Gauge Symmetry

and

Functional RG

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Introduction and Summary

♢ Functional Renormalization Group (FRG): A convincing nonperturbative approach to field theories. ⇒ Dynamics described by RG flow of couplings $g(\Lambda)$ in theory space.

The momentum cutoff Λ often conflicts with symmetries.

Gauge symmetry ?

Even in the presence of Λ, gauge symmetry realized as a quantum symmetry by imposing Ward-Takahashi (WT) identity for Wilson action S_Λ

$$\Sigma_\Lambda \sim \left(\partial S_\Lambda/\partial\phi^A\right)\delta\phi^A - \mathrm{Str}\left(\partial\delta\phi/\partial\phi\right) = 0$$

(1) symmetry tr. $\delta\phi^A$ depend on S_Λ
(2) non-trivial Jacobian factor in functional measure $\mathcal{D}\phi$

$\Sigma_\Lambda = 0$ defines gauge invariant subspace in theory space.

⇒ Expression of symmetry in FRG.

Two actions: Wilson and Effective average actions

	Wilson action	Effective average action
	S_Λ	Γ_Λ
Diagrams	connected	1PI
Gauge symmetry	WT id.	'mod. Slavnov-Taylor id.'
RG flow eq.	Polchinski eq.	Wetterich eq.

$$\Gamma_{I,\Lambda}[\Phi] = S_{I,\Lambda}[\phi] - \frac{1}{2}(\Phi - \phi)\Delta_H^{-1}(\Phi - \phi)$$

$$\Phi - \phi = -\Delta_H \frac{\partial^l S_{I,\Lambda}}{\partial \phi} = -\Delta_H \frac{\partial^l \Gamma_{I,\Lambda}}{\partial \Phi}$$

where Δ_H is the high momentum propagator that allows particles above Λ to propagate.

Our strategy
- Algebraic study of gauge symmetry with the Wilson action.
- Wetterich eq.: one-loop calculation with the full propagator.

We study QED.

♢ Two fundamental equations we have to solve:

i) WT identity $\Sigma_\Lambda = 0$

ii) RG flow eq. $\partial_t \Gamma_\Lambda = \mathrm{Str}\,(\partial_t R)\left[\partial^2\Gamma_\Lambda/\partial\Phi\partial\Phi + R\right]^{-1}$ $(\Lambda\partial_\Lambda = \partial_t)$

We discuss an exact solution to $\Sigma_\Lambda = 0$ for suitably truncated Wilson action in QED.

- Need to introduce higher dimensional interactions with form factors (momentum dependent 4-fermi couplings).
- Take account of full momentum dependence in WT and flow eq. without using derivative (momentum) expansion.

♢ Results

- Exact evaluation of photon 2-point functions
- Relation which corresponds to $Z_1 = Z_2$
- Relations between form factors in 4-fermi couplings and photon propagator

♢ Plan of the talk

[1] Derivation of the WT identity

[2] WT identity in QED

[3] Exact solution to WT identity

[4] Momentum dependent flow eq.

[5] Outlook

Derivation of the WT identity

◇ Consider a gauge-fixed theory described by

$$\mathcal{S}[\varphi] \;=\; \frac{1}{2}\varphi \cdot D \cdot \varphi + \mathcal{S}_I[\varphi], \qquad \varphi \cdot D \cdot \varphi = \int_p \varphi^A(-p) D_{AB}(p) \varphi^B(p).$$

We rewrite its partition function as

$$\begin{aligned} \mathcal{Z}_\varphi[J] &= \int \mathcal{D}\varphi \exp\left(-\mathcal{S}[\varphi] + J \cdot \varphi\right) \\ &= N_J \int \mathcal{D}\phi \exp\left[-\frac{1}{2}\phi \cdot K^{-1} D \cdot \phi + J \cdot K^{-1}\phi\right] \\ &\times \int \mathcal{D}\chi \exp\left[-\frac{1}{2}\chi \cdot (1-K)^{-1} D \cdot \chi - \mathcal{S}_I[\phi + \chi]\right] \end{aligned}$$

where $\chi = \varphi - \phi$. We have introduced an IR cutoff Λ through a positive function that behaves as

$$K(p) \equiv K(p^2/\Lambda^2) \to \quad \begin{cases} 1 & (p^2 < \Lambda^2) \\ 0 & (p^2 > \Lambda^2) \end{cases}$$

For cut-off function, we take e.g. $K(p) = e^{-p^2/\Lambda^2}$.

