

Superstring theory and integration over the moduli space

Kantaro Ohmori

Hongo, The University of Tokyo

June, 24th, 2013@Nagoya

arXiv:1303.7299 [KO,Yuji Tachikawa]

Section 1

introduction

Simplified history of the superstring theory

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- '12: "Superstring Perturbation Theory Revisited"
[[E.Witten](#)]:

Superstring perturbation theory and supergeometry

The work of [KO,Tachikawa]

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- Formulation using supergeometry
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⇒ 2 examples where ambiguity does not appear
- $\mathcal{N} = 0 \subset \mathcal{N} = 1$ embedding [N. Berkovits,C. Vafa '94]
- Topological amplitudes
 - [I. Antoniadis, E. Gava, K.S. Narain, T.R. Taylor '94]
 - [M. Bershadsky, S. Cecotti, H. Ooguri, C. Vafa '94]

1

introduction

2

Supergeometry and Superstring

3

Reduction to integration over moduli

Section 2

Supergroup and Superstring

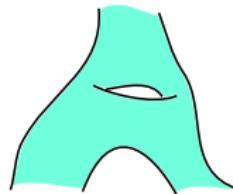
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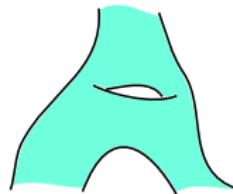
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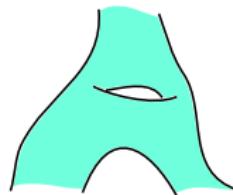
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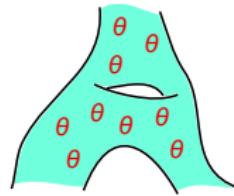
Superstring theory

Superstring theory

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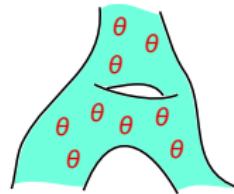
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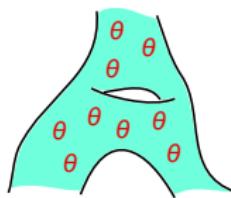
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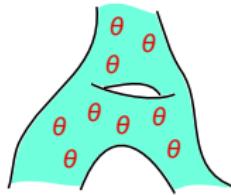
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- $\overset{?}{\Rightarrow} A = \int_{M_{bos,g}} dmd\bar{m} F'(m, \bar{m})$

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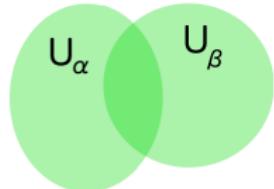
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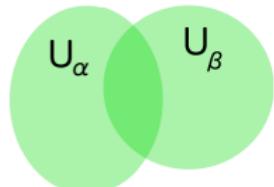
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- $M_{\text{red}} = \{(x_i, \dots, x_p)\} \hookrightarrow M$



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- $$D_\theta = \partial_\theta + \theta \partial_z = F(z|\theta)(\partial_{\theta'} + \theta' \partial_{z'})$$

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- $\theta \in H^0(\Sigma, T\Sigma_{\text{red}}^{*1/2})$
 \Rightarrow split SRS \leftrightarrow Riemann surface with spin str.

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Supermoduli space of SRS with punctures

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- $\dim \mathcal{M}_{\text{super}} = \Delta_e | \Delta_o$
 $= 3g - 3 + n_{NS} + n_R | 2g - 2 + n_{NS} + n_R / 2 \quad (g \geq 2)$

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- if we fix the behavior of Γ near boundary.



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- $\Sigma = (m|\eta_1, \dots, \eta_{\Delta_o})$

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- c.f. $D_\mu \phi^\dagger D^\mu \phi \ni A_\mu A^\mu \phi^\dagger \phi$

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- $\mathsf{F}_{\mathcal{V},g}(\mathbf{m}, \eta, \bar{\mathbf{m}}, \bar{\eta}) = \left\langle \prod_i \mathcal{V}_i \prod_{a=1}^{\Delta_e} \left(\int_{\Sigma_{\text{red}}} d^2 z b \mu^a \right) \prod_{b=1}^{\Delta_e} \left(\int_{\Sigma_{\text{red}}} d^2 z \bar{b} \tilde{\mu}^b \right) \right.$
 $\left. \prod_{\sigma=1}^{\Delta_o} \delta \left(\int_{\Sigma_{\text{red}}} d^2 z \beta \chi^\sigma \right) \prod_{\tau=1}^{\tilde{\Delta}_o} \delta \left(\int_{\Sigma_{\text{red}}} d^2 z \tilde{\beta} \tilde{\chi}^\tau \right) \exp(-S_\chi - S_{\tilde{\chi}} - S_{\chi\tilde{\chi}}) \right\rangle_{\mathbf{m}}$

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$$= \prod_\sigma^{\Delta_o} Y(p_\sigma)$$

$$Y(p_\sigma) = \delta(\beta(p_\sigma)) T_F(p_\sigma)$$

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 \Rightarrow Formulation of FMS
- Valid where $[\delta(z - p_\sigma)]$ spans odd deformations
 \Rightarrow Coordinates given by picture changing formalism is **not** global!

