

# The scalar spectrum of many-flavour QCD

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for the LatKMI collaboration

SCGT12, 2012 December 4, Nagoya

# LatKMI collaboration

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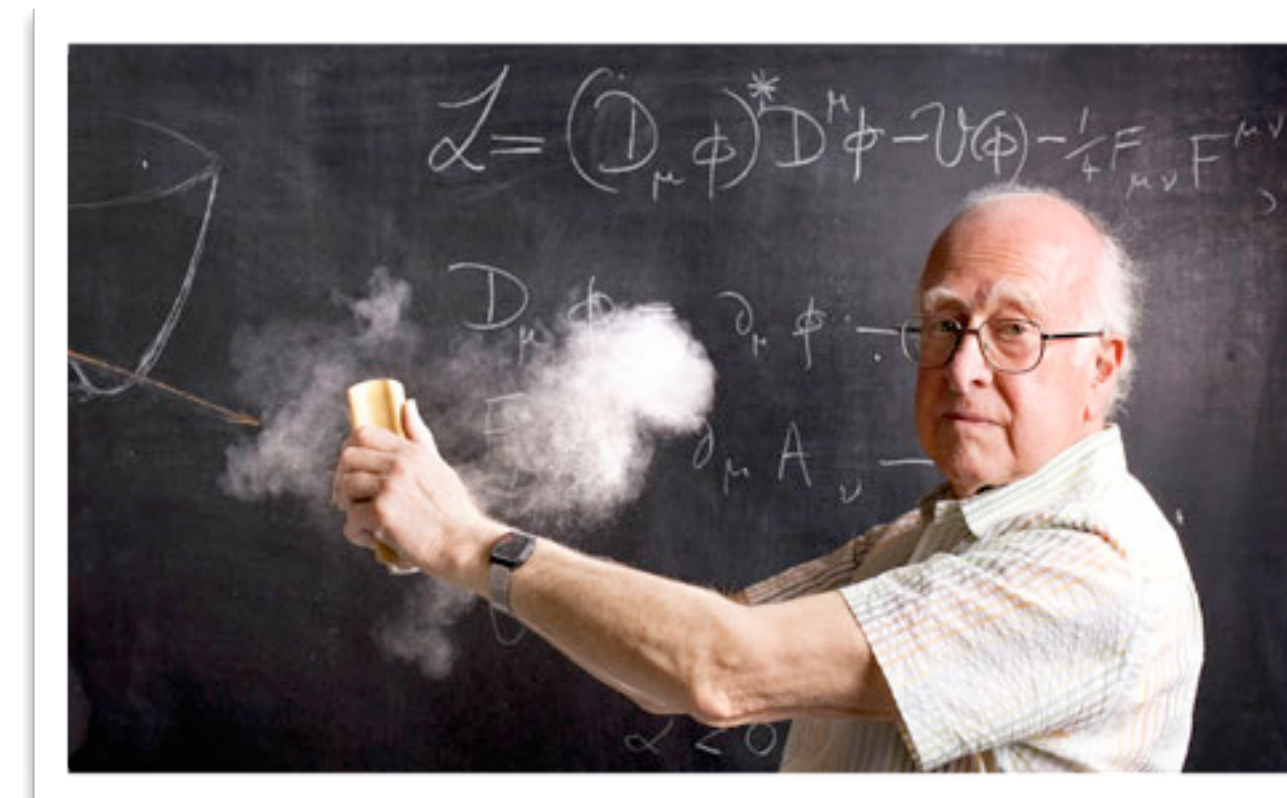
A. Shibata



# Outline

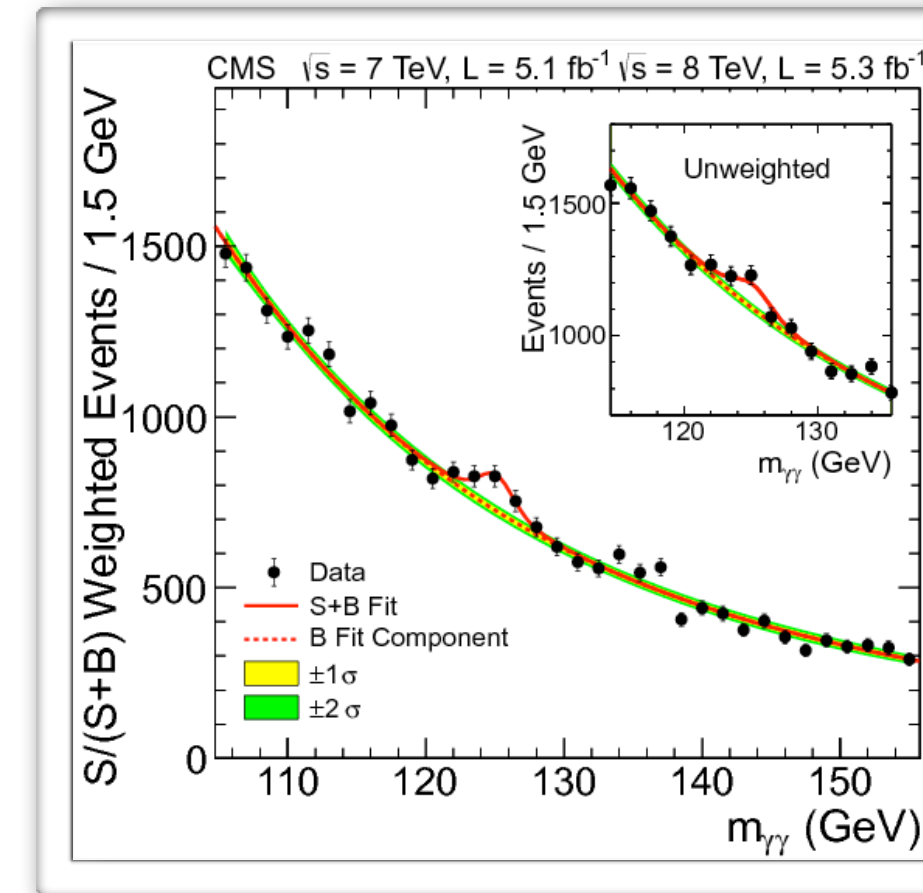
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- Motivations
- Measuring the scalar glueball mass
- Measuring the scalar singlet meson mass
- Results
- Conclusions and outlook

