Looking for New Physics via the Flavor observables

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Physics beyond the Standard model (SM)

The SM of particle physics is the successful theory which explain many experimental results and phenomena



- Structure of matter
 Quarks and Leptons : fundamental constituents of matter
- Mechanism of force propagation
 Force particles : carry forces
- Mechanism of getting mass
 Higgs boson : gives mass to matter

On 4 July 2012, the Higgs boson (=last missing peace of the SM) has been found at LHC (Large Hadron Collider) experiment at CERN

Physics beyond the Standard model (SM)

However, the SM is not satisfactory:

- Neutrino mass
- Cosmological problems (Dark matter, Dark energy, Baryon asymmetry,,,)
- Hierarchy problem
- Origin of the flavor structure
 :

We need to search for New physics beyond the SM

How to search for new physics

Direct search High energy frontier ATLAS/CMS (LHC)

Generate new particles directly with high energy collider

SM particles



NP particles



Indirect searchHigh luminosity frontiere.g.) Belle/Belle II

Search for NP through the virtual effects of new particles







How to search for new physics

Direct search High energy frontier ATLAS/CMS (LHC)

Generate new particles directly with high energy collider

New particles have not been observed yet (~1TeV) The NP scale might be higher than the TeV scale

Indirect search High luminosity frontier e.g.) Belle/Belle I



Search for NP through the virtual effects of new particles

Indirect searches for NP can explore above TeV scale

 \Rightarrow Flavor physics

Flavor physics

Flavor physics: flavor (type of quarks (leptons)) transition processes

Typical example)

Decay processes and CP asymmetries in Kaon and B mesons

Flavor physics

In the SM, flavor transition processes are governed by Cabibbo-Kobayashi-Maskawa (CKM) matrix The unitarity triangle

 $V_{cd}V_{cb}^*$

 Φ 1 : CP asymmetry in

 $B_d^0 \rightarrow J/\psi K_s$ decay

Measured by B factory



Kobayashi and Maskawa won the 2008 Nobel prize!





Flavor Changing Neutral Current

FCNC is strongly suppressed in the SM \Rightarrow high sensitivity to NP

- Loop factors
- GIM mechanism (CKM unitarity)
- CKM hierarchy :



$$\begin{array}{c} V_{ts}^*V_{td}\sim 5\cdot 10^{-4} \ll V_{tb}^*V_{td}\sim 10^{-2} < V_{tb}^*V_{ts}\sim 4\cdot 10^{-2} \\ \textbf{K system} \qquad \qquad \textbf{Bd system} \end{array}$$

Sensitive to high scale NP beyond the energy range accessible at the LHC

$$\Lambda_{\mathrm{NP}} \sim egin{pmatrix} \mathcal{O}(10^5 \mathrm{TeV}) & : K^0 \ \mathcal{O}(10^3 \mathrm{TeV}) & : B_{d,s} \ \end{bmatrix}$$
 [G.Isidori, 1507.00867]

SM works very well



there is still room for NP e.g.) O(10-20%) NP contributions to B-Bbar mixing are still allowed

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Some discrepancies from the SM are reported

- angular distributions in $BR(B \to K^* \mu \mu)$: P'_5 ~3 σ
- lepton flavor non-universality : $R_K \equiv \frac{BR(B^+ \to K^+ \ell \ell)}{BR(B^+ \to K^+ ee)}$ ~2.60
- leptonic decay rates : $BR(B \rightarrow D^* \tau \nu)$ ~3.9 σ
- Direct CP violation in Kaon : ϵ'/ϵ ~2.9



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Direct CP violation in Kaon : ϵ'/ϵ ~2.90

-> discuss later in detail

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It may be a hint of new physics beyond the SM

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- What is ε'/ε anomaly ?
 - CP violation in Kaon
 - Current status of ε'/ε
- New physics implications
 - Correlations with other observables: $K \rightarrow \pi v v$
 - Supersymmetric scenario





