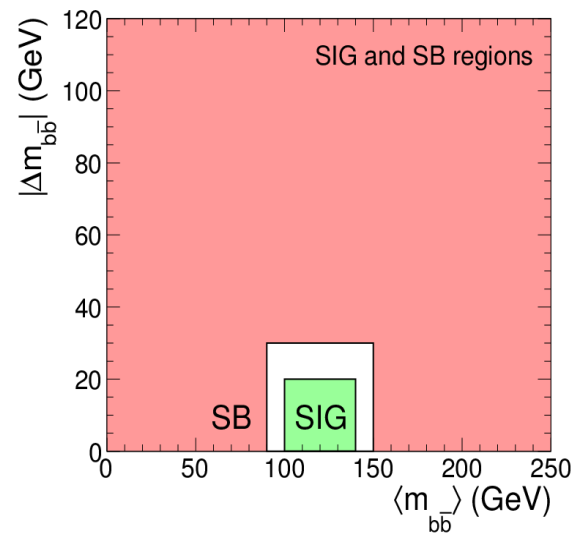
- Ewkinos (neutralinos and charginos) are expected to be light and couple naturally to the Higgs and the vector bosons;
- Several topologies are possible:



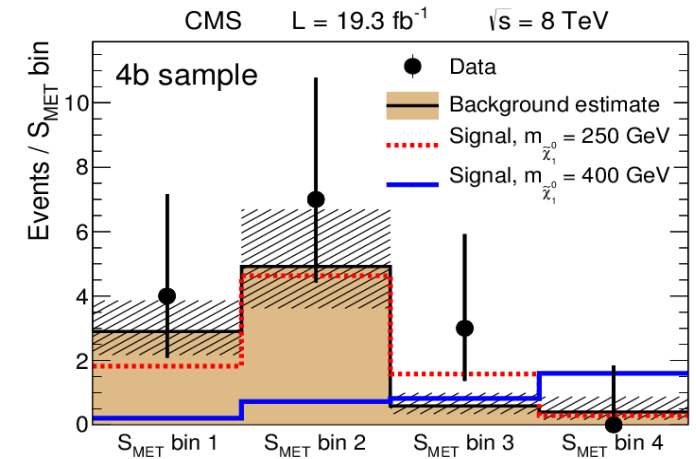
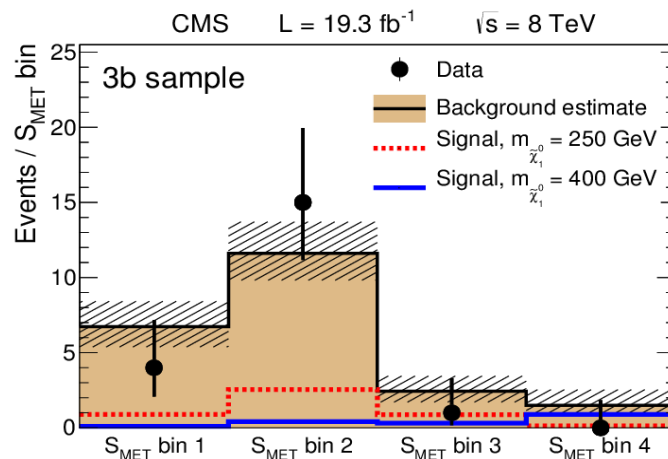
- Basically we search for events containing Higgs (or W and Z) bosons, associated with a non SM source of MET;
- Today I will cover only the topology with two Higgs's in the final state, both of them decaying to a  $b\bar{b}$  pair.

# Ewino production with Higgs in the final state

- Search for events with 4 b-jets, and combine them to make two Higgs candidates;
- Dominant background is from  $t\bar{t}$  events;
- Discriminating variable: MET significance.