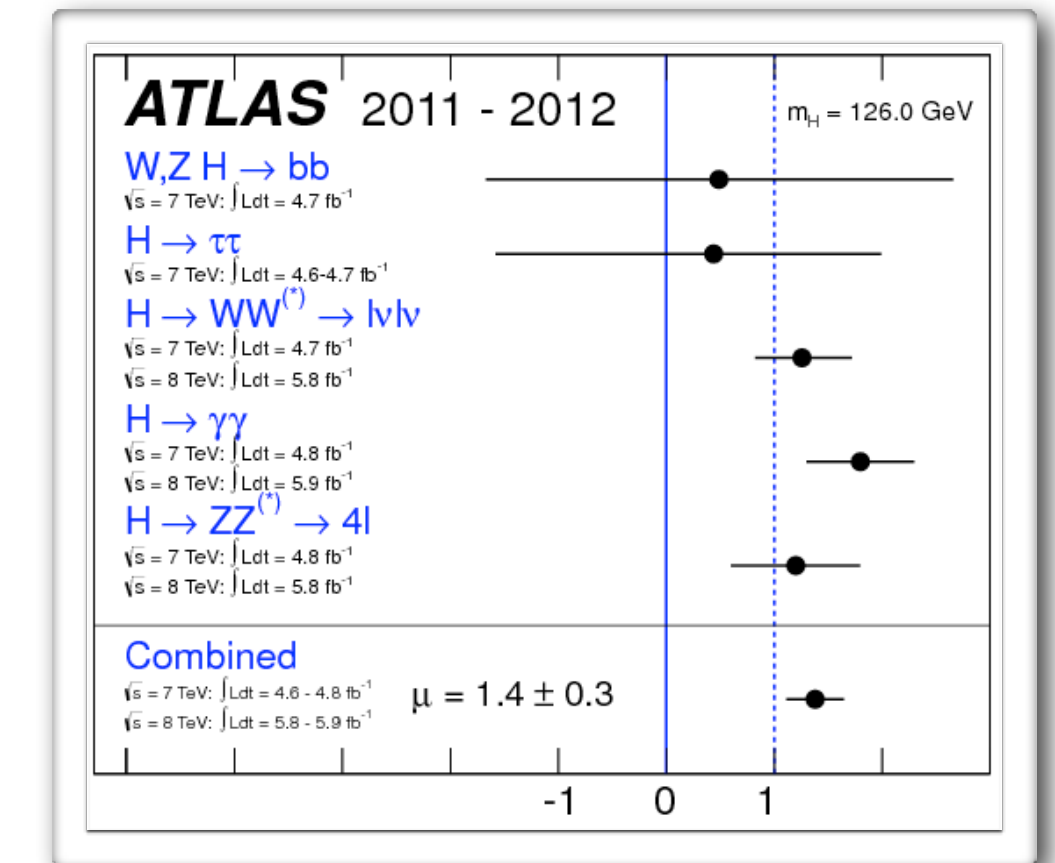


# Motivations

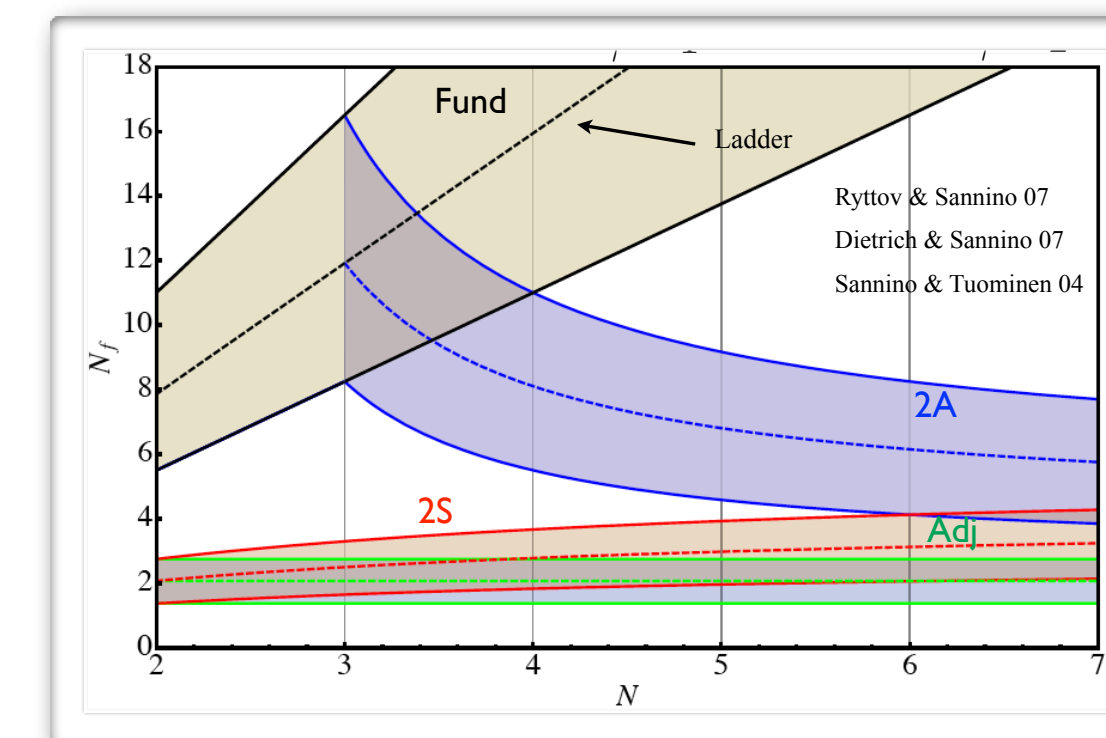
- a Higgs-like particle has been found at LHC
  - $m_H \sim 126$  GeV
  - $J^{PC} = 0^{++}$  (not parity odd at  $2.5\sigma$ )
  - couplings consistent with the Standard Model
- a composite “Higgs” scalar is not ruled out
- this is realised in Walking Technicolor models
- interest in finding the lower edge of the Conformal Window
- a light scalar appears as pseudo-NG boson



Phys.Lett. B716 (2012) 30-61



Phys.Lett. B716 (2012) 1-29



F. Sannino

# QCD with 12 fundamental fermions

- well studied by several groups recently
- however the chiral and continuum physics are still controversial
- additional information can be gained by measuring the scalar spectrum
- spectroscopic techniques already used in QCD can be developed for many flavours
- focus on scalar glueballs and scalar mesons
  - masses, mixing, scaling, ...

Conformal or Walking? Monte Carlo renormalization group studies of SU(3) gauge models with fundamental fermions

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Twelve massless flavors and three colors below the conformal window

Zoltán Fodor<sup>a,e</sup>, Kieran Holland<sup>b</sup>, Julius Kuti<sup>\*c</sup>, Dániel Nógrádi<sup>e</sup>, Chris Schroeder<sup>d</sup>

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Lattice study of infrared behaviour in SU(3) gauge theory with twelve massless flavours

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Lattice study of conformality in twelve-flavor QCD