The Wilson action is defined by

$$S_\Lambda[\phi] = \frac{1}{2}\phi^A (K^A)^{-1}\ D_{AB}\phi^B + S_{I,\Lambda}[\phi]\ ,$$

where the interaction part is given by a functional integral

$$\begin{aligned} \exp\Big[-S_{I,\Lambda}[\phi]\Big] &= \int \mathcal{D}\chi \exp\Big[-\frac{1}{2}\chi\cdot(\Delta_H)^{-1}\cdot\chi - \mathcal{S}_I[\phi+\chi]\Big] \\ \Delta_H &= (1-K)D^{-1} \end{aligned}$$

The partition function for the Wilson action,

$$Z_\phi[J] = \int \mathcal{D}\phi \exp\Big[-S_\Lambda[\phi] + J \cdot K^{-1}\phi\Big] \ ,$$

is related to that for the original one by

$$\mathcal{Z}_\varphi[J] \quad = \quad N_J Z_\phi[J] \ ,$$

where the normalization factor is given by

$$N_J \quad = \quad \exp\frac{1}{2}\left[-(-)^{\epsilon(J_A)} J_A \left(\frac{1-K}{K}\right)^A \left(D^{-1}\right)^{AB} J_B\right] \ .$$

♢ We define the WT operator

$$\Sigma_\Lambda[\phi] = K^A \left[\frac{\partial^r S_\Lambda}{\partial \phi^A} \delta\phi^A - (-)^{\epsilon_A} \frac{\partial^l \delta\phi^A}{\partial \phi^A} \right],$$

$$\Sigma_\Lambda[\phi] = 0$$

signals for the presence of BRST (quantum) symmetry.

To find $\delta\phi$, take functional average of the WT op. for the original theory with standard BRST symmetry

$$\begin{aligned} \Sigma[\varphi] &= \frac{\partial^r \mathcal{S}}{\partial \varphi^A} \delta\varphi^A - (-)^{\epsilon_A} \underbrace{\frac{\partial^l \delta\varphi^A}{\partial \varphi^A}}_{=0} \\ \delta\varphi^A &= R^A{}_B \varphi^B \, . \end{aligned}$$

where $R^A{}_B$ are field independent coeffients, and $\delta\varphi^A$ stand for *classical* (conventional) BRST transformations for linear symmetry. Use them as a "seed" for quantum symmetry.

Through relations

$$\int \mathcal{D}\varphi\ \delta\varphi \exp\left(-\mathcal{S}[\varphi] + J\cdot\varphi\right) = N_J \int\ \mathcal{D}\phi\ K^{-1}\delta\phi\ \ \exp\left(-S_\Lambda[\Phi] + J\cdot K^{-1}\Phi\right)$$

$$\int \mathcal{D}\varphi\ \Sigma[\varphi] \exp\left(-\mathcal{S}[\varphi] + J\cdot\varphi\right) = N_J \int\ \mathcal{D}\phi\ \Sigma_\Lambda[\phi]\ \ \exp\left(-S_\Lambda[\Phi] + J\cdot K^{-1}\Phi\right)\ ,$$

we find

$$\delta\phi^A = R^A{}_B\left[\phi^B\right]_\Lambda, \qquad \left[\phi^B\right]_\Lambda = \phi^B - (\Delta_H)^{BC}\frac{\partial^l S_{I,\Lambda}}{\partial\phi^C}$$

where $\left[\phi^A\right]_\Lambda$ are "composite operators" for fields ϕ^A. They obey RG flow equations:

$$\partial_t\mathcal{O}_\Lambda[\phi] = -\frac{\partial^r S_{I,\Lambda}}{\partial\phi^A}(\partial_t\Delta_H)^{AB}\frac{\partial^l\mathcal{O}_\Lambda}{\partial\phi^B} + \frac{1}{2}(-)^{\epsilon_A(1+\epsilon_\mathcal{O})}\left(\partial_t\Delta_H\right)^{AB}\frac{\partial^l\partial^r\mathcal{O}_\Lambda}{\partial\phi^B\partial\phi^A}\ .$$