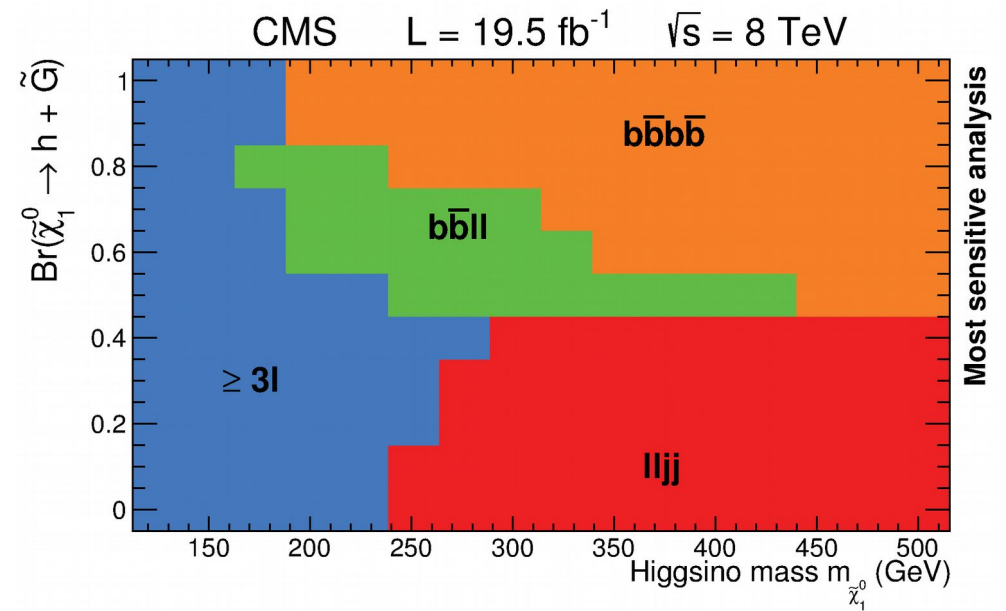
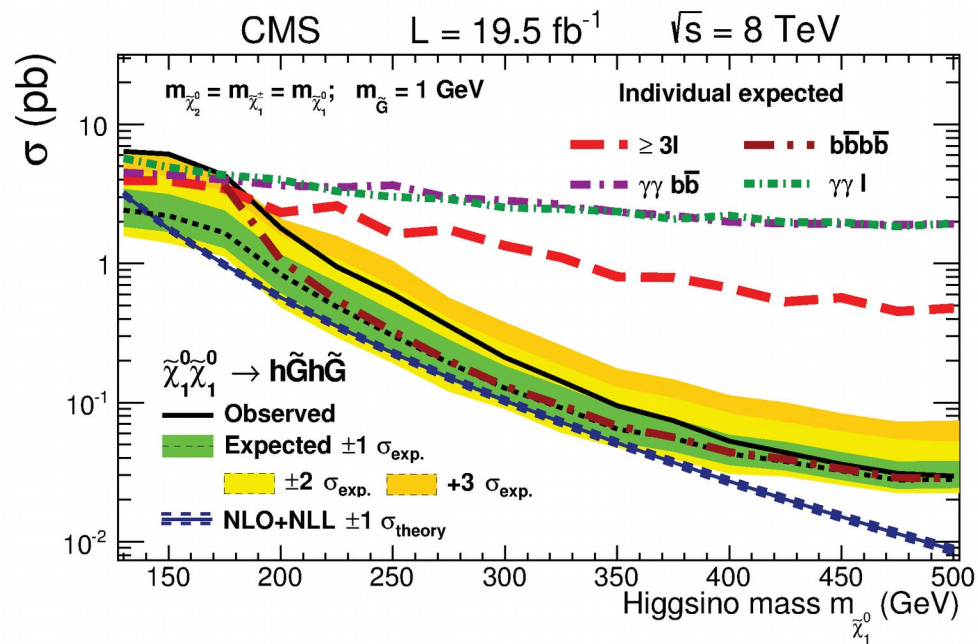