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Toshihide Maskawa,<sup>1</sup> Kei-ichi Nagai,<sup>1</sup> Hiroshi Ohki,<sup>1</sup>  
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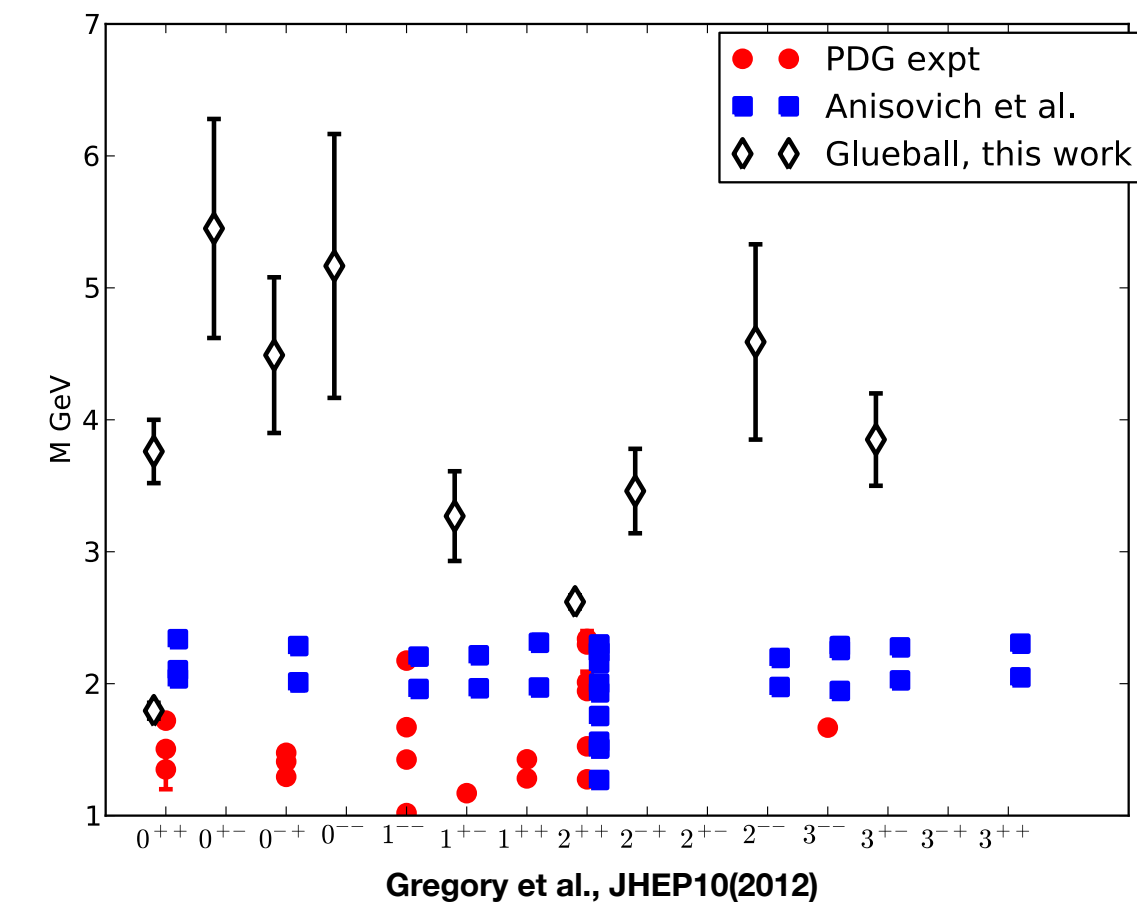
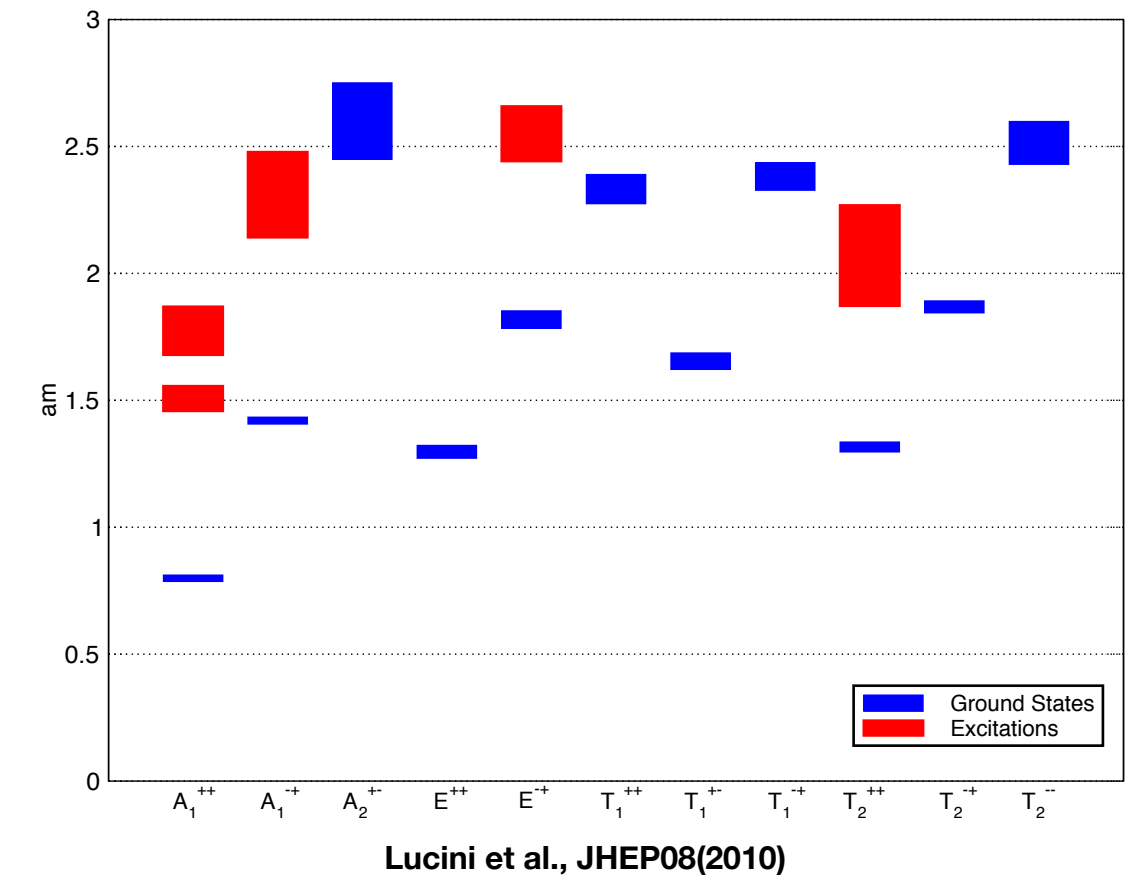
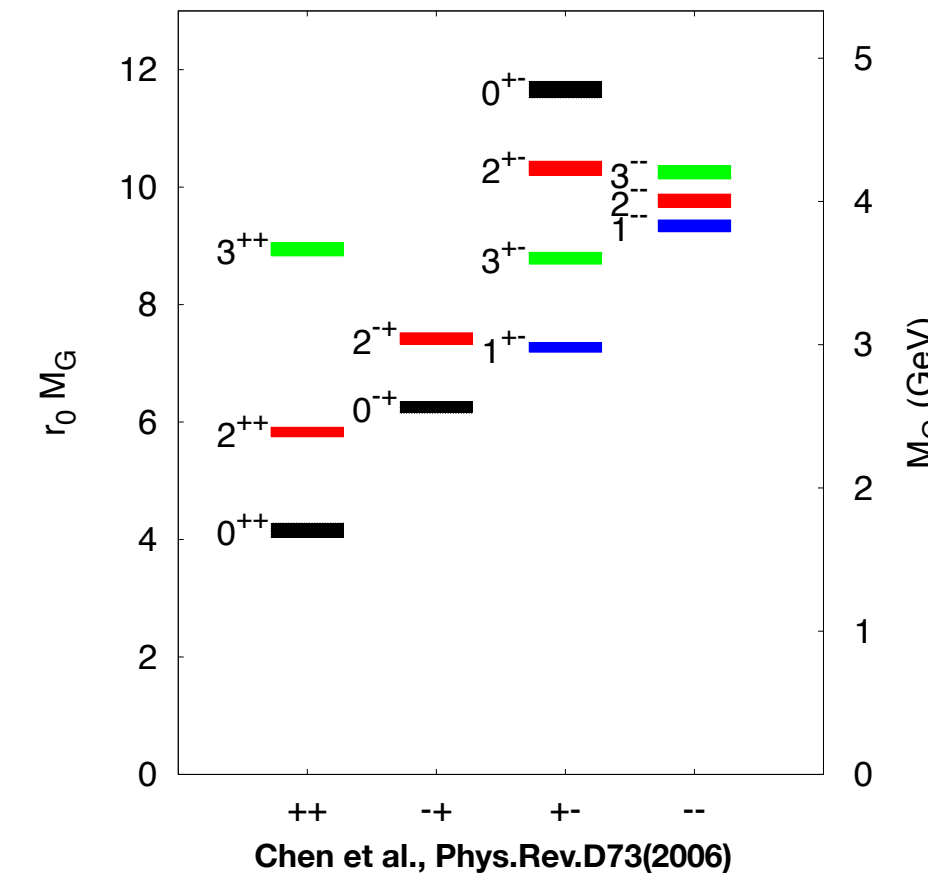
# Lattice simulations

- SU(3) gauge field + 12 dynamical fermions
  - gauge action → tree level Symanzik
  - fermion action → HISQ (Highly Improved Staggered Quarks)
  - modified HMC based on MILC code (v7)
- fixed gauge coupling →  $\beta = 4.0$
- two volumes →  $V=18^3, 24^3$
- five bare fermion masses →  $am_f = 0.06 - 0.16$
- long trajectories → >15000 MD units

L	T	$am_f$	traj.	confs.	bins
18	24	0.06	15445	3090	103
18	24	0.08	22200	4440	148
18	24	0.10	15300	3060	102
18	24	0.12	21500	4300	172
18	24	0.16	18450	3690	123
24	32	0.06	28000	14000	700
24	32	0.08	14320	7160	358
24	32	0.10	8960	4480	224

# Glueball spectroscopy

- the glueball spectrum is well studied in SU(N) Yang-Mills theories
- the lightest glueballs are starting to be investigated in full LQCD
- yet to be done in many-flavour QCD
  - typically heavy (mass determined by  $\Lambda_{\text{QCD}}$ )
  - typically noisy (require large statistic)



# Glueball spectroscopy: operators

- eigenstates of the Hamiltonian are classified according to the irreducible representations of the cubic group

$$\{A_1(1), A_2(1), E(2), T_1(3), T_2(3)\}$$

- suitable gauge-invariant operators must be constructed that respect the symmetries

