# The Jet Energy Profile: A BSM Analysis Tool

R. Sekhar Chivukula SCGT 15, March 2015





## The Jet Energy Profile: A BSM Analysis Tool

- Dijets at the LHC: searching for new resonances
- Benchmark Resonances
- The Jet Energy Profile (JEP)
- Measuring the JEP Examining New Resonances
- Conclusions

Based on RSC, EHS, & N. Vignaroli arXiv: 1412.3094

## The LHC Produces Jets Copiously



New particles decaying to dijets can be produced to very large masses!

# LHC Dijet Data



ATLAS, arXiv:1407.1376

CMS, arXiv:1501.04198

### Possible Dijet Resonances

initial state	J	$SU(3)_C$	$SU(2)_L$	$U(1)_Y$	$ Q_e $	B
QQ	0	$\overline{f 3}\oplus f 6$	$1 \oplus 3$	$\frac{1}{3}$	$\frac{4}{3}, \frac{2}{3}, \frac{1}{3}$	$\frac{2}{3}$
QU	1	$\overline{f 3}\oplus f 6$	2	$\frac{5}{6}$	$\frac{4}{3}, \frac{1}{3}$	$\frac{2}{3}$
QD	1	$\overline{f 3}\oplus f 6$	2	$-\frac{1}{6}$	$\frac{2}{3}, \frac{1}{3}$	$\frac{2}{3}$
	0	$\overline{f 3}\oplus f 6$	1	$\frac{4}{3}$	$\frac{4}{3}$	$\frac{2}{3}$
DD	0	$\overline{f 3}\oplus f 6$	1	$-\frac{2}{3}$	$\frac{2}{3}$	$\frac{2}{3}$
UD	0	$\overline{3} \oplus 6$	1	$\frac{1}{3}$	$\frac{1}{3}$	$\frac{2}{3}$
QA	$(\frac{1}{2}, \frac{3}{2})$	$3 \oplus ar{6} \oplus 15$	2	$\frac{1}{6}$	$\frac{2}{3}, \frac{1}{3}$	$\frac{1}{3}$
UA	$\frac{1}{2}, \frac{3}{2}$	$3 \oplus \mathbf{ar{6}} \oplus 15$	1	$\frac{2}{3}$	$\frac{2}{3}$	$\frac{1}{3}$
DA	$\frac{1}{2}, \frac{3}{2}$	$3\oplusar{6}\oplus15$	1	$\frac{1}{3}$	$\frac{1}{3}$	$\frac{1}{3}$
AA	0, 1, 2	$1 \oplus 8 \oplus 8 \oplus 10 \oplus \overline{10} \oplus 27$	1	0	0	0
$Q\bar{Q}$	(1)	$1 \oplus 8$	$1 \oplus 3$	0	1, 0	0
$Q\bar{U}$	0	$1 \oplus 8$	2	$-\frac{1}{2}$	1,0	0
$Q\bar{D}$	0	$1 \oplus 8$	2	$\frac{1}{2}$	1,0	0
$U\bar{U}, \ D\bar{D}$	(1)	$1 \oplus 8$	1	0	0	0
$U\bar{D}$	1	$1 \oplus 8$	1	1	1	0

(NB: Colored resonances cannot decay solely to leptons)

Han, Lewis, Liu arXiv:1010.4309

If a new resonance is discovered, and decays only to dijets, can we determine what it is and how it decayed?

To illustrate: consider three benchmark possibilities

QQ Resonances: A Coloron

### Color Octet Vector Resonances

#### Gauge bosons from extended color groups:

Classic Axigluon: P.H. Frampton and S.L. Glashow, Phys. Lett. B 190, 157 (1987). Topgluon: C.T. Hill, Phys. Lett. B 266, 419 (1991). Flavor-universal Coloron: R.S. Chivukula, A.G. Cohen, & E.H. Simmons, Phys. Lett. B 380, 92 (1996). Chiral Color with  $g_L \neq g_R$ : M.V. Martynov and A.D. Smirnov, Mod. Phys. Lett. A 24, 1897 (2009). New Axigluon: P.H. Frampton, J. Shu, and K. Wang, Phys. Lett. B 683, 294 (2010).