Again, good agreement between data and expected SM backgrounds



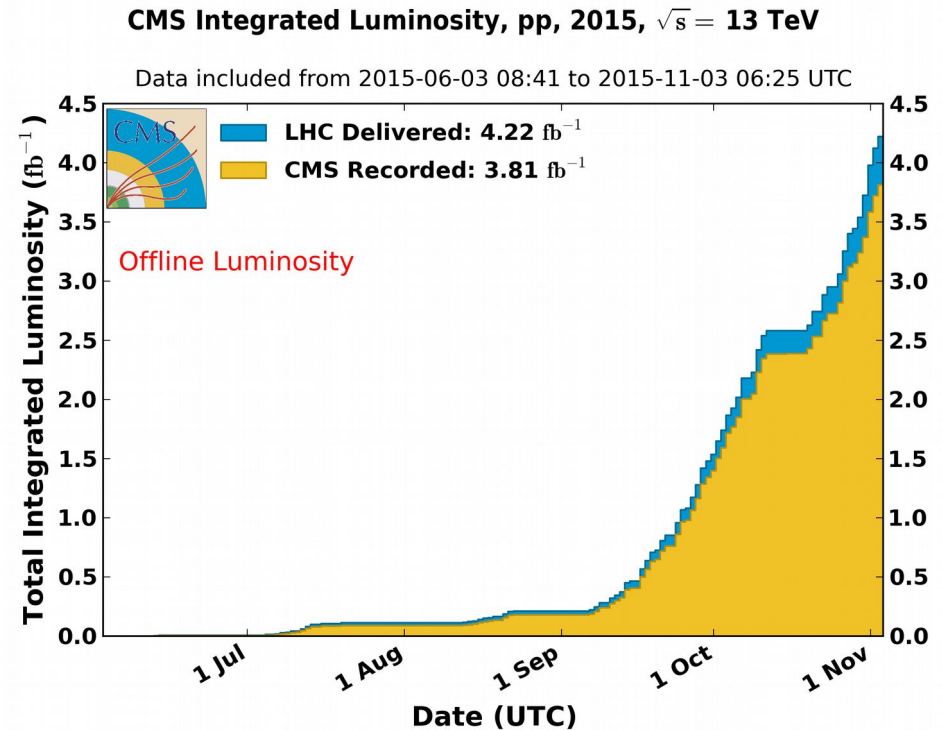
# Ewino production with Higgs in the final state

- Not yet able to get significant exclusion with Run1 data, but we are getting really close...