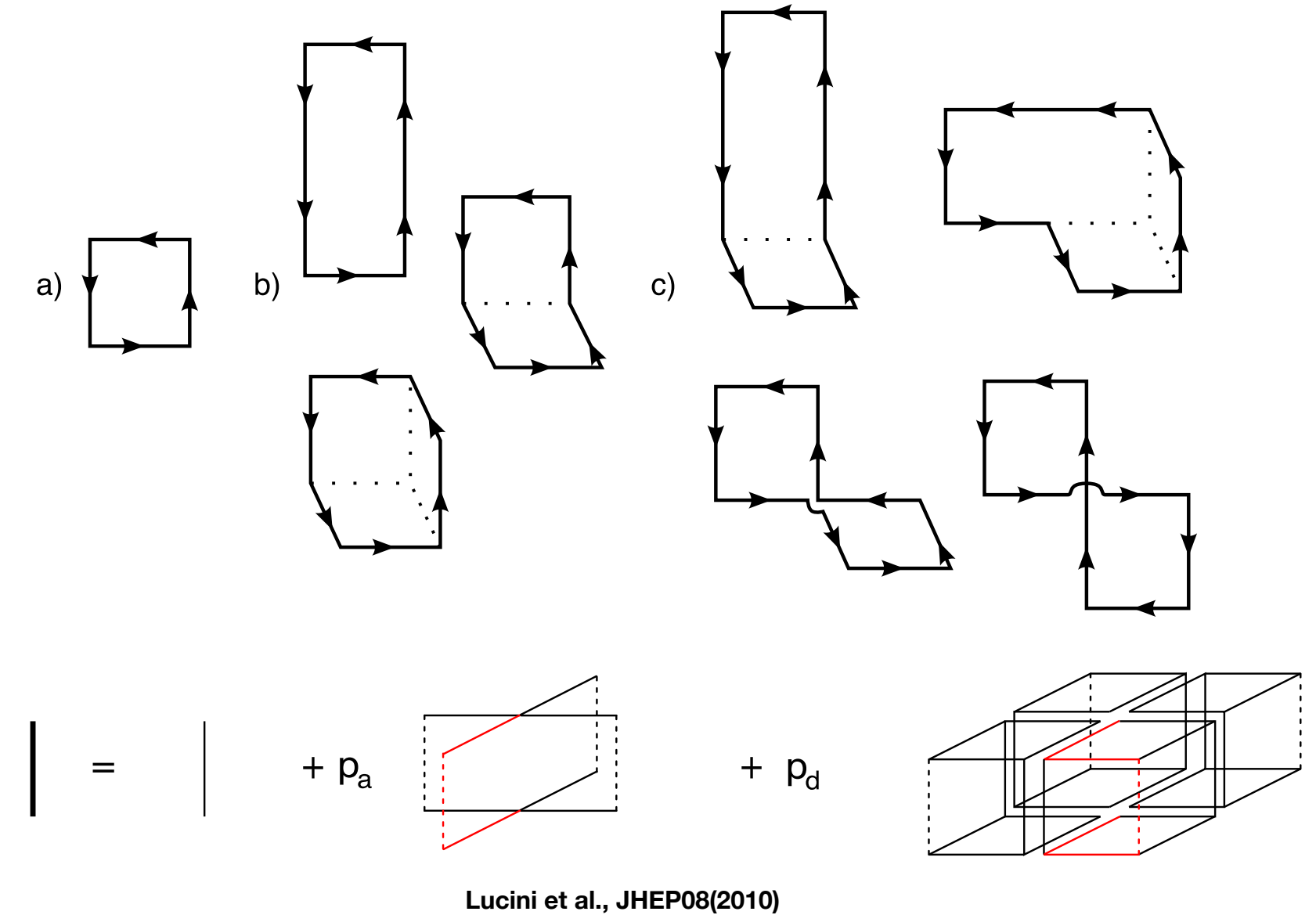
$$\mathcal{O}_G(t) = \frac{1}{L^3} \sum_{x \in L^3} \text{Tr} \left( \prod_{l \in \mathcal{W}(x)} U_l \right) \quad \mathcal{O}_G^{(R)}(t) = \sum_{\alpha=1}^{24} a_{\alpha}^{(R)} \mathcal{R}_{\alpha} [\mathcal{O}_G(t)]$$

- vacuum contributions must be subtracted in the scalar case

$$\mathcal{O}^{(A_1)}(t) - \langle 0 | \mathcal{O}^{(A_1)} | 0 \rangle$$

- improved operators are obtained by blocking and smearing algorithms

$J$	$A_1$	$A_2$	$E$	$T_1$	$T_2$
0	1	0	0	0	0
1	0	0	0	1	0
2	0	0	1	0	1
3	0	1	0	1	1
4	1	0	1	1	1





# Glueball spectroscopy: variational analysis

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- basis of operators →

$$\{\mathcal{O}_1(t), \dots, \mathcal{O}_n(t)\}$$

- matrix of correlators →

$$C_{ij}(t) = \sum_{\tau} \langle 0 | \mathcal{O}_i^{\dagger}(\tau + t) \mathcal{O}_j(\tau) | 0 \rangle$$

- generalized eigenvalue problem →

$$C_{ij}(t) v_j^{\alpha} = \lambda^{\alpha} v_i^{\alpha} \quad \Phi_{\alpha}(t) = \sum_{i=1}^n v_i^{\alpha} \mathcal{O}_i(t)$$

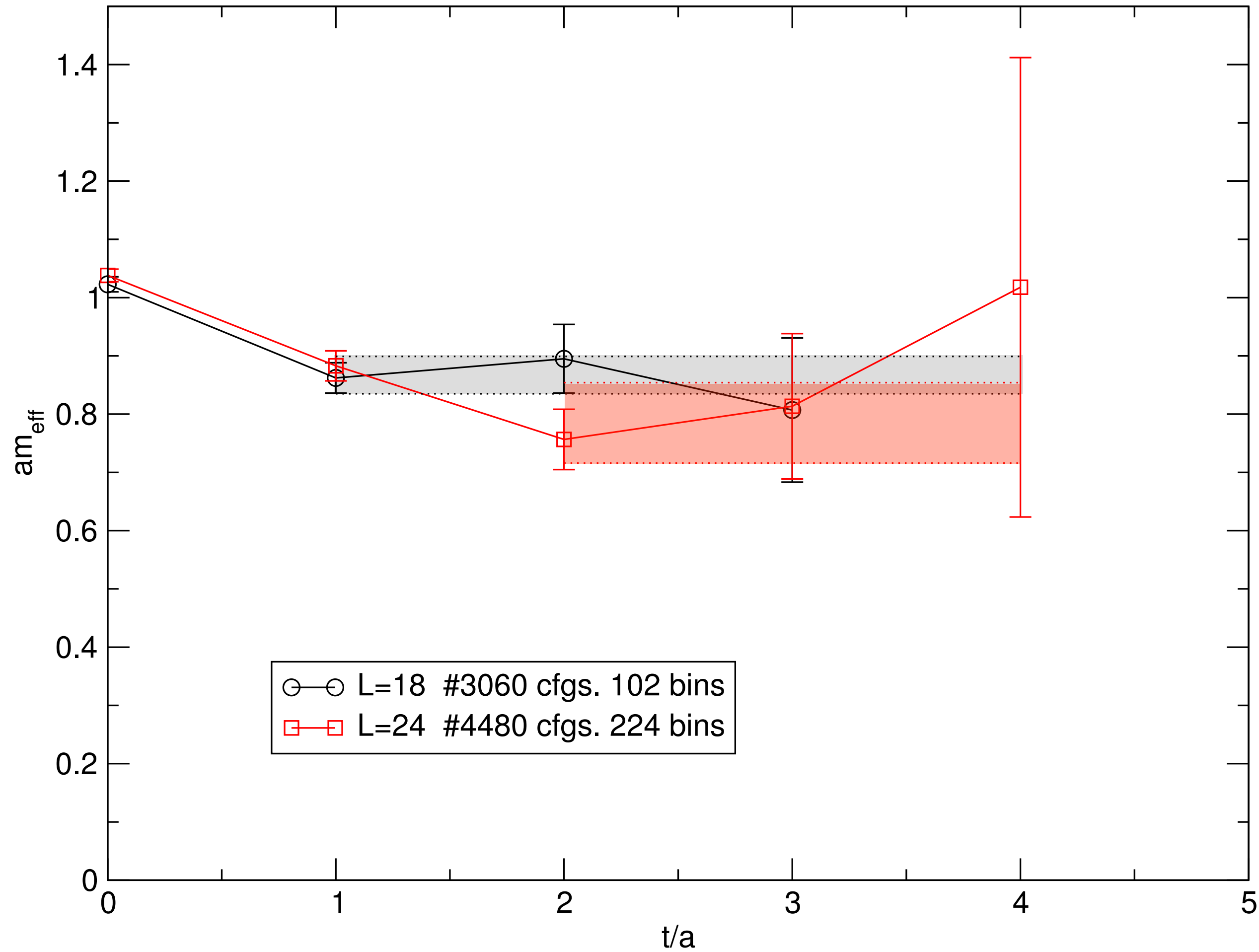
- ground state correlator fit →

$$\langle \Phi_{\alpha}^{\dagger}(t) \Phi_{\alpha}(0) \rangle = |c_{\alpha}|^2 \left( e^{-m_{\alpha} t} + e^{-m_{\alpha}(T-t)} \right)$$

- the effective mass plateau is used to determine the fitting window on the correlator
- excited states contributions are given by  $\alpha > 0$