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CP violation in Kaon



 $\begin{array}{l} \pi: \mathsf{pseudo \ scalar} \ (0^{-+}) \\ \Rightarrow \mathsf{CP} \ \mathsf{odd} \end{array}$

$$\begin{split} |K_L\rangle \neq \frac{1}{\sqrt{2}} \begin{bmatrix} |K^0\rangle + |\bar{K}^0\rangle \end{bmatrix} & \text{CP}|K^0\rangle = -|\bar{K}^0\rangle \\ = \frac{1}{\sqrt{2}} \begin{bmatrix} (1+\epsilon)|K^0\rangle + (1+\epsilon)|\bar{K}^0\rangle \end{bmatrix} \\ K_1 = \frac{1}{\sqrt{2}}(K^0 - \bar{K}^0) & :\text{CP even} \\ K_2 = \frac{1}{\sqrt{2}}(K^0 + \bar{K}^0) & :\text{CP odd} \\ = |K_2\rangle + \epsilon \ |K_1\rangle \\ \text{CP odd} & \text{even} \end{split}$$

CP violation in Kaon



Indirect : ε (mixing)

$$\epsilon_K \simeq \frac{\mathrm{Im} M_{12}^K}{\Delta M^K}$$
$$(\epsilon_K)_{\mathrm{exp}} = (2.228 \pm 0.011) \times 10^{-3}$$



K-Kbar mixing

Direct : ε' (decay)

$$\left|\frac{\eta_{00}}{\eta_{+-}}\right|^{2} \simeq 1 - 6 \operatorname{Re}\left(\frac{\varepsilon'}{\varepsilon}\right) \qquad \eta_{00} = \frac{A(K_{\mathrm{L}} \to \pi^{0}\pi^{0})}{A(K_{\mathrm{S}} \to \pi^{0}\pi^{0})} \\ \eta_{+-} = \frac{A(K_{\mathrm{L}} \to \pi^{+}\pi^{-})}{A(K_{\mathrm{S}} \to \pi^{+}\pi^{-})} \\ (\epsilon'/\epsilon)_{\mathrm{exp}} = (16.6 \pm 2.3) \times 10^{-4}$$

K->ππ decay

Direct CP violation : ε'/ε $A_{0,2} = A(K^0 \to (\pi\pi)_{I=0,2})$ ϵ_K' $\frac{\omega}{\sqrt{2} |\epsilon_K|_{\exp} \operatorname{Re} A_0}$ ImA_0 $\operatorname{Im} A_2$ ϵ_K ω Dominated by QCD penguin EW penguin u, du, d \bar{u}, \bar{d} g \bar{u}, d \overline{d}

Direct CP violation : ε'/ε $A_{0,2} = A(K^0 \to (\pi\pi)_{I=0,2})$ ϵ'_K $-\frac{\omega}{\sqrt{2} |\epsilon_K|_{\exp} \operatorname{Re} A_0}$ $\frac{1}{\omega}$ $\mathrm{Im}A_2$ $\mathrm{Im}A_{\mathrm{f}}$ ϵ_K Dominated by QCD penguin EW penguin u, d \overline{u}, d $\Delta I = 1/2$ rule $\frac{\text{Re}A_0}{\text{Re}A_2} \equiv \frac{1}{\omega} = 22.46$

In SM, there is accidental cancellation between ImA_0 and ImA_2 due to the enhancement factor $1/\omega$

Direct CP violation : ε'/ε $A_{0,2} = A(K^0 \to (\pi\pi)_{I=0,2})$ $= -\frac{\omega}{\sqrt{2} |\epsilon_K|_{\text{exp}} \text{Re}A_0}$ ϵ'_K _ $\left(\mathrm{Im}A_0 - \frac{1}{\omega}\mathrm{Im}A_2\right)$ ϵ_K **EW** penguin QCD penguin $\frac{g \overleftarrow{b} \overline{u}, \overline{d}}{\overleftarrow{\gamma}, t, \overleftarrow{\zeta}} d$ $O_8 = \frac{3}{2} (\bar{s}_\alpha d_\beta)_{V-A} \sum e_q (\bar{q}_\beta q_\alpha)_{V+A}$ $O_6 = (\bar{s}_\alpha d_\beta)_{V-A} \sum (\bar{q}_\beta q_\alpha)_{V+A}$

ϵ'/ϵ anomaly

Using the first lattice result, ϵ'/ϵ has been calculated in the SM as

SM
$$(\varepsilon'/\varepsilon)_{\mathrm{SM}} = (1.9 \pm 4.5) \times 10^{-4}$$
 [Buras et.al '15]

Exp
$$(\epsilon'/\epsilon)_{
m exp} = (16.6 \pm 2.3) \times 10^{-4}$$
 [NA48, KTeV]

2.9σ difference

New physics providing enhancement of ε'/ϵ ?