# LHC Run2

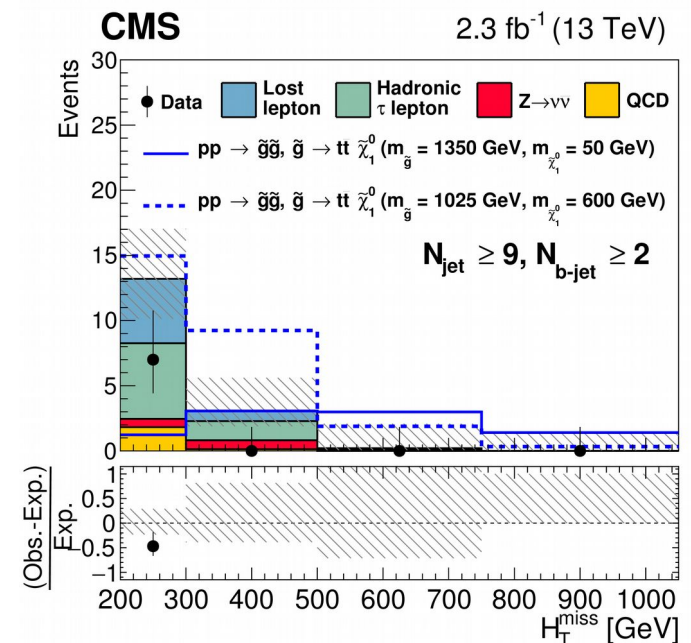
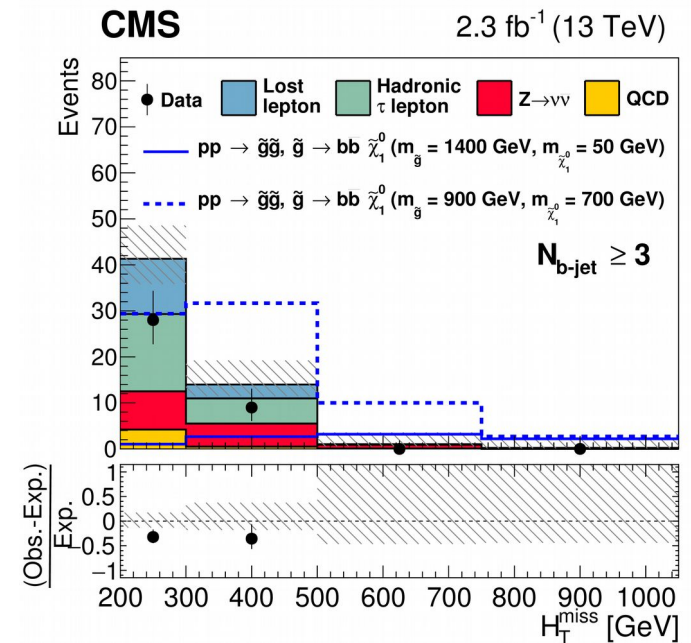
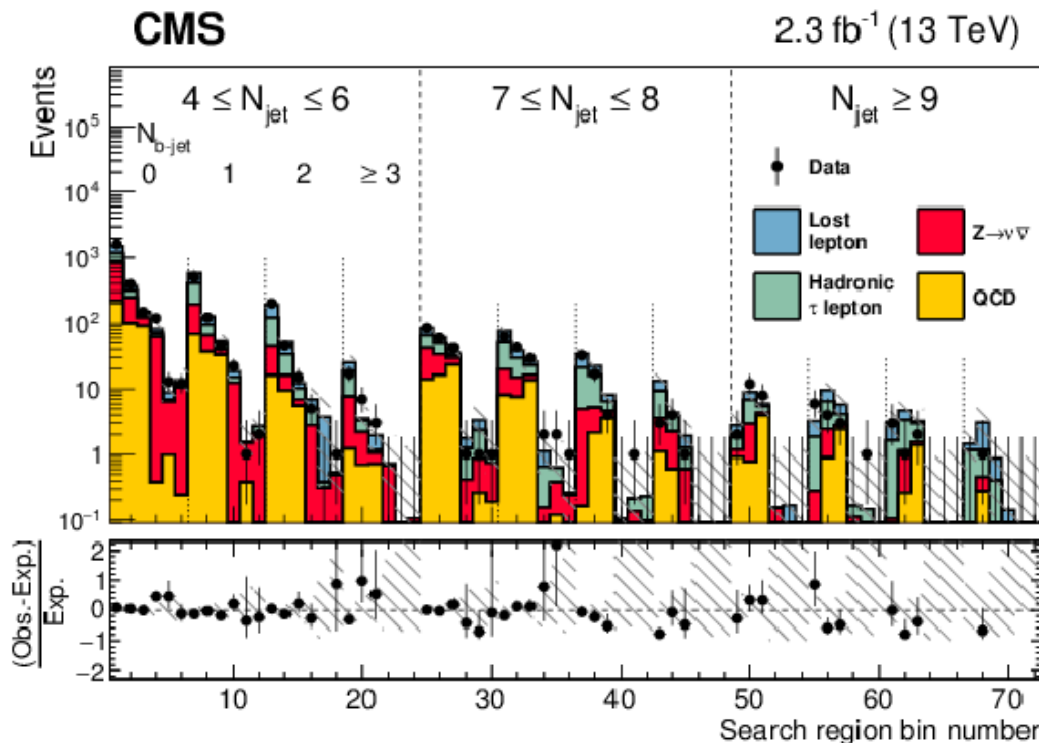
- What I have shown so far is based on the  $\sim 20 \text{ fb}^{-1}$  of the 8 TeV 2012 Run;
- After a long shutdown, the LHC resumed operations in 2015, raising the collision energy to 13 TeV;
- The LHC delivered  $\sim 4.2 \text{ fb}^{-1}$ , which, after accounting for data-taking inefficiency, data quality requirements, and the problem with the magnet, reduce to  $\sim 2.3 \text{ fb}^{-1}$  available for analysis;
- Anyway, thanks to the higher cross-sections this data sample is sufficient to increase the sensitivity of many Run1 analyses.