# Results: scalar glueball

Scalar glueball effective mass  $\beta=4.0$   $am_f=0.10$



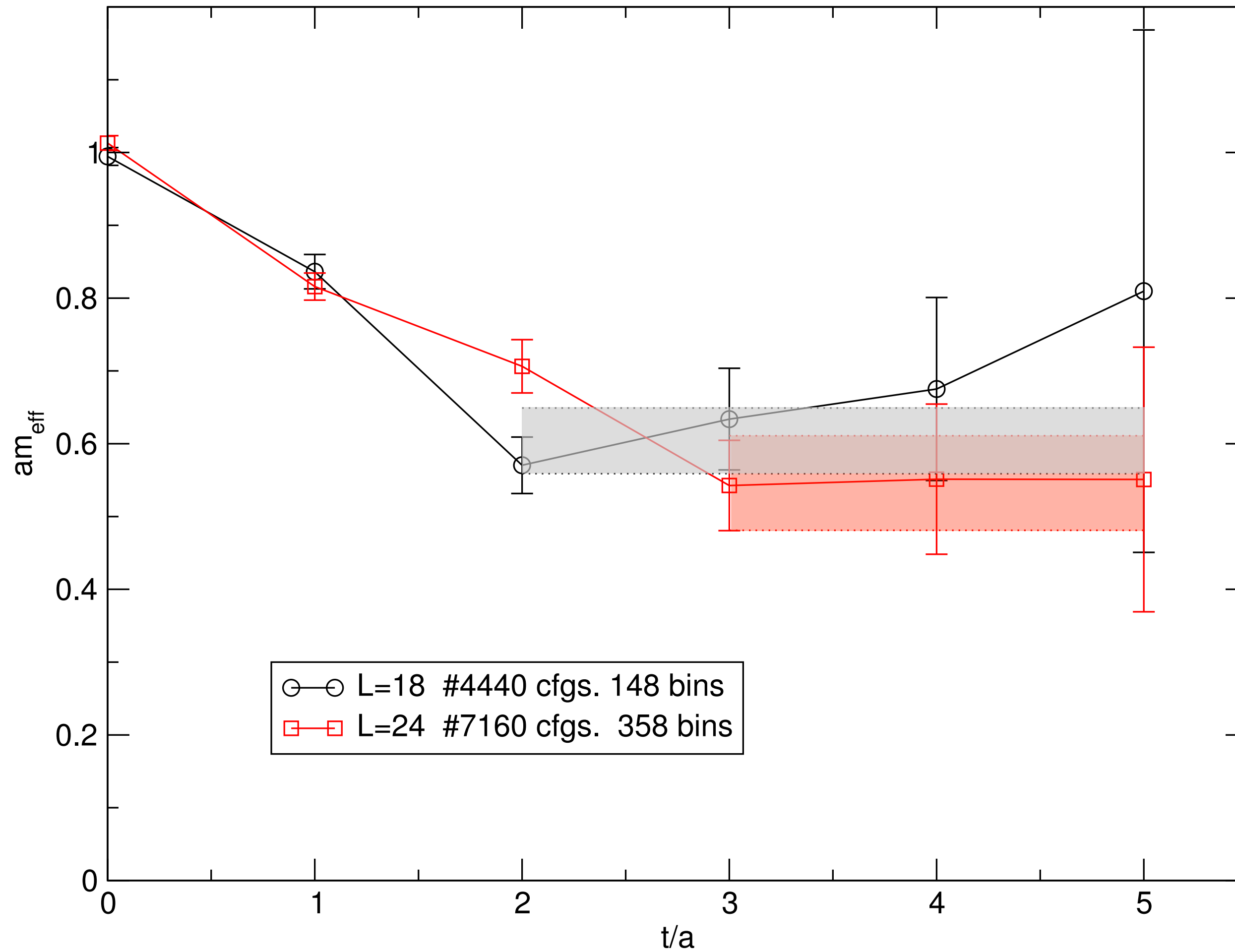
$$m_{0^{++}} > m_{\pi}$$

$L=18 \rightarrow am=0.867(32)$

$L=24 \rightarrow am=0.785(69)$

# Results: scalar glueball

Scalar glueball effective mass  $\beta=4.0$   $am_f=0.08$

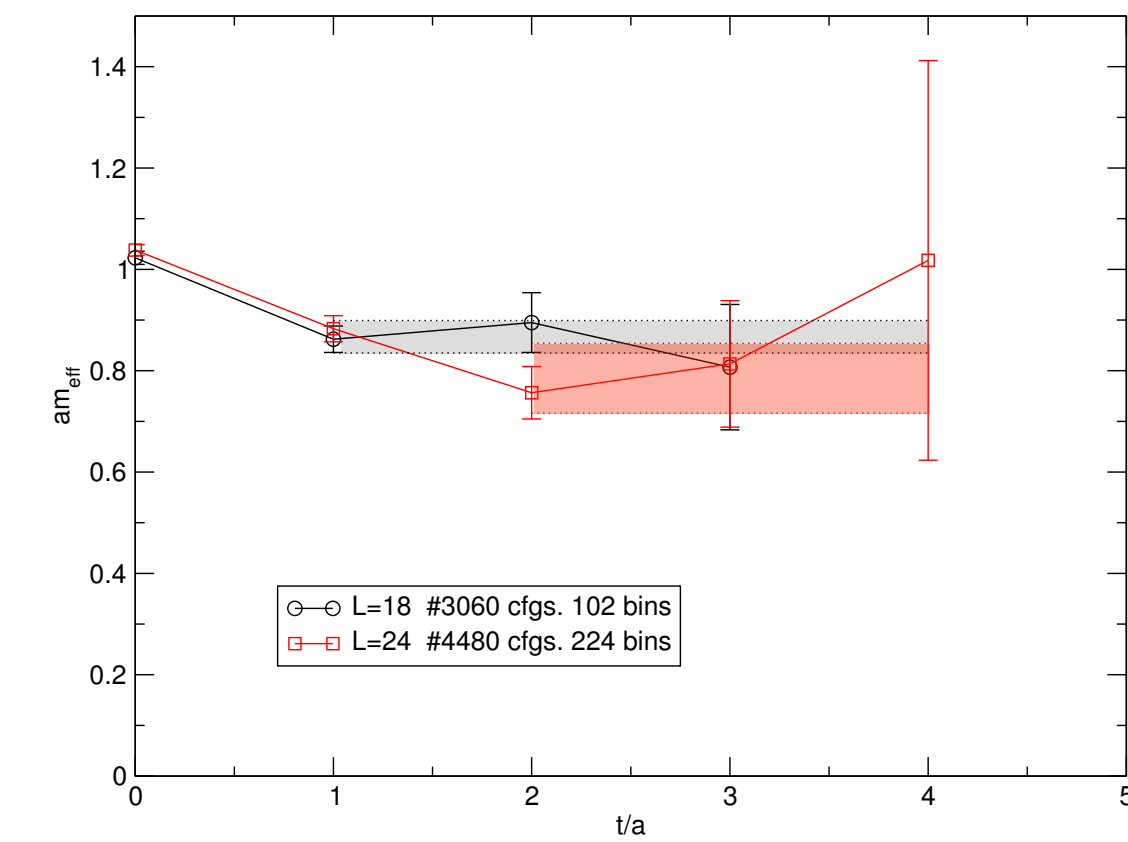


$$m_{0++} \gtrsim m_{\pi}$$

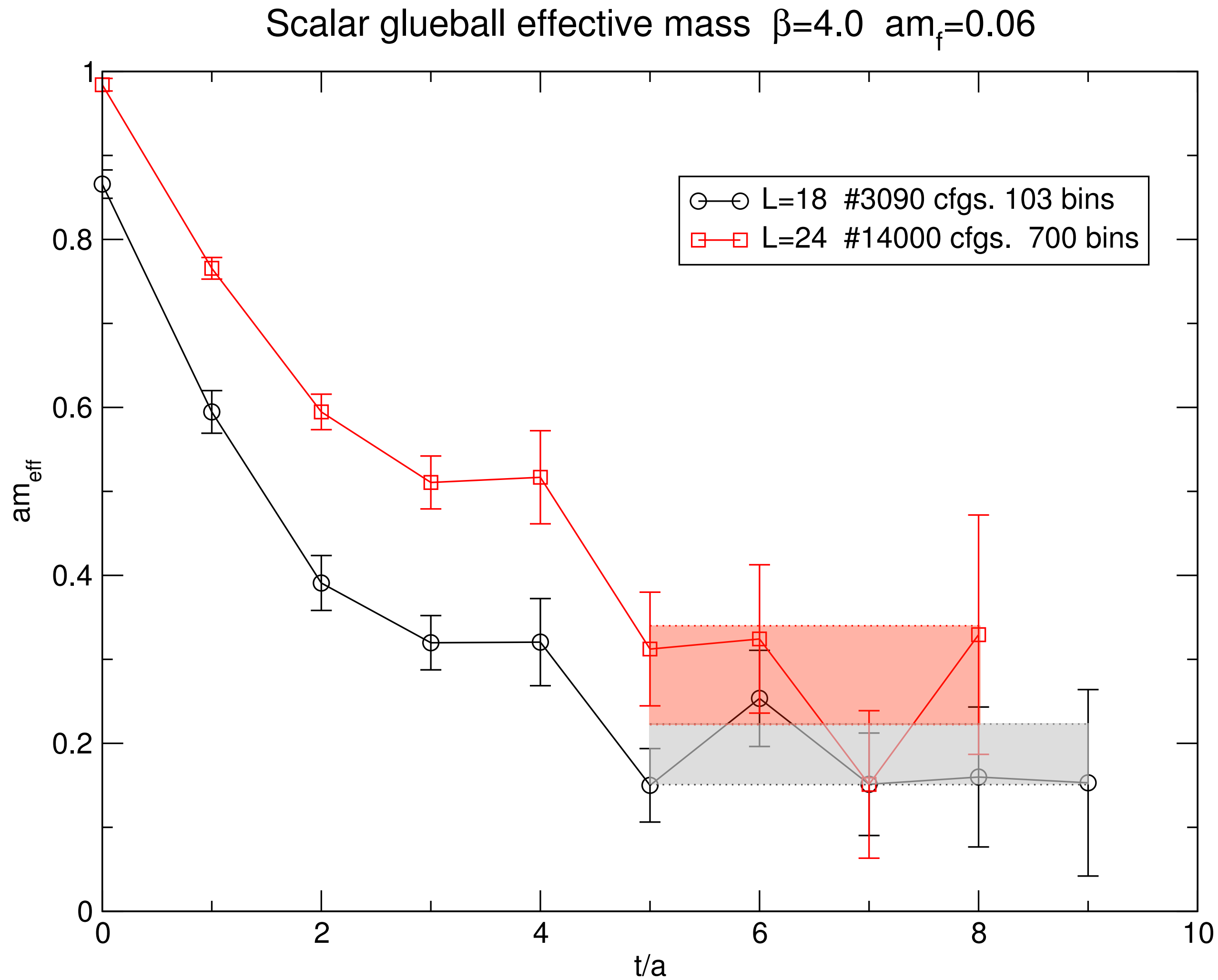
L=18  $\rightarrow am=0.604(45)$

L=24  $\rightarrow am=0.546(65)$

Scalar glueball effective mass  $\beta=4.0$   $am_f=0.10$



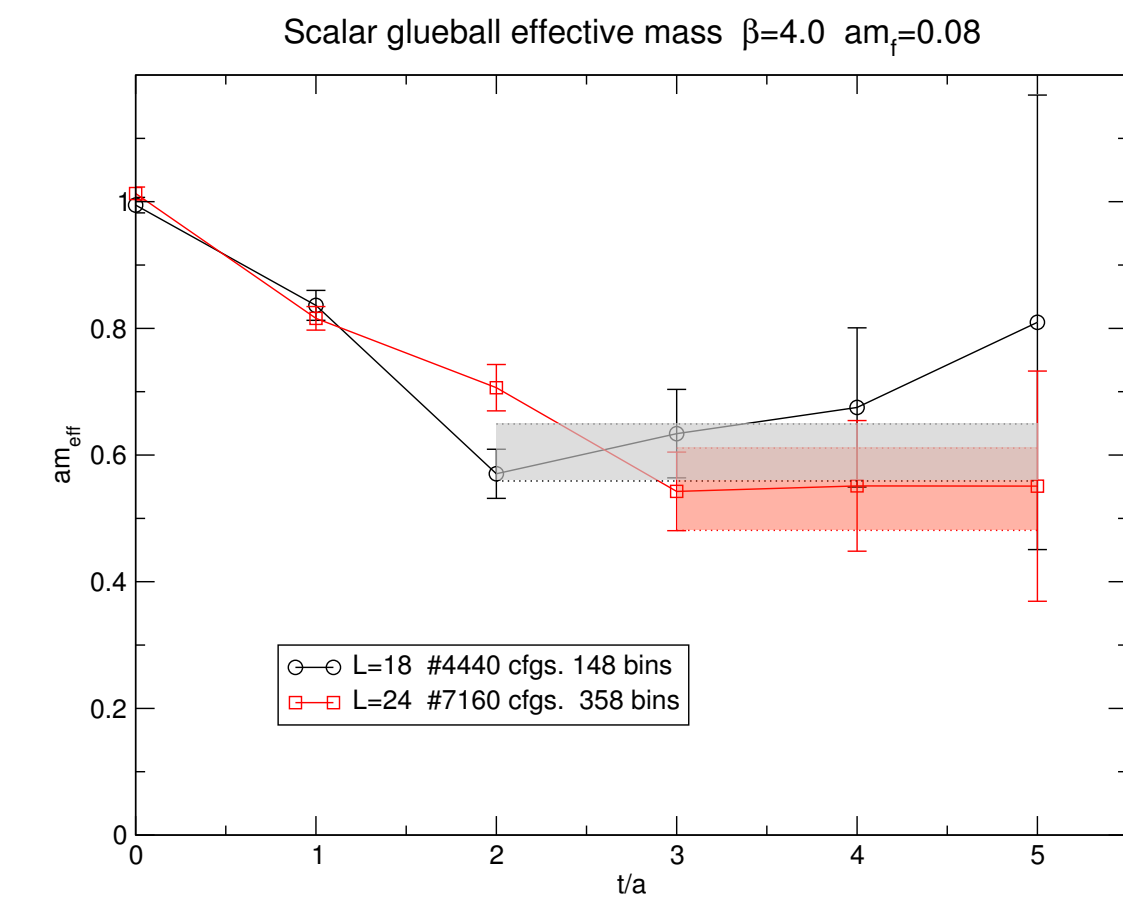
# Results: scalar glueball



$$m_{0++} \lesssim m_{\pi}$$