#### Similar color-octet states:

KK gluon: H. Davoudiasl, J.L. Hewett, and T.G. Rizzo, Phys. Rev. D63, 075004 (2001) B. Lillie, L. Randall, and L.-T. Wang, JHEP 0709, 074 (2007). Techni-rho: E. Farhi and L. Susskind, Physics Reports 74, 277 (1981).

#### More exotic colored states:

Color sextets, colored scalars, low-scale scale string resonances... T. Han, I. Lewis, Z. Liu, JHEP 1012, 085 (2010). Coloron Models: Gauge Sector



SU(3)<sub>1</sub> x SU(3)<sub>2</sub> color sector with  $M^2 = \frac{u^2}{4} \begin{pmatrix} h_1^2 & -h_1h_2 \\ -h_1h_2 & h_2^2 \end{pmatrix}$ unbroken subgroup: SU(3)<sub>1+2</sub> = SU(3)<sub>QCD</sub>

 $h_1 = \frac{g_s}{\cos\theta} \qquad h_2 = \frac{g_s}{\sin\theta}$ 

gluon state:  $G^A_\mu = \cos \theta A^A_{1\mu} + \sin \theta A^A_{2\mu}$ couples to:  $g_S J^\mu_G \equiv g_S (J^\mu_1 + J^\mu_2)$   $M_G = 0$ 

coloron state:  $C^A_\mu = -\sin\theta A^A_{1\mu} + \cos\theta A^A_{2\mu}$   $M_C = \frac{u}{\sqrt{2}}\sqrt{h_1^2 + h_2^2}$ couples to:  $g_S J^\mu_C \equiv g_S (-J^\mu_1 \tan\theta + J^\mu_2 \cot\theta)$ 

Quarks' SU(3)<sub>1</sub> x SU(3)<sub>2</sub> charges impact phenomenology

### Matter Couplings

SU(3)1	SU(3) <sub>2</sub>	model	pheno.	
Benchmark	(t,b) <sub>L</sub> q <sub>L</sub> t <sub>R</sub> ,b <sub>R</sub> q <sub>R</sub>	coloron	dijet	
<b>Q</b> R	(t,b) <sub>L</sub> q <sub>L</sub> t <sub>R</sub> ,b <sub>R</sub>			
t <sub>R</sub> ,b <sub>R</sub>	(t,b) <sub>L</sub> q <sub>L</sub> q <sub>R</sub>			
q∟	(t,b) <sub>L</sub> t <sub>R</sub> ,b <sub>R</sub> q <sub>R</sub>			
q∟ t <sub>R</sub> ,b <sub>R</sub>	(t,b) <sub>L</sub> q <sub>R</sub>	new axigluon	dijet, At <sub>FB,</sub> FCNC	
<b>q</b> L <b>Q</b> R	(t,b) <sub>L</sub> t <sub>R</sub> ,b <sub>R</sub>	topgluon	dijet, tt, bb, FCNC, R <sub>b</sub>	
t <sub>R</sub> ,b <sub>R</sub> q <sub>R</sub>	(t,b)∟ q∟	classic axigluon	dijet, At <sub>FB</sub>	
q <sub>L</sub> t <sub>R</sub> ,b <sub>R</sub> q <sub>R</sub>	(t,b)∟			

q = u,d,c,s

Estimated LHC Reach: Signal Jet Selection

- p<sub>T</sub> > 30 GeV, |η| < 2.5</li>
- t-channel rejection:  $|\Delta \eta| < 1.3$
- Inspired by CMS cuts, arXiv:1501.04198
- Acceptance rates: 50% 60% for benchmark models