# Searches for gluino pair production

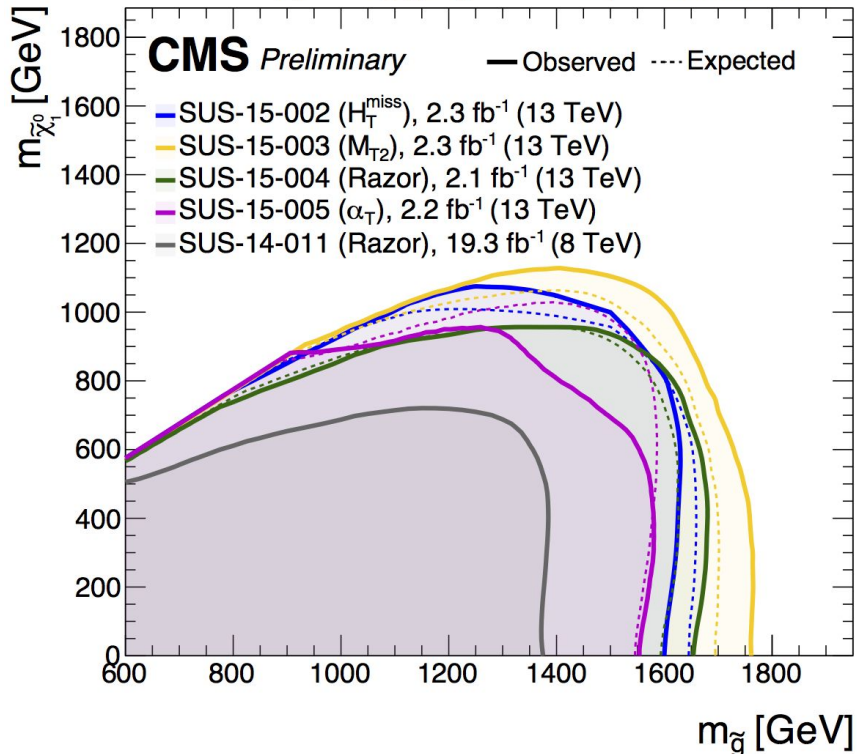
- 4-dimensional analysis in missing  $H_T$ ,  $H_T$ , jet multiplicity and b-jet multiplicity;
- 72 search regions in total.



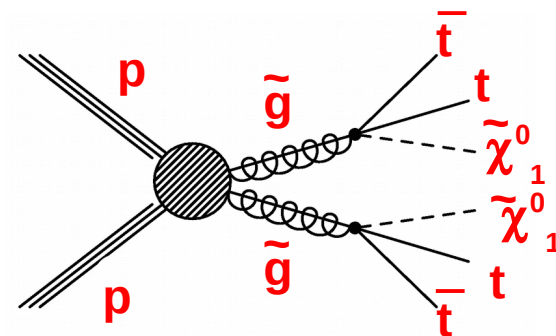
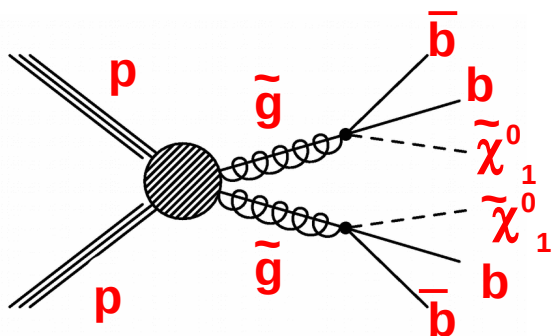
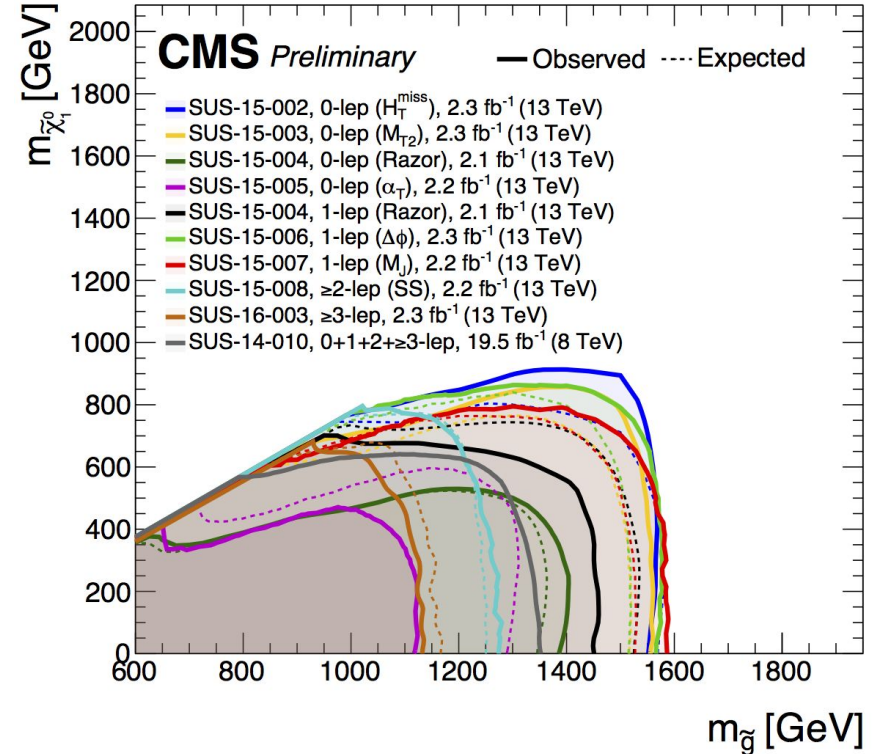
# Searches for gluino pair production

Run1 limits are already superseded!