$L=18 \rightarrow am=0.187(36)$

$L=24 \rightarrow am=0.281(59)$

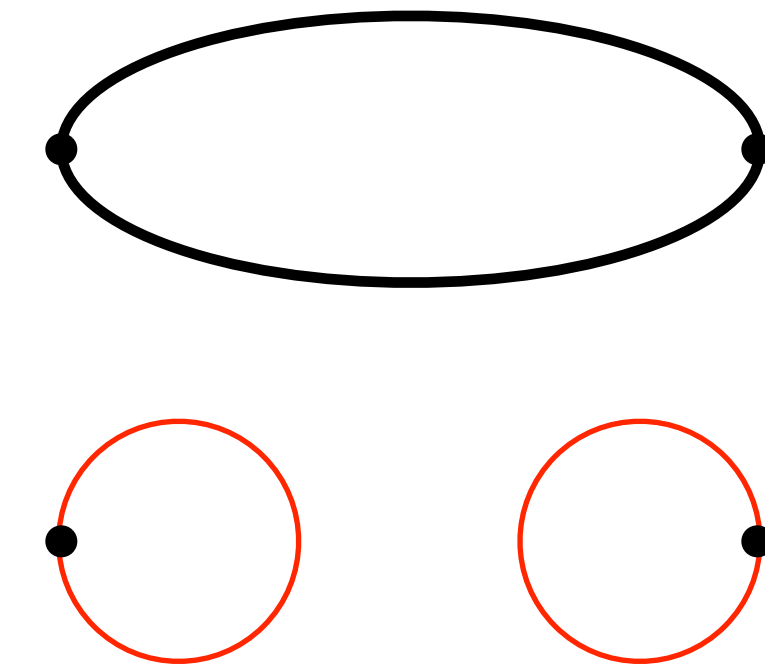
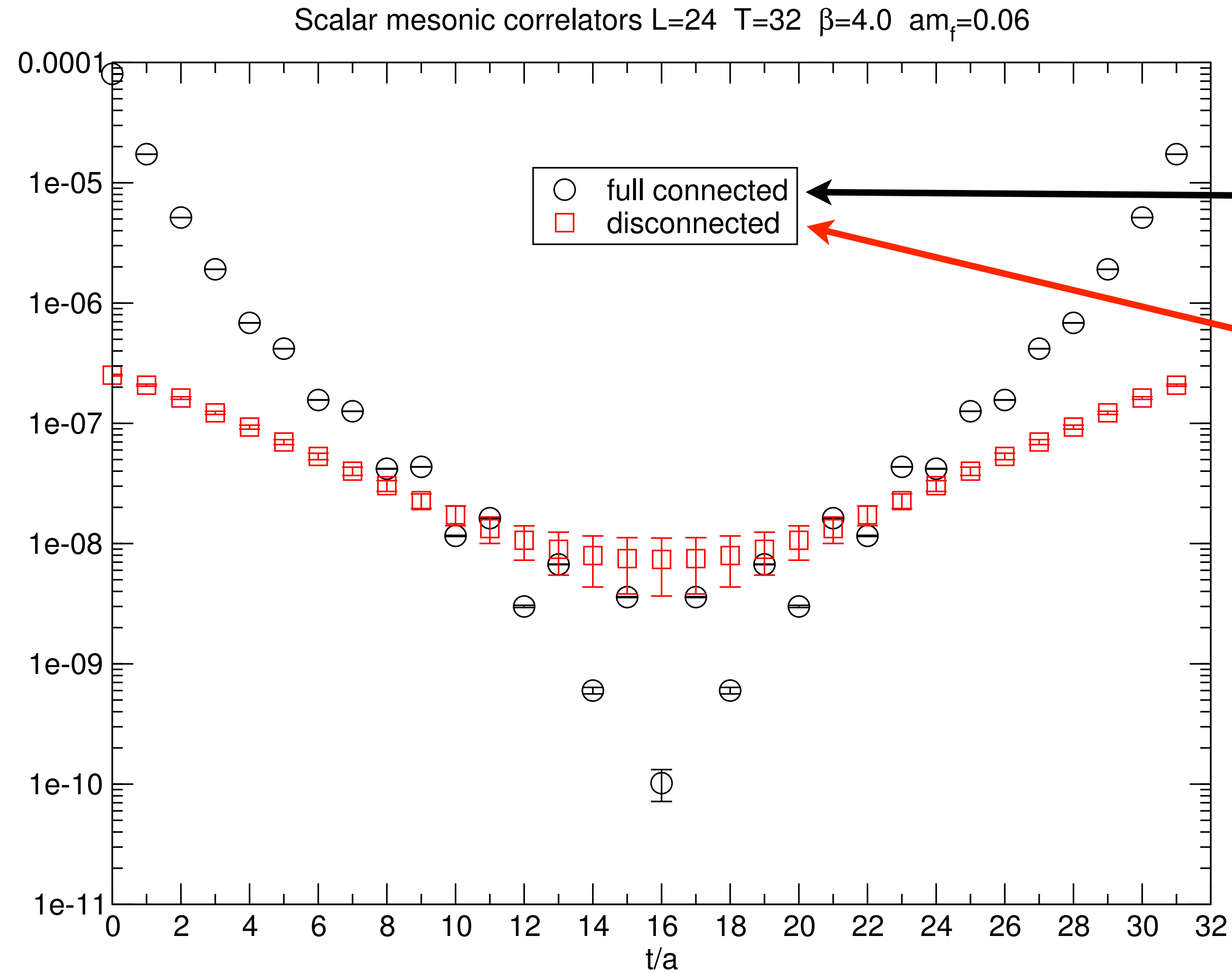


# Meson spectroscopy: scalar flavour-singlet

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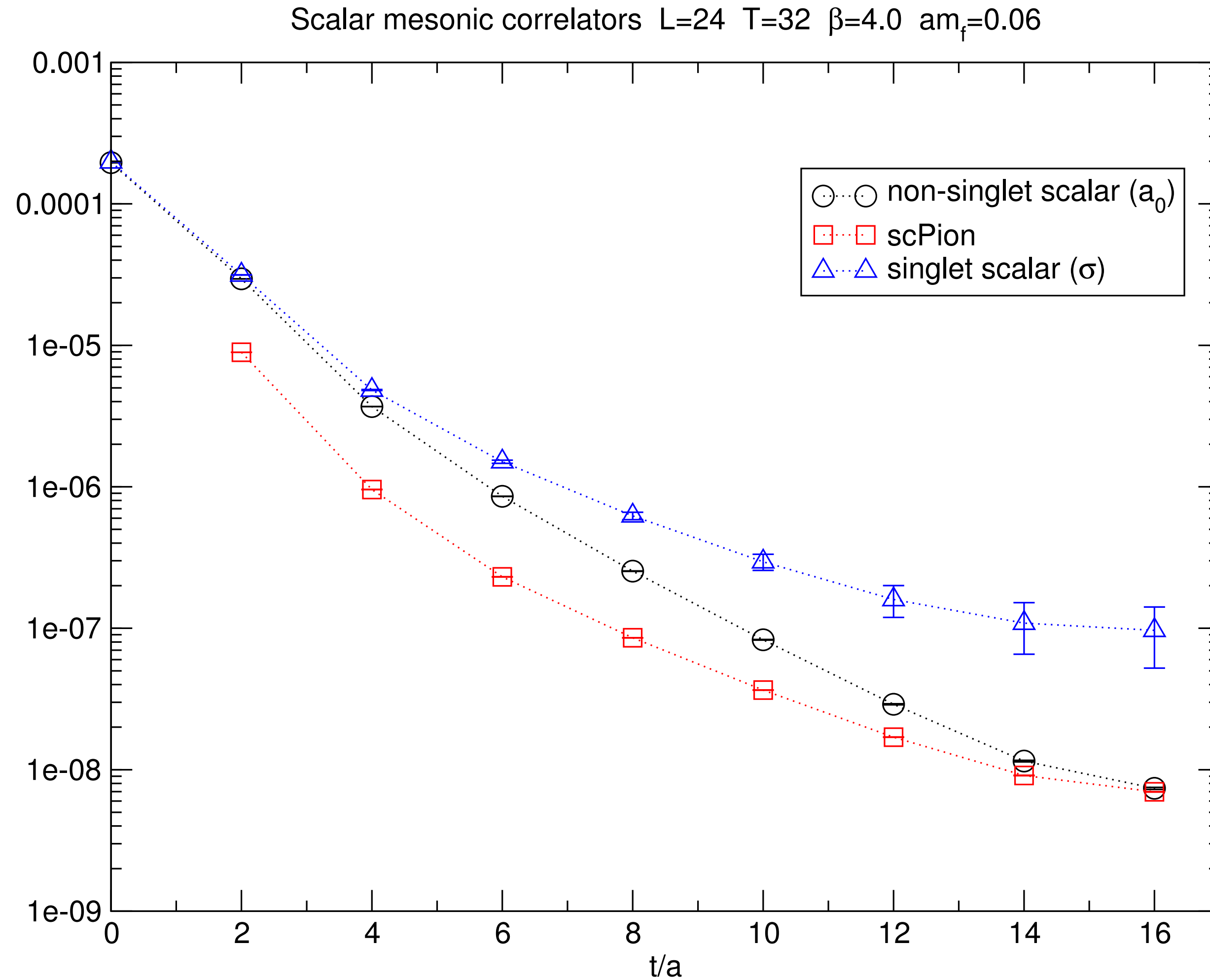
- $N_f=12$  staggered formalism  $\rightarrow$  the scalar interpolating mesonic operator couples to two states
  - $0^-$  state  $\rightarrow$  is the scPion (non-goldstone), degenerate with the  $\pi$  (goldstone) state in our simulations
  - $0^+$  state  $\rightarrow$  contains the scalar non-singlet  $a_0$  and the scalar flavour-singlet  $f_0$  (or  $\sigma$  in QCD)
- disconnected quark loops must be evaluated in the latter case (together with the vacuum contribution)
  - for this we use stochastic gaussian sources and follow Gregory et al., PRD77:065019(2008)
- thanks to the large statistic we are able to get a signal for the scalar flavour-singlet state (*preliminary*)

# Results: scalar flavour-singlet meson



- connected and disconnected correlator measured on 14000 configurations
- 64 stochastic gaussian sources used for the disconnected piece on each configuration
- 2 stochastic gaussian sources used for the connected piece on each configuration