### 14 TeV LHC Reach: Flavor Universal Coloron



# Qg Resonances: An "Excited" Quark

# Excited Quarks & Heavy Vector Partners

- Composite Quark Models
- Composite Higgs
- Extra-Dimensional Models

### Benchmark: Doublet Partner of First-Generation

$$\mathcal{L}_{int} = \frac{1}{2\Lambda} \bar{q}_R^* \sigma^{\mu\nu} \left[ g_S f_S \frac{\lambda^a}{2} G^a_{\mu\nu} + g f \frac{\tau}{2} \cdot \mathbf{W}_{\mu\nu} + g' f' \frac{Y}{2} B_{\mu\nu} \right] q_L + \text{H.c.}$$

$$\Gamma(q^* \to qg) = \frac{1}{3} \alpha_S f_S^2 \frac{m_{q*}}{\Lambda^2}$$

Following ATLAS & CMS, take:  $\Lambda = m_{q^*}$ 

Baur, Spira, Zerwas: PRD 42 (1990) 815.

### 14 TeV LHC Reach: Excited First-Generation Doublet



gg Resonance: Color Singlet or Octet Scalars

# Colored Scalars

- Models with an extended color sector
- Dynamical EWSB with colored constituents
- Extra-Dimensional Models

### Benchmark Model: Color Octet Scalar

$$\mathcal{L}_{S_8} = g_S d^{ABC} \frac{k_S}{\Lambda_S} S_8^A G_{\mu\nu}^B G^{C,\mu\nu}$$

$$\Gamma(S_8) = \frac{5}{3} \alpha_S \frac{k_S^2}{\Lambda_S^2} m_{S_8}^3$$

Following ATLAS & CMS, take:  $\Lambda_S = m_{S_8}$ 

Han, Lewis, Liu arXiv:1010.4309

### 14 TeV LHC Reach: Color Octet Scalar



# The Jet Energy Profile

# Gluons radiate more than Quarks



Quarks:  $C_F=4/3$  Gluons:  $C_A=3$ 

Question: How does this tendency manifest after showering in a real detector?

See, for example, Ellis, Kunszt, Soper PRL 69 (1992) 3615

# Integrated Jet Shape

Average fraction of jet p<sub>T</sub> lying within a sub-cone of radius r:



Expect quarks form "tighter" jets than gluons, for fixed  $p_T$ 

$$Y(r) = \frac{1}{N^{\text{jet}}} \sum_{\text{jets}} \frac{p_T(0, r)}{p_T(0, R)}, \qquad 0 \le r \le R,$$

# CMS Measurements



Combination of kinematic and compositional effects!

> Good agreement between data and MC ("tuning" required)

CMS, JHEP 1206 (2012) 160

# ATLAS Measurements



Note difference for quark vs. gluon jets, and change consistent with increasing quark-jet fraction

ATLAS, PRD 83 (2011) 052003

A dijet resonance changes the quark/gluon composition in the resonance region: Can we see this using measurements of the jet energy profile?

### Analytical Tool for Understanding Quark/Gluon Jets: NLL Resummation



Collins, Soper, Sterman, PRD 71 (2005) 112002 Li, Li, Yuan, PRL 107 (2011) 152001 PRD 87 (2013) 074025

### Model vs. CMS Data



Li, Li, Yuan, PRD 87 (2013) 074025

# Limitations of JEP Model

- NLL resummation model used to estimate sensitivity of jet energy profile measurements to presence of a new dijet resonance.
- Model includes two phenomenological parameters, which will need to fixed; will be done once LHC data is available. (Will use Tevatron "tune" for statistical analysis.)
- Note: not really dependent on model *correctly predicting* JEP — rather, as a <u>model of size of change</u> in resonance region.
- Will show statistical discrimination present will leave systematic errors to the experts, the experimenters!