$pp \rightarrow \tilde{g}\tilde{g}, \tilde{g} \rightarrow b\bar{b}\tilde{\chi}_1^0$  Moriond 2016



$pp \rightarrow \tilde{g}\tilde{g}, \tilde{g} \rightarrow t\bar{t}\tilde{\chi}_1^0$  Moriond 2016



# Searches for direct stop pair production

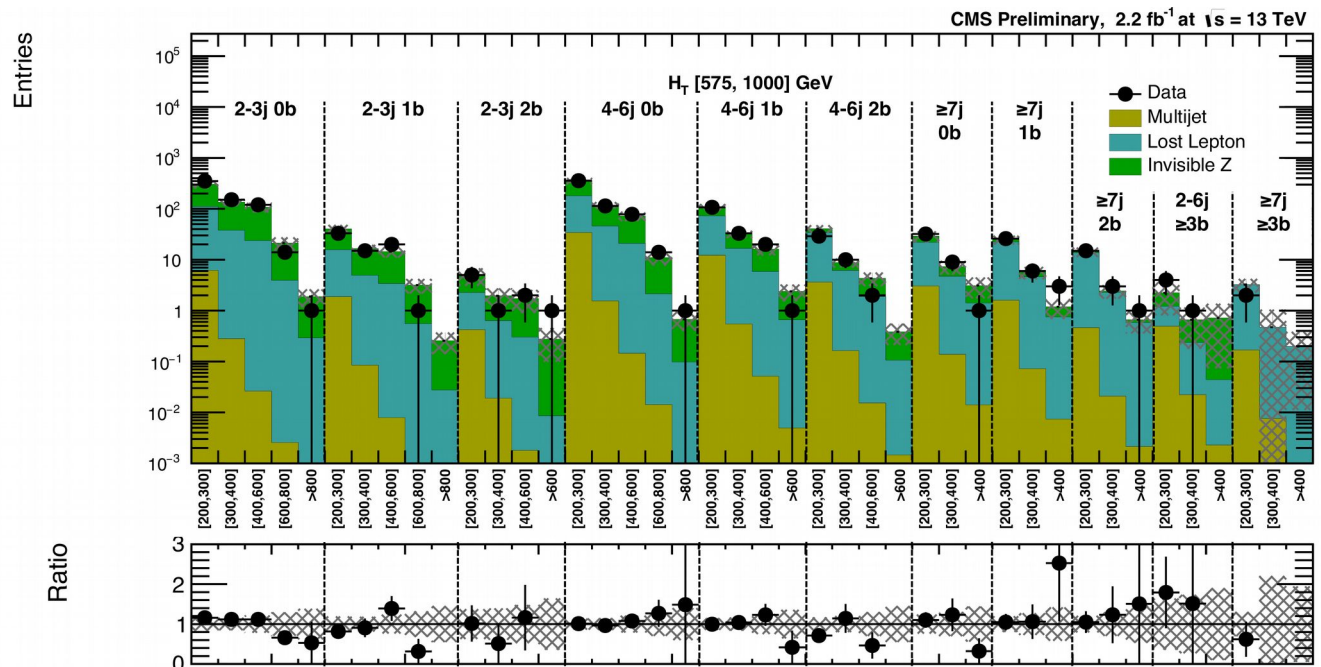
- Analysis exploiting the  $M_{T2}$  variable, designed to be robust against jet  $p_T$  mismeasurements:

$$M_{T2}(m_{\tilde{\chi}}) = \min_{\vec{p}_T^{\tilde{\chi}(1)} + \vec{p}_T^{\tilde{\chi}(2)} = \vec{p}_T^{\text{miss}}} \left[ \max \left( M_T^{(1)}, M_T^{(2)} \right) \right]$$