We also obtain general expression of the WT op. for linear gauge symmetry

$$\Sigma_\Lambda[\phi] = K^A\left\{\frac{\partial^r S_\Lambda}{\partial\phi^A}R^A{}_B\left[\phi^B\right]_\Lambda + (-)^{\epsilon_A}R^A{}_B\left(\Delta_H\right)^{BC}\frac{\partial^l\partial^r S_{I,\Lambda}}{\partial\phi^C\partial\phi^A}\right\}\ .$$

WT identity for QED

♢ Consider the Wilson action $S_\Lambda[\phi] = S_{0,\Lambda} + S_{I,\Lambda}$ for the fields $\phi^A = (a_\mu,\ \bar\psi_{\hat\alpha},\ \psi_\alpha,\ c,\ \bar c)$.

The kinetic part of the Wilson action is given by

$$
\begin{aligned}
S_{0,\Lambda} &= \frac{1}{2}(K^A)^{-1} Z_A \phi^A D_{AB} \phi^B \\
&= \int_p K^{-1}(p) \left[\frac{Z_3}{2} a_\mu(-p) p^2 \left\{ \delta_{\mu\nu} - \left(1 - (Z_3 \xi_0)^{-1}\right) \frac{p_\mu p_\nu}{p^2} \right\} a_\nu(p) + \bar c(-p) i p^2 c(p) \right] \\
&+ \int_p K^{-1}(p) Z_2 \bar\psi(-p) \not{p} \psi(p) \ ,
\end{aligned}
$$

where we have introduced the renormalization constants, $Z_2,\ Z_3$. The classical BRST tr.

$$
\begin{aligned}
\delta_{cl}\ a_\mu(p) &= -i p_\mu c(p), & \delta_{cl}\ \bar c(p) &= \xi_0^{-1} p_\mu a_\mu(p) \\
\delta_{cl}\ \psi(p) &= -i\ e_0 \int_q \psi(q) c(p-q), & \delta_{cl}\ \bar\psi(-p) &= i\ e_0 \int_q \bar\psi(-q) c(q-p) \ ,
\end{aligned}
$$

fix the coefficients $R^A{}_B$ in our general formula for quantum symmetry. Here, $e_0,\ \xi_0$ are gauge coupling and gauge fixing parameters which are constants.

The WT operator for QED is constructed as

$$
\begin{aligned}
\Sigma_\Lambda[\phi] \quad = \quad & \int_p \Big\{ \frac{\partial S_\Lambda}{\partial a_\mu(p)}(-ip_\mu)c(p) + \frac{\partial^r S_\Lambda}{\partial \bar{c}(p)}\xi_0^{-1} p_\mu a_\mu(p) \Big\} \\
& -i\ e_0 \int_{p,\ q} \Big\{ \frac{\partial^r S_\Lambda}{\partial \psi_\alpha(q)} \frac{K(q)}{K(p)} \psi_\alpha(p) - \frac{K(p)}{K(q)} \bar{\psi}_{\hat{\alpha}}(-q) \frac{\partial^l S_\Lambda}{\partial \bar{\psi}_{\hat{\alpha}}(-p)} \Big\} c(q-p) \\
& -i\ e_0 \int_{p,\ q} U_{\beta\hat{\alpha}}(-q,p) \Big\{ \frac{\partial^l S_\Lambda}{\partial \bar{\psi}_{\hat{\alpha}}(-p)} \frac{\partial^r S_\Lambda}{\partial \psi_\beta(q)} - \frac{\partial^l \partial^r S_\Lambda}{\partial \bar{\psi}_{\hat{\alpha}}(-p) \partial \psi_\beta(q)} \Big\} c(q-p)\ ,
\end{aligned}
$$

where

$$
U(-q,p) = Z_2^{-1} \left[K(q) \frac{1-K(p)}{\not{p}} - K(p) \frac{1-K(q)}{\not{q}} \right]
$$