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Summary

ϵ'/ϵ beyond SM



NP in ImA₂ is favored because of $\Delta I=1/2$ enhancement factor : $1/\omega \sim 22$

The type of (ImA₂)_{NP} :



ϵ'/ϵ beyond SM



NP in ImA₂ is favored because of $\Delta I=1/2$ enhancement factor : $1/\omega \sim 22$



Other signal ? $\rightarrow K \rightarrow \pi \nu \nu$

Kaon rare decay : $K_L \rightarrow \Pi^0 VV$ and $K^+ \rightarrow \Pi^+ VV$

- Features of $K \rightarrow \pi \nu \nu$ decay
 - Rare decay : $BR_{SM} \sim 10^{-11}$



- Theoretically clean : hadronic matrix element can be removed
- Experiments are in progress

$$BR(K_L \to \pi^0 \nu \bar{\nu})_{exp} < 2.6 \times 10^{-8} \ (90\% \text{C.L.})$$
$$BR(K^+ \to \pi^+ \nu \bar{\nu})_{exp} = (1.73^{+1.15}_{-1.05}) \times 10^{-10}$$

Koto @ J-Parc, <u>running</u>



aim to measure BR($K_L \rightarrow \pi^0 vv$) around the SM sensitivity

NA62 @ CERN <u>running</u>



aim to measure O(100) events(~10% precision) BR($K^+ \rightarrow \pi^+ \nu \nu$) by 2018

Kaon rare decay : $K_L \rightarrow \Pi^0 VV$ and $K^+ \rightarrow \Pi^+ VV$

• $K \rightarrow \pi v v$ and Unitarity triangle

Determination of CPV phase (η) directly



/α

 $K^{\dagger} \rightarrow \pi^{\dagger} \nu \overline{\nu}$ (NA62)

Phase 2

Phase 1

 $K_{1} \rightarrow \pi^{0} v \overline{v}$ (KOTO

Unitarity triangle fit independently of B physics



Prospects:NA62: BR($K^+ \rightarrow \pi^+ vv$) at 10% accuracy (End 2018) KOTO: Phase 1: aim to measure BR($K_L \rightarrow \pi^0 vv$) at 10% accuracy Phase 2: BR($K_L \rightarrow \pi^0 vv$) at 10% accuracy

Correlations b/w Kaon observables in

[A.J.Buras, D.Buttazzo and R.Knegjens, JHEP1511(2015)166 A.J.Buras, JHEP1604(2016)071 C.Bobeth, A.J.Buras, A.Celis and M.Jung, 1703.04753]



	ε'/ε	ε_K	$K_L \to \pi^0 \nu \bar{\nu}$	$K^+ \to \pi^+ \nu \bar{\nu}$	$K_L \to \mu^+ \mu^-$	ΔM_K
ImΔ	*	*	*	*		*
$\operatorname{Re}\Delta$		*		*	*	*

 $K_{L} \rightarrow \pi^{0} \nu \nu$: Direct CPV and depends on only Im part \Rightarrow Strong correlation with ϵ' / ϵ



 $\begin{array}{ll} \text{Only RH (or LH) scenario} & \Rightarrow K_L \rightarrow \pi^0 \nu \nu \text{ is suppressed} \\ \text{RH + LH scenario} & \Rightarrow K_L \rightarrow \pi^0 \nu \nu \text{ can be enhanced} \\ \end{array}$