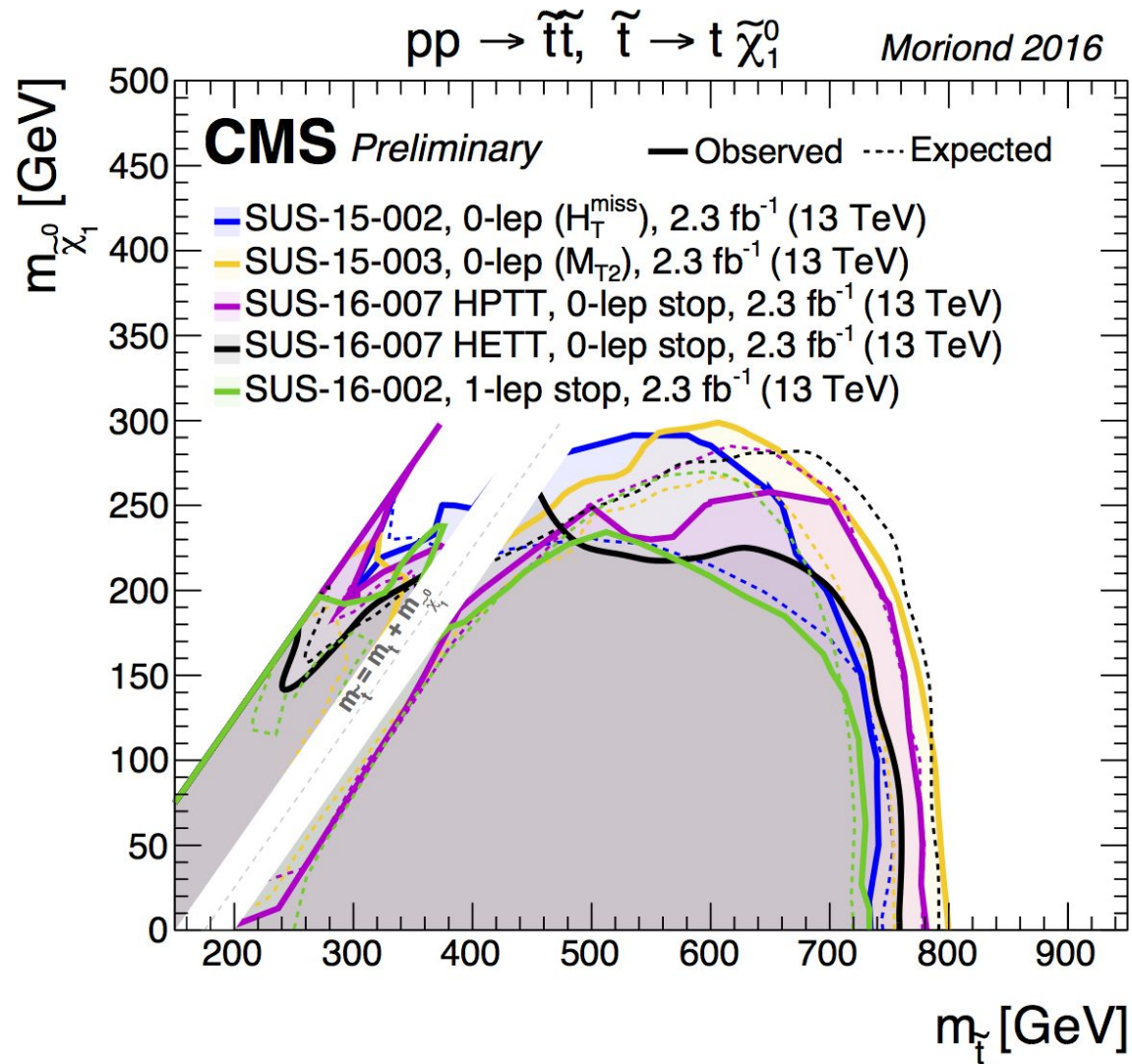
$$(M_T^{(i)})^2 = (m^{\text{vis}(i)})^2 + m_{\tilde{\chi}}^2 + 2 \left( E_T^{\text{vis}(i)} E_T^{\tilde{\chi}(i)} - \vec{p}_T^{\text{vis}(i)} \cdot \vec{p}_T^{\tilde{\chi}(i)} \right)$$

- Very powerful for signatures with two (pseudo-)jets;