Exact solution to WT identity

$$S_{I,\Lambda}[\phi] = \Gamma_{I,\Lambda}[\Phi] + \frac{1}{2}(\Phi - \phi) \cdot (1 - K)^{-1} D \cdot (\Phi - \phi)$$

To construct interaction part $S_{I,\Lambda}[\phi]$, we first specify its 1PI part, namely effective average action $\Gamma_{I,\Lambda}[\Phi]$, imposing for simplicity chiral symmetry on the fermionic sector. We introduce some form factors in 4-fermi interactions:

$$\begin{aligned}
\Gamma_{I,\Lambda}[\Phi] = &\int_p \left[\frac{Z_3}{2} A_\mu(-p) \mathcal{M}_{\mu\nu}(p) A_\nu(p) + Z_2 \sigma(p) \bar{\Psi}(-p) \not{p} \Psi(p) \right] \\
&- e Z_2 Z_3^{1/2} \int_{p,q} \bar{\Psi}(-p) \not{A}(p-q) \Psi(q) + \frac{Z_2^2}{2\Lambda^2} \int_{p_1,\cdots,p_4} (2\pi)^4 \delta^4(p_1 + p_2 + p_3 + p_4) \\
&\times \left\{ h_S(s,u) \left[\left(\bar{\Psi}\Psi\right)^2 - \left(\bar{\Psi}\gamma_5\Psi\right)^2 \right] + h_V(s,u) \left[\left(\bar{\Psi}\gamma_\mu\Psi\right)^2 + \left(\bar{\Psi}\gamma_5\gamma_\mu\Psi\right)^2 \right] \right. \\
&\left. + \frac{1}{\Lambda^2} (p_1 + p_4)_\mu (p_2 + p_3)_\nu h_{V'}(s,u) \left[\left(\bar{\Psi}\gamma_\mu\Psi\right) \left(\bar{\Psi}\gamma_\nu\Psi\right) + \left(\bar{\Psi}\gamma_5\gamma_\mu\Psi\right) \left(\bar{\Psi}\gamma_5\gamma_\nu\Psi\right) \right] \right\}
\end{aligned}$$

Here $\quad \mathcal{M}_{\mu\nu}(p) = P^T_{\mu\nu} \mathcal{T}(p) + P^L_{\mu\nu} \mathcal{L}(p), \qquad P^T = \delta_{\mu\nu} - p_\mu p_\nu / p^2, \quad P^L = p_\mu p_\nu / p^2$

and s, u are Mandelstam variables.

$S_{I,\Lambda}[\phi]$ is constructed by using the Legendre transformation :

$$
\begin{aligned}
S_{I,\Lambda}[\phi] &= \Gamma_{I,\Lambda}[\Phi] + \frac{1}{2}(\Phi - \phi)\cdot(1-K)^{-1}D\cdot(\Phi-\phi) \\
&= \Gamma_{I,\Lambda}[\phi] + \frac{Z_3}{2}\int_p a_\mu(-p)\left[\sum_{n=1}(-)^n\Big[\big(\mathcal{M}\Delta_H\big)^n\Big]_{\mu\lambda}(p)\ \mathcal{M}_{\lambda\nu}(p)\right]a_\nu(p) \\
&\quad -eZ_2Z_3^{1/2}\int_{p,q}\sum_{n=1}(-)^n\Big[\big(\mathcal{M}\Delta_H\big)^n\Big]_{\mu\nu}(p-q)a_\nu(p-q)\left(\bar{\psi}\gamma_\mu\psi\right) \\
&\quad -\frac{Z_2^2e^2}{2}\int_{p_1,\cdots,p_4}(2\pi)^4\delta^4(p_1+p_2+p_3+p_4)\left(\bar{\psi}\gamma_\mu\psi\right)\left(\bar{\psi}\gamma_\nu\psi\right)\left(\mathbf{\Delta}_G\right)_{\mu\nu}(p_1+p_2)
\end{aligned}
$$