# Measuring the Jet Energy Profile

Warm-Up: Signal Only  

$$\psi_{jj}(r) = \psi_1(r) + \psi_2(r)$$
  
1, 2 = q, g  
 $\psi_S(r) = f\psi_{\bar{q}q}(r) + (1 - f)\psi_{gg}(r)$   
 $f = 0.0 \ (C_{\mu}^A)$   
 $= 0.5 \ (q^*)$   
 $= 1.0 \ (S_8^A)$ 

 $\Delta f=0.1\Rightarrow 5\sigma$  discrimination

# Measuring "f"

Benchmark Resonance Parameters: 4 TeV C (tan $\theta$ =0.60), q\* (f<sub>S</sub>=0.4), S<sub>8</sub> (k<sub>S</sub>=0.65) (not excluded, observable with 30 fb<sup>-1</sup> @ 14 TeV)

- MC Simulations of <u>signal events</u>
  - Madgraph v.5, Pythia v.6
- Fastjet clustering, anti-k<sub>T</sub>, R=0.5
- Consider events with  $|M_{jj}-M| < \Gamma/2$
- Examine  $\psi_{S}(r)$  in each MC event
- Accumulate to determine statistical error in  $\psi_{S}(r)$
- Find uncertainty is Gaussian (Poisson errors)
- Statistical Errors Only

$$(\delta\psi_S(0.1))^2 \approx \frac{\sigma^2(0.1)}{S}$$
 with  $\sigma(0.1) \approx 0.4$ 

# Signal Plus Background



*Errors* in  $\Psi_{\rm S}(r)$  and  $\Psi_{\rm B}(r)$ scale the same way:  $(\delta\psi_{OBS}(r))^2 \approx \frac{\sigma^2(r)}{S+B}$   $(\delta\psi_B(r))^2 \approx \frac{\sigma^2(r)}{B}$ 

## Uncertainties including Background

$$\psi_{OBS}(r) = \frac{S}{S+B}\psi_S(r) + \frac{B}{S+B}\psi_B(r) ,$$

$$\psi_S(r) = \psi_{OBS}(r) + \frac{B}{S}(\psi_{OBS}(r) - \psi_B(r)) .$$

$$(\delta\psi_S)^2 \approx \frac{\sigma^2}{S} \left[1 + 2\frac{B}{S}\right] + \frac{(\psi_S - \psi_B)^2}{S}$$

Dilution due to background

Statistical Uncertainty due to signal

### JEP Benchmark Measurements

 $L = 100 \text{ fb}^{-1}$ M = 4 TeV $\Psi_{ij}(\mathbf{r})$ 2.0 JEP Measurement **Discrimination Power:** qq 1.5 qg  $\sigma$  (L=100 fb<sup>-</sup>  $L(5\sigma)$ gg 1.0 5.485  $\bar{q}q$  - qgAverage values determined 4.7110qg - ggby JEP NLL model; 0.5 30 10 $\bar{q}q$  - gguncertainties by MC simulations 0.0  $- r_{0.5}$ 0.1 0.2 0.3 0.4

## JEP Reach: Colorons



Grey Region: 5 $\sigma$  resonance discovery reach Contours:  $\Delta f \leq 0.1, 0.05, 0.02, 0.005$ Benchmark Point in Red

# JEP Reach q\*



Grey Region: 5 $\sigma$  resonance discovery reach Contours:  $\Delta f \leq 0.1, 0.05, 0.02, 0.005$ Benchmark Point in Red

# JEP Reach S<sub>8</sub>



Grey Region: 5 $\sigma$  resonance discovery reach Contours:  $\Delta f \leq 0.1, 0.05, 0.02, 0.005$ Benchmark Point in Red

## Comments

- Statistical discriminating power of JEPs strong can potentially cover entire discovery region
- Systematic issues will dominate area of active and growing interest at ATLAS and CMS.
- Potential improvements: trimming, pruning, grooming - although JEP less sensitive to underlying event than jet mass, etc.

# Conclusions

- The run 2 (and beyond) LHC reach for new dijet resonances extends to very high energies - beyond 4-5 TeV.
- If a new resonance is discovered, and decays only to dijets, the JEP provides a potentially powerful tool to determine the nature of the resonance.
- The statistical power is there, the ultimate utility will depend on a detailed understanding of detector-dependent systematic issues.