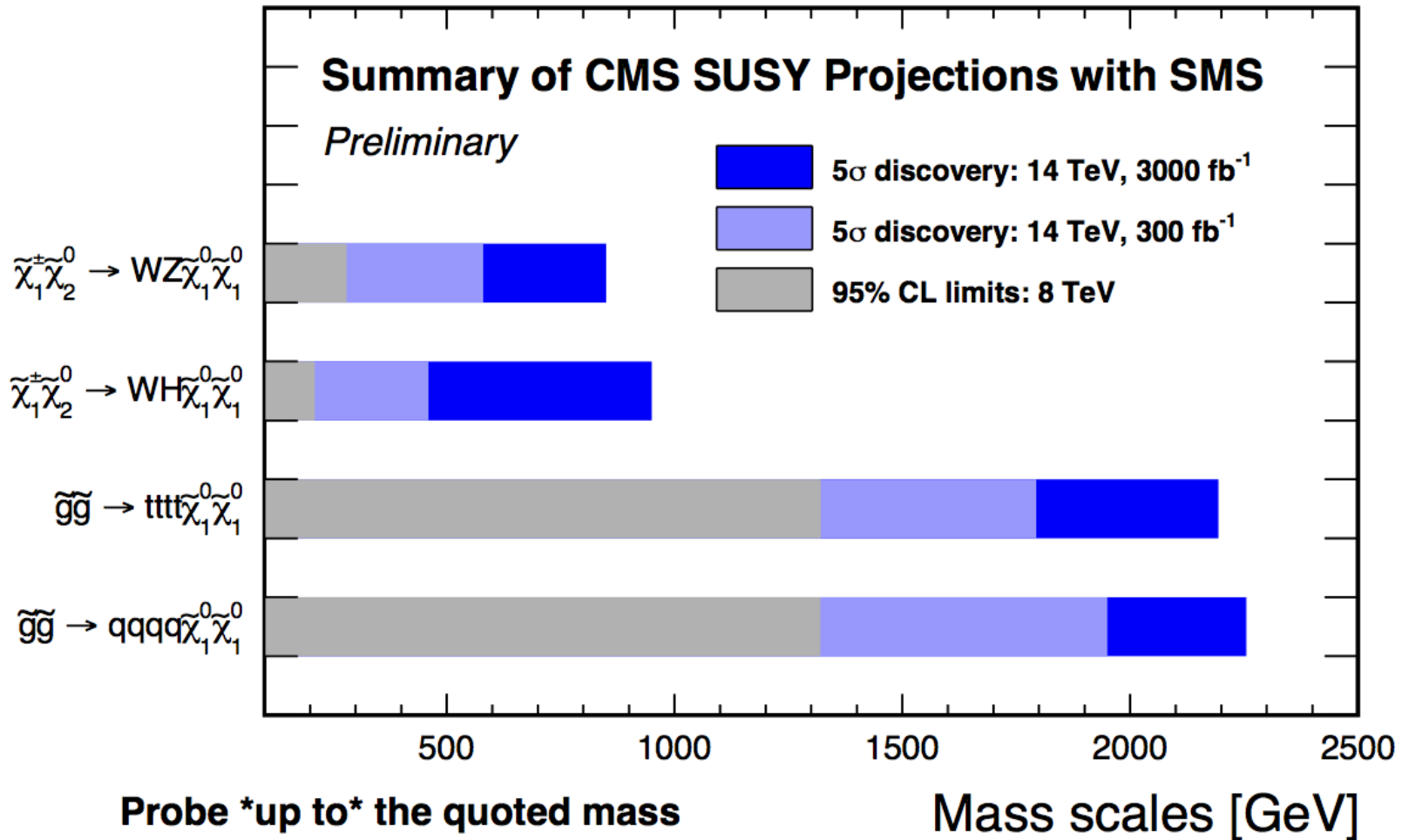
**CMS-SUS-15-003**



# Searches for direct stop pair production



# Ultimate LHC sensitivity



# Conclusions

- Run1 of the LHC has greatly extended our sensitivity to discoveries of SUSY signatures, I covered only a tiny part of all the searches;
- Unfortunately no evidence so far, the Natural SUSY paradigm is under a lot pressure;
- Still high hopes to find SUSY, but probably we will have to give up/relax something (minimality, naturalness, R-parity conservation, ... );
- LHC Run2 and HL LHC will extend our sensitivity even further, hopefully some of the excesses we have seen in the recent past will stay;
- And we should not forget that SUSY might show up somewhere else (Flavor Physics,  $(g-2)_\mu$ , direct DM searches...).