where additional terms to $\Gamma_{I,\Lambda}$ are 1P reducible contributions. $\left(\mathbf{\Delta}_G\right)_{\mu\nu}(p) = Z_3^{-1}[P^T_{\mu\nu}T(p) + P^L_{\mu\nu}L(p)]$ is full photon propagator constructed with photon 2-point functions

$$
T(p) = \frac{1-K}{p^2+(1-K)\mathcal{T}(p)}, \qquad L(p) = \frac{\xi(1-K)}{p^2+\xi(1-K)\mathcal{L}(p)}, \qquad \xi = Z_3\xi_0
$$

Substitute $S_\Lambda = S_{0,\Lambda} + S_{I,\Lambda}$ into the WT identity $\Sigma_\Lambda[\phi] = 0$, which can be expanded in polynomial of ϕ^A. Consider two terms $a_\mu \times c$ and $\bar\psi \times \psi \times c$ in this expansion. For simplicity, we assume $\sigma(p) = 0$ for fermionic 2-point function. From $a_\mu(p) \times c(-p)$ term, we have

$$Z_3 p_\nu \mathcal{L}(p) = -e_0\, e Z_2 Z_3^{1/2} \int_q \mathrm{Tr}[U(-p-q, q)\gamma_\nu]$$

From $\bar\psi(p) \times \psi(-q) \times c(q-p)$ term, we have second WT relation

$$\begin{aligned}
&\left(e_0 Z_2 - e Z_2 Z_3^{1/2}\right)(\not{p} - \not{q}) - 2e_0 Z_2^2 \int_k \Bigg[\frac{1}{\Lambda^2}\Big\{\big(h_S - 2h_V\big)[k^2, (p+q)^2] \\
&-2h_V[(p+q)^2, k^2]\Big\} + e^2 T(k^2) + \frac{1}{\Lambda^4}\Big\{2(p-q)^2 h_{V'}[k^2, (p+q)^2] \\
&+k^2 h_{V'}[(p+q)^2, k^2]\Big\}\Bigg] U(-q-k, p+k) \\
&-e_0 Z_2^2 \int_k \left[\frac{2}{\Lambda^4} h_{V'}[(p+q)^2, k^2] + e^2 \frac{1}{k^2}\{T(k^2) - L(k^2)\}\right] \not{k}\, U(-q-k, p+k) \not{k} = 0\ .
\end{aligned}$$

This constraint splits into two conditions: $\mathrm{constant} \times (\not{p} - \not{q})$ and one-loop part. They should separately vanish. The first one gives

$$e_0 = Z_3^{1/2} e \ .$$

This corresponds to the well-known WT relation in the standard realization of gauge symmetry in QED: $Z_1 = Z_2$ for $Z_1 = Z_2 Z_3^{1/2} Z_e$ with $e = Z_e \, e_0$.

On the other hand, one-loop part gives

$$\frac{1}{\Lambda^2}\left\{\left(h_S - 2h_V\right)[k^2, (p-q)^2] - 2h_V[(p-q)^2, k^2]\right\}$$
$$= e^2\left\{T[(p-q)^2] - L[(p-q)^2]\right\} - \frac{e^2}{2}\left\{T(k^2) + L(k^2)\right\}$$
$$\frac{1}{\Lambda^4} h_{V'}[(p+q)^2, k^2] = -\frac{e^2}{2k^2}\left\{T(k^2) - L(k^2)\right\}$$

These are relations between 4-fermi interactions and photon propagator.

Note that derivative expansion will give , $1 - Z_3^{1/2} e/e_0 \simeq [h_S(0,0) - 4h_V(0,0)]$.