Correlations b/w Kaon observables in



Supersymmetry

supersymmetry : one of the most attractive candidates of new physics



Bounds from direct searches have been pushed up beyond 1 TeV

$$m_{\tilde{g}}, \ m_{\tilde{q}} \gtrsim 1 \text{TeV}$$

Higgs mass can be realized due to heavy stop

suggestion of high scale susy

Supersymmetry

quark & squark mass matrices cannot be diagonalized simultaneously

$$\mathcal{M}_{\tilde{u}}^{2} = \begin{pmatrix} m_{\tilde{Q}}^{2} + m_{u}^{2} + \cos 2\beta \, m_{Z}^{2} \left(\frac{1}{2} - \frac{2}{3}\sin^{2}\theta_{W}\right) & \frac{\upsilon_{2}}{\sqrt{2}}T_{U}^{*} - \mu m_{u}\cot\beta \\ \frac{\upsilon_{2}}{\sqrt{2}}T_{U}^{T} - \mu^{*}m_{u}\cot\beta & m_{\tilde{U}}^{2T} + m_{u}^{2} + \frac{2}{3}\cos 2\beta \, m_{Z}^{2}\sin^{2}\theta_{W} \end{pmatrix}$$

squark interactions depend on mixing matrix which is different from V_{CKM}



 after diagonalizing quark mass matrix, off diagonal elements of squark mass matrix give flavor violating effects

LR mixing

$$\mathcal{M}_{\tilde{u}}^2 = \operatorname{diag}(m_{\tilde{q}}^2) + m_{\tilde{q}}^2 \begin{pmatrix} \delta_{LL}^u \\ \delta_{RL}^u \end{pmatrix}$$

J

$$\begin{split} (\delta^u_{LR})_{ij} &= \frac{\frac{v_2}{\sqrt{2}} (T_U)^*_{ij}}{m_{\tilde{q}}^2} \\ \tilde{t}_R & \cdots \swarrow \tilde{t}_{LR} \quad \tilde{u}_L \end{split}$$

T_{D,U} : Trilinear scalar

coupling

Z penguin in MSSM



They do not decouple even if SUSY particles are heavy as long as a product of the mass insertion (MI) parameters (δ_{LR}^{q})* (δ_{LR}^{q}) is fixed.

not assume trilinear coupling \propto yukawa large trilinear coupling -> large contribution

 \Rightarrow vacuum stability

Vacuum stability

large trilinear couplings make the vacuum unstable



charge-color breaking (CCB) vacua

 $\Rightarrow \delta_{\tiny LR,RL}$ are constrained by vacuum stability

 The Z-penguin contributions are constrained by requiring the lifetime of the EW vacuum is longer than the age of the universe

Bound from Vacuum stability



 $m_{squark} \sim 5 \text{ TeV} \Leftrightarrow (\delta_{LR}^{u})_{i3} < 0.07$

ε'/ε

 $|(T_U)_{13}| = |(T_U)_{23}| \quad m_{\tilde{q}} \equiv m_{\tilde{Q}_i} = m_{\tilde{U}_3} \quad \tan\beta = 50$

* The CP-violating phase of Zsd coupling is taken to be maximal



 ε'/ε can be explained for the SUSY scale < 4–6 TeV

 ϵ'/ϵ vs. $K_L \rightarrow \pi^0 \nu \nu$



 $B(K_L \rightarrow \pi^0 vv)$ is predicted to be less than 60% of the SM prediction

Comments on anomaly in B sector

Z model is not favored by anomalies in b \rightarrow s transitions (P₅', R(K), R(K*),,,), which suggest negative C₉^{NP}

$$C_9^{NP} < 0 \qquad O_9 = (\bar{s}_L \gamma_\mu b_L) (\bar{\mu} \gamma^\mu \mu)$$

In Z model, it is hard to produce large C_9^{NP} due to smallness of the vector coupling to charged lepton

$$b \rightarrow s$$

 $\mu \rightarrow \mu \qquad (1 - 4 s_w^2)$

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Summary

Flavor physics offers a powerful probe of NP beyond the SM

There is 2.9 σ anomaly in direct CP violation in Kaon ε'/ε

We study SUSY model with large trilinear coupling

chargino Z penguin contribution

$$\epsilon'/\epsilon \Leftrightarrow$$
 SUSY scale < 4-6 TeV

 \Leftrightarrow BR(K_L-> $\pi\nu\nu$) < 0.6 * SM

The measurements of $K_L \rightarrow \pi^0 v v$ will be important test of this model