Integration over a supermanifold

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- M : $\dim 1|2$ supermanifold,

$M = U_1 \cup U_2, U_1 \ni (m|\eta_1, \eta_2), U_2 \ni (m'|\eta'_1, \eta'_2)$

$M = V_1 \coprod V_2, U_i \supset V_i$

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Integration over a supermanifold

- M : $\dim 1|2$ supermanifold,

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(with holomorphic factorization)

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- Infrared regularization

Section 3

Reduction to integration over moduli

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- $\omega = i_* \alpha \wedge P.D. [\mathcal{M}_{\text{spin}}]$
- ω does not depend on the locations of picture changing operators

Topological amplitudes in string theory

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- F-term of 4d effective field theory is related to topological string.

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- $\mathcal{A}_g = (g!)^2 F_g = \int_{\mathcal{M}_{bos}} \dots$

Vertex operators

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- Graviphoton: $\mathbf{V}_T = \Sigma \times \text{spacetime} \times (\text{ghost})$
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 $\mathbf{A}_{\pm\pm} : U(1)_R \text{ charge } (\pm 1, \pm 1)$
⇐ Coupling \mathbf{A} to $\mathcal{N} = (2, 2)$ supergravity:
 $\int_{\Sigma_{\text{red}}} \mathbf{A}_{\pm\pm} \chi^\mp \tilde{\chi}^\mp \rightarrow \chi = \chi^+ = \chi^-$

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- nonzero contribution
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- $\dim \Gamma = 6g - 6 | 6g - 6 \Rightarrow$ saturation η 's
- The correlation function does not depend of the choice of χ^σ

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- Special cases exist.
- Amplitudes which is equivalent topological amplitudes are the case.

Prospects

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- Another application of supergeometry

End.

Berkovits and Vafa embedding

[N.Berkovits, C.Vafa '94]

- Worldsheet supersymmetric model equivalent to bosonic string theory
⇒ Bosonic string : An (unstable) vacuum of superstring theory
- $\langle \mathbf{V}_1 \cdots \mathbf{V}_n \rangle_{\text{bos}} = \langle \mathbf{V}'_1 \cdots \mathbf{V}'_n \rangle_{\text{super}}$
- $\Leftrightarrow \int_{\mathcal{M}_{\text{bos}}} \cdots = \int_{\mathcal{M}_{\text{super}}} \cdots$

Supermoduli space of SRS with punctures

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- even deformation:

$$H^1(\Sigma_{\text{red}}, T\Sigma_{\text{red}}) + \text{positions of punctures}$$

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- $\dim \mathcal{M}_{\text{super}} = \Delta_e | \Delta_o$

$$= 3g - 3 + n_{\text{NS}} + n_{\text{R}} | 2g - 2 + n_{\text{NS}} + n_{\text{R}}/2 \quad (g \geq 2)$$

Superconformal transformation

- $\nu_f = f(z)(\partial_\theta - \theta\partial_z)$
- $V_g = g(z)\partial_z + 1/2\partial_z g\theta\partial_\theta$
- $G_n = \nu_{z^{n+1/2}}$
- $L_m = -V_{z^{m+1}}$

Superconformal transformation near R vertex

- $\nu_f = f(z)(\partial_\theta - z\theta\partial_z)$
- $V_g = z(g(z)\partial_z + 1/2\partial_z g\theta\partial_\theta)$
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- $\int d\eta_2 d\eta_1 f(\eta_1, \eta_2) = d$

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- R puncture: $\{z = z_1\} \subset \Sigma$: submanifold of Σ
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- $\theta' = \sqrt{z - z_1}\theta \Rightarrow D_\theta^* = \sqrt{z}(\partial_{\theta'} + \theta'\partial_z)$
 θ' : antiperiodic around z_1

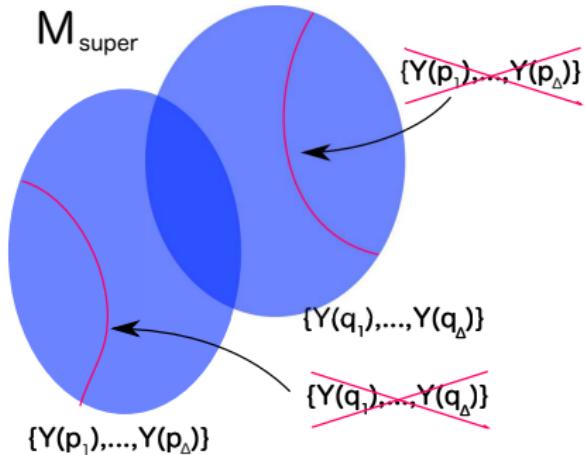
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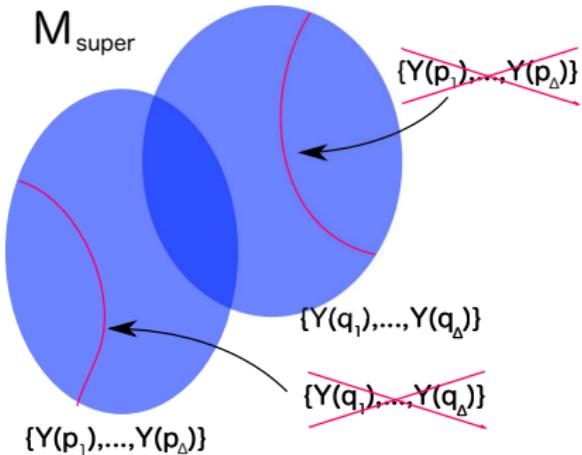
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is **not** global!



Subtlety for picture changing formalism

- Valid where $[\delta(p_\sigma)]$ spans odd deformations
⇒ Coordinates given by picture changing formalism
is **not** global!



- movement of PCOs ⇒ BRST exact term \Leftrightarrow exact form on patches