Remarkably, longitudinal component of photon 2-point function $\mathcal{L}$ can be evaluated exactly for a specific cutoff function $K(p) = e^{-p^2/\Lambda^2}$ using some formula for the modified Bessel functions:

$$\int_0^\pi d\theta\ e^{2pk\cos\theta}\sin^2\theta = \frac{\pi}{2pk} I_1(2pk), \qquad \int_0^\infty dk e^{-k^2} I_1(2pk) = \frac{p}{2}\ {}_1F_1(1,2;p^2)$$

we obtain

$$\mathcal{L}(p^2) = -e^2 \frac{\Lambda^2}{2\pi^2 \bar{p}^4}\left[1 - \exp(-\bar{p}^2/2) - \bar{p}^2\left(1 - \frac{1}{2}\exp(-\bar{p}^2/2)\right)\right]$$

where we have used $e_0 = Z_3^{1/2} e$ to eliminate e_0, and $\bar{p}^2 = p^2/\Lambda^2$. To fix transverse part $\mathcal{T}$, we use RG flow equations.

Momentum dependent flow equations

◇ For photon 2-point functions $\propto e^2$ in RG flow equation we have

$$\frac{Z_3}{2}\int_p A_\mu(-p)\left[P^T_{\mu\nu}\left\{2\eta_A p^2 - 2\mathcal{T}(p)\right\} + (-2)P^L_{\mu\nu}\mathcal{L}(p)\right]A_\nu(p)$$
$$= -e^2 Z_3 \int_{p,q} 2K'(q)\,(1-K(p+q))^2\,\frac{1}{(p+q)^2}\mathrm{Tr}\left[\not{A}(-p)(\not{p}+\not{q})\not{A}(p)\not{q}\right]$$

Rhs can be exactly evaluated to give

$$\mathrm{rhs} = -Z_3\frac{e^2}{2\pi^2}\int_p A_\mu(-p)\left[P^T_{\mu\nu}\frac{1}{4p^4}\left\{4-\left(4+2p^2-p^4\right)\exp(-p^2/2)\right\}\right.$$
$$\left. -P^L_{\mu\nu}\frac{1}{p^4}\left\{1-p^2-\left(1-\frac{p^2}{2}\right)\exp(-p^2/2)\right\}\right]A_\nu(p)$$

Since p^2 term in transverse part here generates well-known anomalous dimension for photon field $\eta_A = e^2/12\pi^2$, we subtract it to find $\mathcal{T}$

$$
\begin{aligned}
\mathcal{T}(p^2) - \eta_A p^2 &= \frac{e^2}{8\pi^2 p^4}\Big\{4 - \big(4 + 2p^2 - p^4\big)\exp(-p^2/2)\Big\} \\
\mathcal{T}(p^2) &= \frac{\Lambda^2 e^2}{8\pi^2 \bar{p}^4}\Big\{4 + \frac{2\bar{p}^6}{3} - \big(4 + 2\bar{p}^2 - \bar{p}^4\big)\exp(-\bar{p}^2/2)\Big\}\,.
\end{aligned}
$$

$\mathcal{L}$ appeared here is exactly the same as the one obtained by WT identity.

In this way, we fix photon 2-point functions.

Note that the same constant mass term appears in both $\mathcal{T}$ and $\mathcal{L}$

$$
\mathcal{T} = \mathcal{L} = \frac{3e^2}{16\pi^2}\Lambda^2 + \mathcal{O}(\bar{p}^2)
$$

Outlook

♢ $\Sigma_\Lambda = 0$ (almost) determines S_Λ.

All 4-fermi couplings expressed in terms of e^2 and photon 2-point functions ?

$\Leftarrow$ careful analysis of flow eq.

♢ Exact evaluation of photon 2-point functions is interesting but only possible in QED with simplified fermionic sector.

$\Rightarrow$ For more complicated cases such as YM theory, need to develop suitable approximation method which replaces derivative expansion.

Taking account of momentum dependence in WT identity and RG flow eq. will give new insights into FRG !