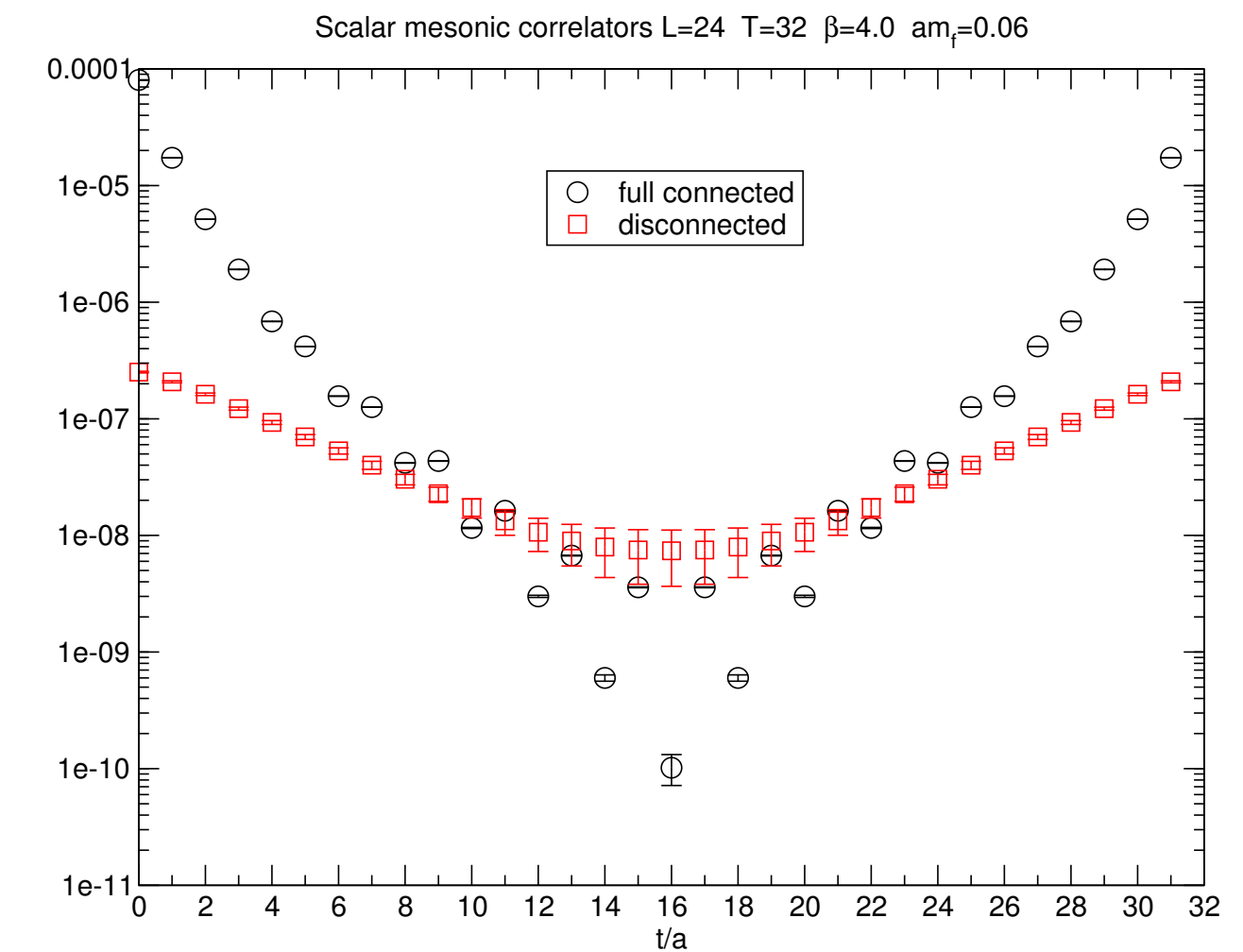
# Results: scalar flavour-singlet meson



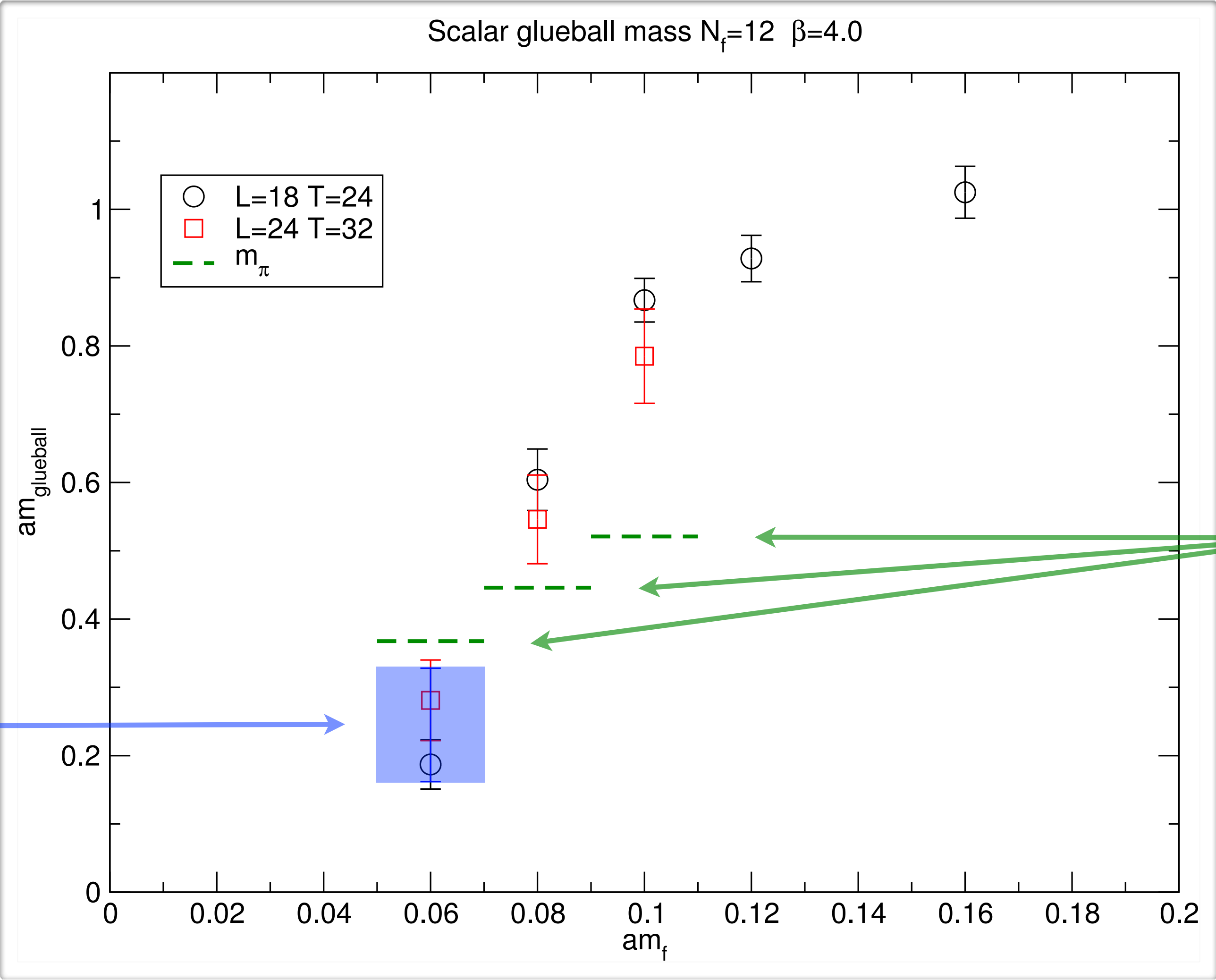
$$C_{0+}(t) = 2C(t) + C(t+1) + C(t-1)$$

$$C_{0-}(t) = 2C(t) - C(t+1) - C(t-1)$$

$$C_\sigma(t) = -C_{0+}(t) + 3D_{0+}^2(t)$$



# Results: summary



scalar flavour-singlet  
meson mass

Goldstone pion mass



# Conclusions

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- first study of scalar spectrum in many-flavour QCD, focused on  $N_f=12$ 
  - scalar glueball and scalar meson interpolating operators are used
  - mesonic disconnected diagrams are evaluated for the scalar flavour-singlet case
- for light bare quark masses there is a light state in the scalar channel with  $m_{SC} < m_\pi$
- both glueball operators and mesonic operators couple to this state
  - important to understand the mixing of such contributions
- starting point of a more careful analysis of light scalar states in (near-)conformal theories
  - important non-perturbative study in phenomenological models of Walking Technicolor

# Outlook

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- identify and reduce systematic errors
  - larger volumes at light quark masses are necessary to carefully estimate finite size errors
  - possible discretisation effects need to be studied by simulating at different lattice spacing
  - longer trajectories will help improving the identification of effective mass plateaux
- lighter quark masses where the scalar state is light are needed to identify its scaling properties
- a variational analysis including both gluonic and mesonic operators in the scalar channel is in progress
  - can give informations about the mixing between the two kind of operators
- different number of flavours → study the scalar state in a candidate model for the “walking” scenario

*Thank you*